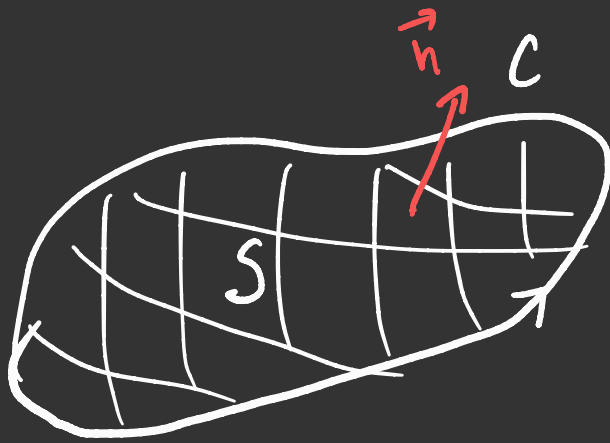


Stokes Formula



orientation is counter clockwise with respect to the normal (going in direction $\Theta \rightarrow \oplus$)

vector field $\vec{F}(\vec{r}) = (P, Q, R)$.

Stokes formula:

$$\oint_C \vec{F} \cdot d\vec{r} = \iint_S \text{curl } \vec{F} \cdot \vec{n} \, dS$$

Left hand side is defined geometrically; independent of coordinates.

The same holds for the right-hand-side.

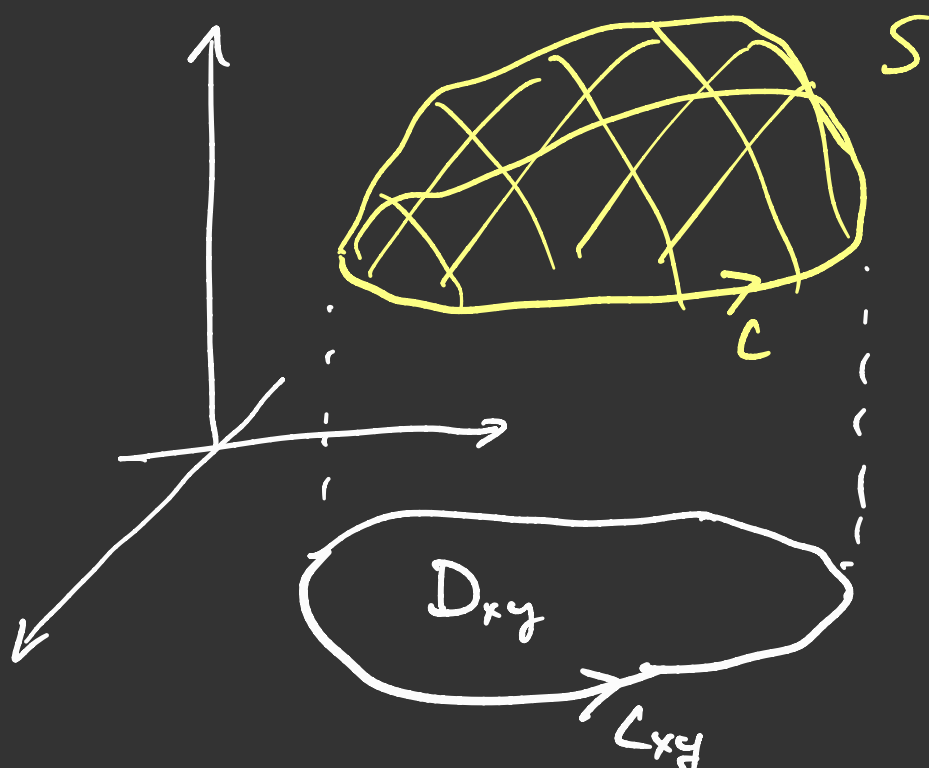
This formula is the basis for Maxwell's equations of electrodynamics. It is also the basis for some of hydrodynamics.

Why is this formula true?

We start from 2-d Stokes formula

Consider $\vec{F}(x,y,z) = (P(x,y), Q(x,y), 0)$

$$\text{curl } F = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \partial_x & \partial_y & \partial_z \\ P & Q & 0 \end{vmatrix} = (0, 0, \partial_x Q - \partial_y P)$$



$$S: z = f(x,y) \\ (x,y) \in D$$

$$\oint_C \vec{F} \cdot d\vec{r} = \oint_C P(x,y) dx + Q(x,y) dy + 0 dz$$

$$= \oint_C P(x,y) dx + Q(x,y) dy$$

Since independent of z

$$= \oint_{C_{xy}} P(x,y) dx + Q(x,y) dy$$

On other hand, since $\hat{n} = (-f_x, -f_y, 1)$ for $z = f(x,y)$

$$\iint_S \text{curl } \vec{F} \cdot \hat{n} dS = \iint_{D_{xy}} \text{curl } \vec{F}(x,y,f(x,y)) \cdot (-f_x, -f_y, 1) dA$$

