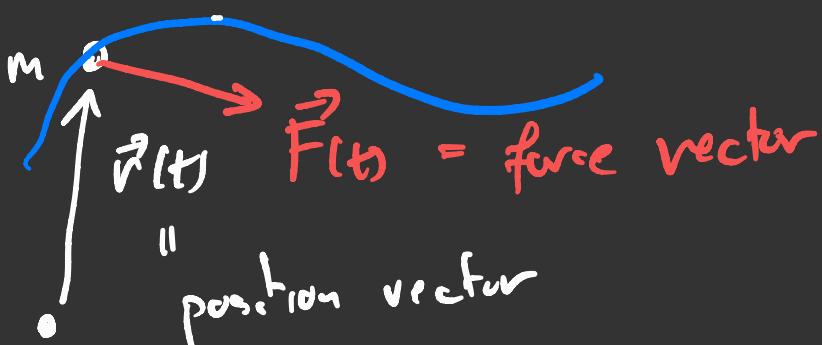


Some dynamical problems.

① 2d Newton's Law : mass m



$$\vec{F}(t) = m \vec{a}(t) = m \vec{r}''(t)$$

↑
acceleration.

Vectorial formulation of Newton's law

Examples : Motion in gravitational field

Assume the earth is horizontal and flat, and the force of gravity acts vertically:



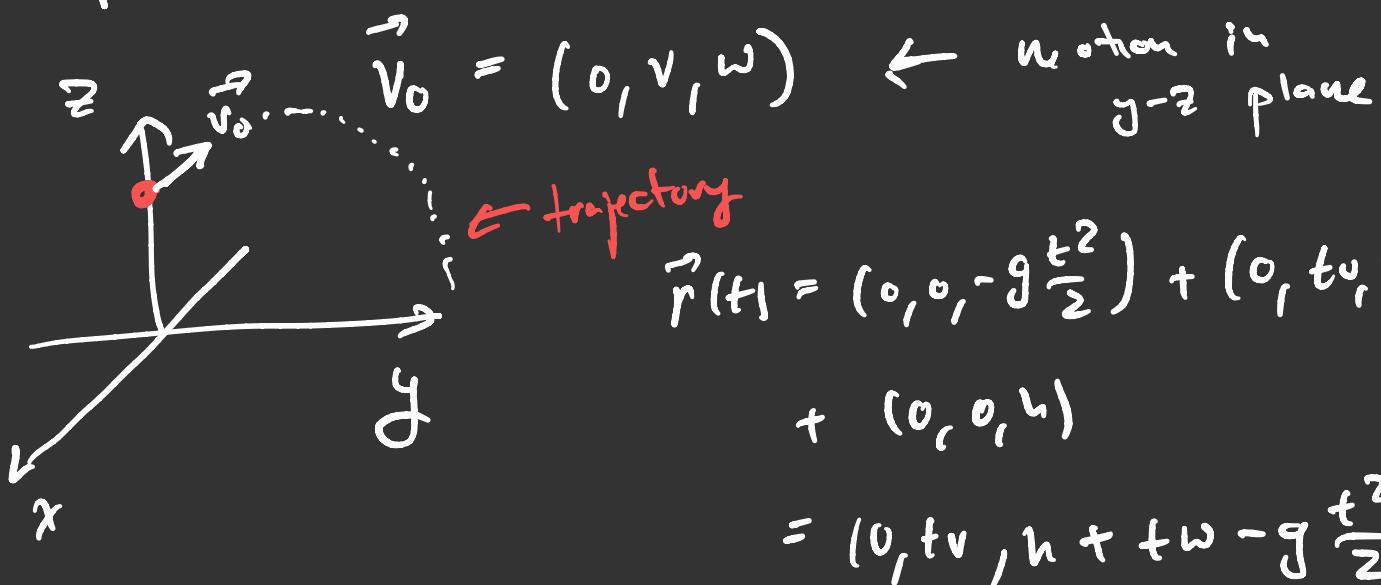
$$\vec{a} = \frac{\vec{F}}{m} = -g \hat{k} = (0, 0, -g)$$