$$= \iint_{D_{xy}} (Q_x(x,y) - P_y(x,y)) dA_{xy}$$

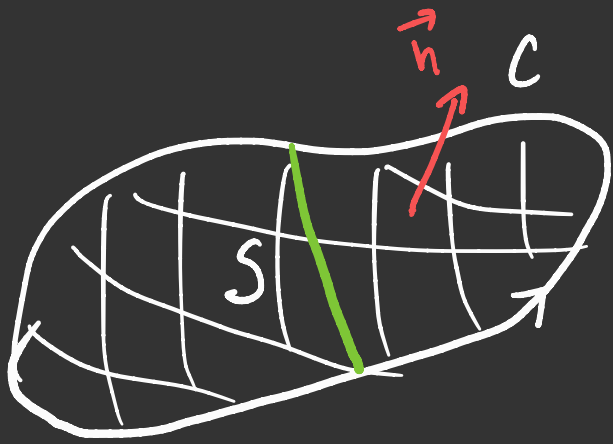
by Green's formula

$$= \oint_{C_{xy}} P dx + Q dy$$

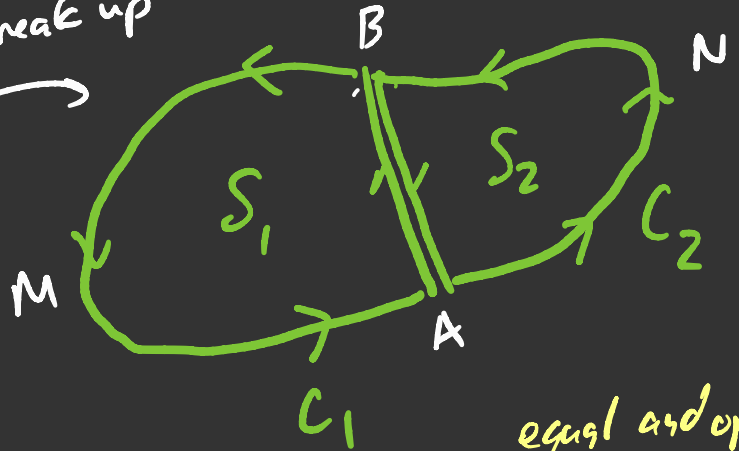
Thus, the Stokes formula for "2d fields" is proved
Applies also to:

$$\vec{G}(\vec{r}) = (0, Q(y,z), R(y,z)) \quad \vec{H}(\vec{r}) = (P(x,z), 0, Q(x,z))$$

Stokes Formula: general case



break up



equal and opposite

$$\oint_{C_1} \vec{F} \cdot d\vec{r} + \oint_{C_2} \vec{F} \cdot d\vec{r} = \int_{BMA} + \int_{AB} + \int_{ANB} + \int_{BA}$$

$$= \int_{BMA} + \int_{ANB}$$

$$= \oint_C \vec{F} \cdot d\vec{r}$$

Also:

$$\iint_S = \iint_{S_1} + \iint_{S_2}$$

by additivity.

Our goal is to prove

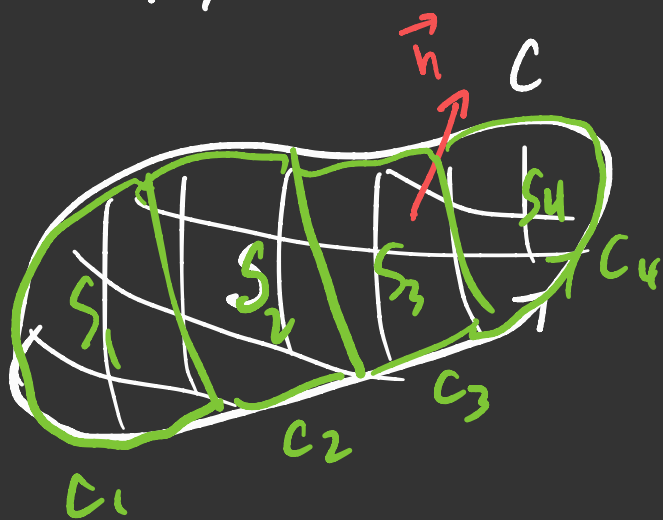
$$\oint_C \vec{F} \cdot d\vec{r} = \iint_S \text{curl } \vec{F} \cdot \vec{n} \, dS$$