$$\begin{aligned}\vec{a} = \frac{d}{dt} \vec{v}(t) &\Rightarrow \vec{v}(t) = \int \vec{a}(t) dt + \vec{v}_0 \\ &= - \int g \hat{k} dt + \vec{v}_0 \\ &= -gt \hat{k} + \vec{v}_0\end{aligned}$$

$$\begin{aligned}\vec{v}(t) &= \frac{d}{dt} \vec{r}(t) \\ \vec{r}(t) &= \int \vec{v}(t) dt + \vec{r}_0 \\ &= -g \frac{t^2}{2} \hat{k} + t \vec{v}_0 + \vec{r}_0\end{aligned}$$

$$\boxed{\vec{r}(t) = \vec{v}_0 + t \vec{v}_0 - g \frac{t^2}{2} \hat{k}}$$

Example: $\vec{r}_0 = (0, 0, H)$ height = H



Let's find equation for the trajectory in the y-z plane, e.g. $z = z(y)$

Note $y = t v \Rightarrow t = \frac{y}{v}$

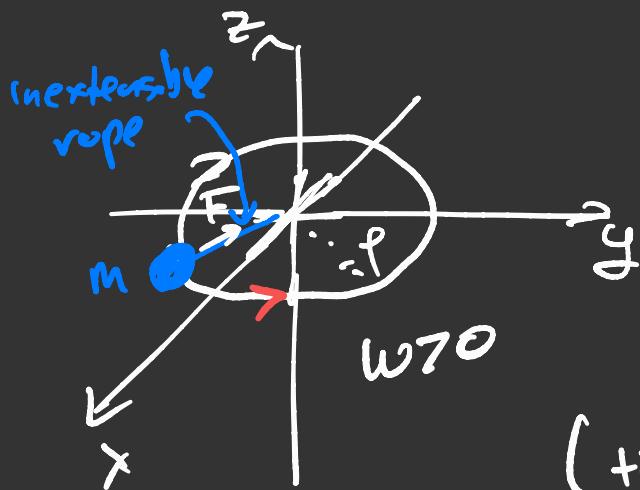
$$z = H + t w - g \frac{t^2}{2}$$

$$= H + \frac{w}{v} y - \frac{g}{2v^2} y^2$$

upside-down parabola.

2) Circular motion

$r > 0 \leftarrow$ radius
 $\omega = \text{frequency}$



$$\vec{r}(t) = (\rho \cos \omega t, \rho \sin \omega t, 0)$$

$$\text{period } T = \frac{2\pi}{\omega}$$

(time it takes to make one turn)

Rope pulls body to the center, so
 force \vec{F} is directed in. Let's find \vec{F} :

$$\vec{F} = m \vec{a} = m \vec{r}''(t)$$

while

$$\vec{r}'(t) = \omega (-\rho \sin(\omega t), \rho \cos(\omega t), 0)$$

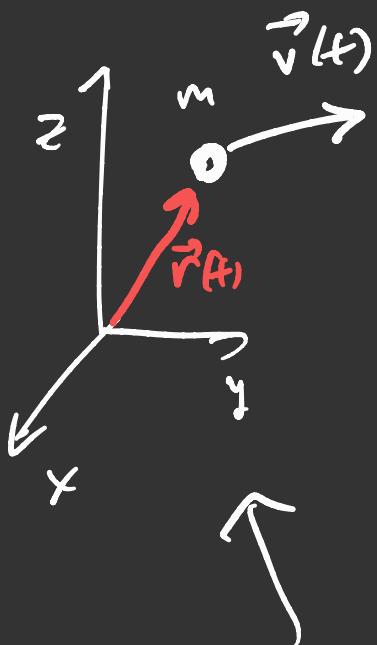
$$\begin{aligned} \vec{r}''(t) &= -\omega^2 (\rho \sin(\omega t), \rho \cos(\omega t), 0) \\ &= -\omega^2 \vec{r}(t). \end{aligned}$$

Thus

$$\vec{F}(t) = -m \omega^2 \vec{r}(t)$$

centrifugal force

3) Angular momentum and torque.



Angular Momentum

$$\vec{M} = m \vec{r} \times \vec{v}$$

perpendicular to \vec{r}, \vec{v}
and directed so that
 $(\vec{r}, \vec{v}, \vec{M})$ form a right triple

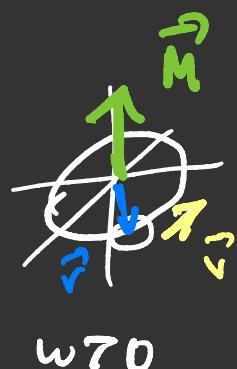
In this case, \vec{M} points inside the page.

In previous example

$$\vec{r} = (\rho \cos(\omega t), \rho \sin(\omega t), 0)$$

$$\vec{v} = (-\rho \sin(\omega t), \rho \cos(\omega t), 0)$$

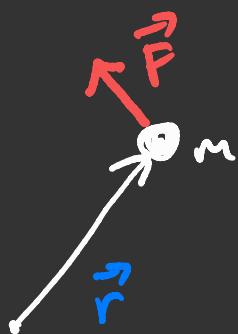
$$\vec{M} = m \vec{r} \times \vec{v} = m \begin{vmatrix} \vec{i} & \vec{j} & \vec{k} \\ \rho \cos \omega t & \rho \sin \omega t & 0 \\ -\rho \sin \omega t & \rho \cos \omega t & 0 \end{vmatrix}$$



$$= m \begin{vmatrix} \rho \cos(\omega t) & \rho \sin(\omega t) & \vec{k} \\ -\rho \sin(\omega t) & \rho \cos(\omega t) & \end{vmatrix}$$

$$= m \rho^2 \omega (\cos^2(\omega t) + \sin^2(\omega t)) \vec{k} = m \rho^2 \omega \vec{k}$$

Torque



Torque

$$\vec{\tau} = \vec{r} \times \vec{F}$$

If a force is acting on a body, it changes its angular momentum.
How does it change? Analogue of 2nd Newton's Law.

Theorem : If the mass is moving along a trajectory $\vec{r}(t)$ under the influence of force \vec{F} , then

$$\frac{d}{dt} \vec{M}(t) = \vec{\tau}(t)$$

where $\vec{M} = m \vec{r} \times \vec{v}$ is the angular momentum.

(Mass) \times Velocity = linear momentum.

\vec{M} is angular momentum

Proof:

$$\vec{M} = \vec{r} \times (m \vec{v})$$

$$\begin{aligned}
 \frac{d}{dt} \vec{M}(t) &= m \vec{r}' \times \vec{v} + m \vec{r} \times \vec{v}' \\
 &= m \vec{v} \times \vec{v} \quad \text{(cancel)} + m \vec{r} \times \vec{r}'' \\
 &= \vec{r} \times (m \vec{r}'') \\
 &= \vec{r} \times \vec{F} \\
 &= \vec{\tau}.
 \end{aligned}$$

(Leibnitz)
 $\vec{a} \times \vec{a} = 0$
 for all \vec{a}
 Pf: change
 places
 $\vec{u} \times \vec{w} = -\vec{w} \times \vec{u}$

Remarkable particular case:

If $\vec{\tau}(t) = 0$, then $\vec{M}(t) = \text{const.}$

$$\left(\frac{d}{dt} \vec{M}(t) = \vec{0} \right) \quad \text{so} \quad \vec{M}(t) = \vec{M}_0$$

What's an example of this? Central force

$$\vec{F}(t) = k(t) \vec{r}(t)$$

$$\begin{aligned}
 \vec{\tau} &= \vec{r} \times \vec{F} \\
 &= k \vec{r} \times \vec{r} = 0.
 \end{aligned}$$

Example:

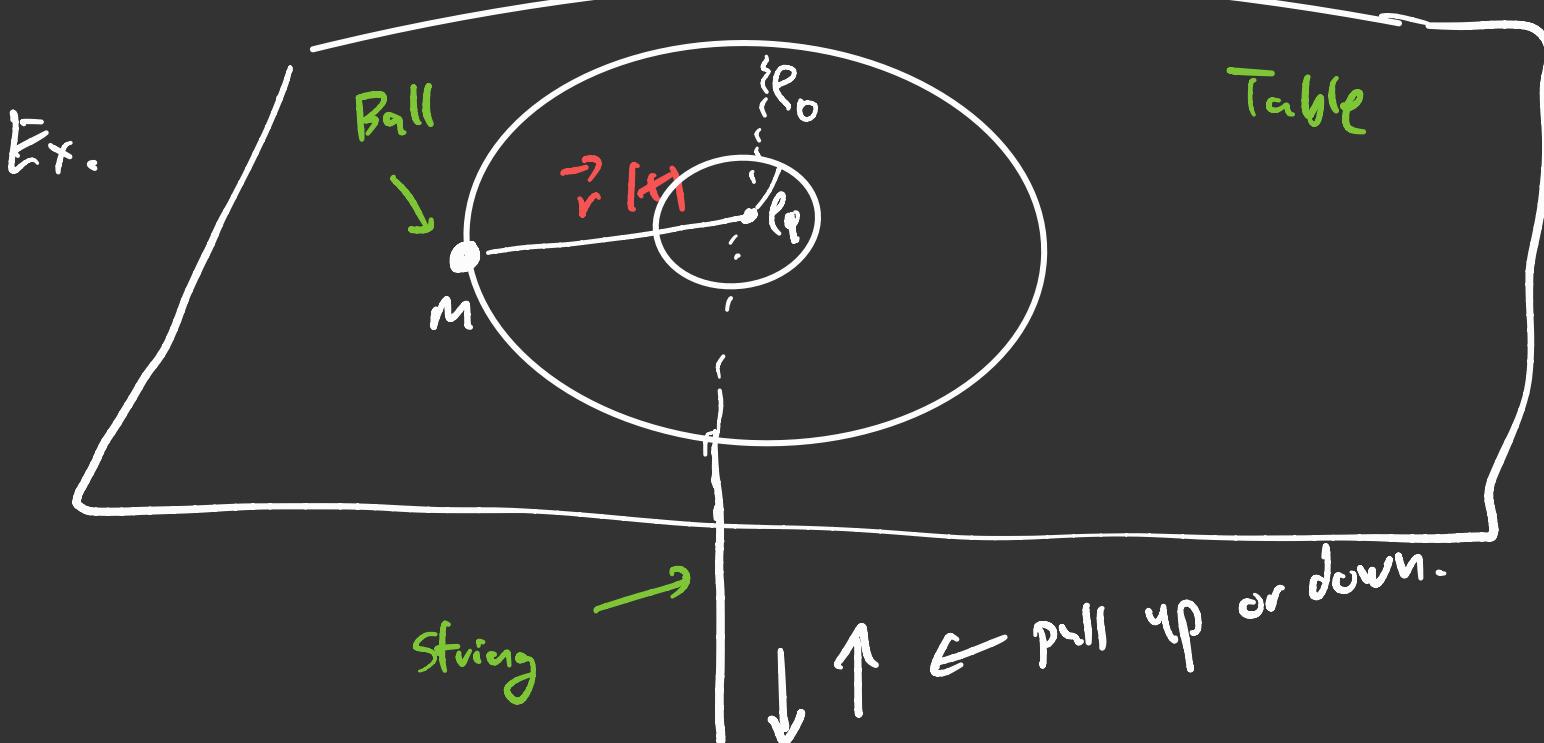
Inverse square law

$$\|\vec{F}\| \sim \|\vec{r}\|^{-2}$$

Plane & M_p

$$\vec{F} = -\frac{GM_p M_s}{\|\vec{r}\|^3} \vec{r}$$

Thus, if \vec{F} is a central force, then
 $\vec{M} = m \vec{r} \times \vec{v}$ is constant.



Thus we can prescribe $\|\vec{r}(t)\|$ arbitrarily.

$$\vec{M} = m \vec{r}(t) \times \vec{v}(t) = \text{constant}$$

$$\vec{M} = m r^2 \vec{\omega} \hat{k} \quad \text{angular momentum.}$$

$$\vec{M}_0 = m r_0^2 \vec{\omega}_0 \hat{k} \quad \vec{M}_1 = m r_1^2 \vec{\omega}_1 \hat{k}$$

$$\Rightarrow \frac{\omega_1}{\omega_0} = \frac{r_0^2}{r_1^2} \Leftrightarrow$$

$$\boxed{\omega_1 = \frac{r_0^2}{r_1^2} \omega_0}$$

If you reduce length by half, $r_1 = \frac{1}{2} r_0$, then $\omega_1 = 4 \omega_0$ and $\|v_1\| = 2 \|v_0\|$.