We know now that, if

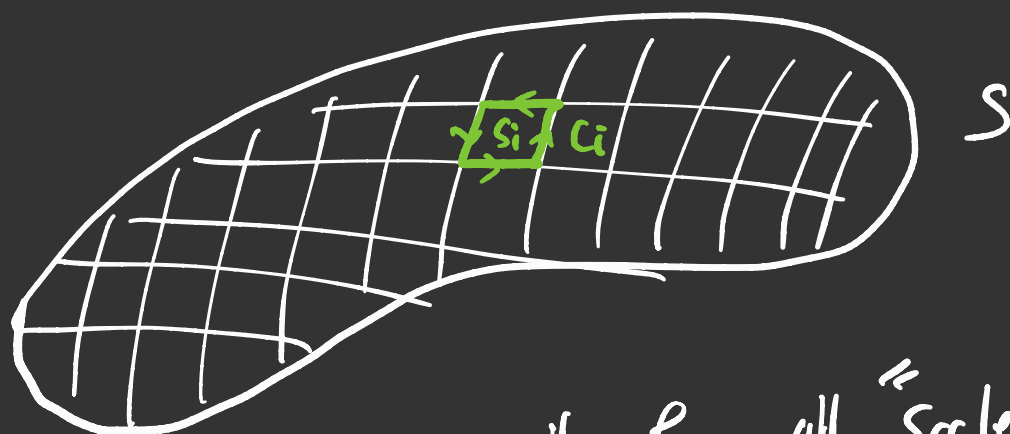
$$\oint_{C_1} \vec{F} \cdot d\vec{r} = \iint_{S_1} \text{curl } \vec{F} \cdot \vec{n} \, dS$$

$$\oint_{C_2} \vec{F} \cdot d\vec{r} = \iint_{S_2} \text{curl } \vec{F} \cdot \vec{n} \, dS$$

Then, we achieve our goal. Thus, as we did to "prove" Green's theorem, we reduced our problem to two subproblems. As before, we continue...

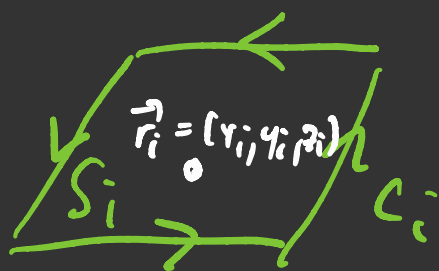


$$\oint_C \vec{F} \cdot d\vec{r} = \sum_{i=1}^N \oint_{C_i} \vec{F} \cdot d\vec{r} \quad \iint_S \text{curl} \vec{F} \cdot \vec{n} dS = \sum_{i=1}^N \iint_{S_i} \text{curl} \vec{F} \cdot \vec{n} dS$$



If we can prove it for all "scales", then by the above (ignoring controlling errors) we have proved Stokes theorem.

Consider just one small piece



Our functions P, Q, R are assumed differentiable and are thus approximated by linear functions.

$$P(x, y, z) \approx P(\vec{r}_i) + P_x(\vec{r}_i)(x - x_i) + P_y(\vec{r}_i)(y - y_i) + P_z(\vec{r}_i)(z - z_i)$$

$$= P(\vec{r}_i) - P_x(\vec{r}_i)x_i - P_y(\vec{r}_i)y_i - P_z(\vec{r}_i)z_i$$

$$+ P_x(\vec{r}_i)x + P_y(\vec{r}_i)y + P_z(\vec{r}_i)z$$

$$:= \bar{P}_i + P_x(\vec{r}_i)x + P_y(\vec{r}_i)y + P_z(\vec{r}_i)z$$

Similarly, in S_i ,

$$Q(\vec{r}) \approx \bar{Q}_i + Q_x(\vec{r}_i)x + Q_y(\vec{r}_i)y + Q_z(\vec{r}_i)z$$

$$P(\vec{r}) \approx \bar{P}_i + P_x(\vec{r}_i)x + P_y(\vec{r}_i)y + P_z(\vec{r}_i)z$$

Now for the integral:

$$\oint_{C_i} P dx + Q dy + R dz$$

$$= \oint_{C_i} (\bar{P}_i + P_x x + P_y y + P_z z) dx \\ + (\bar{Q}_i + Q_x x + Q_y y + Q_z z) dy \\ + (\bar{R}_i + R_x x + R_y y + R_z z) dz$$

Now, many terms are zero:

$$\oint_{C_i} (\bar{P}_i dx + \bar{Q}_i dy + \bar{R}_i dz)$$

$$= \oint_{C_i} \frac{\partial}{\partial x}(\bar{P}_i x) dx + \frac{\partial}{\partial y}(\bar{Q}_i y) dy + \frac{\partial}{\partial z}(\bar{R}_i z) dz$$

$$= \oint_{C_i} \nabla (\bar{P}_i x + \bar{Q}_i y + \bar{R}_i z) \cdot d\vec{r} = 0$$

← potential field

Similarly

$$\oint_{C_i} P_x x dx + Q_y y dy + R_z z dz$$

potential
field

$$= \oint_{C_i} \nabla \left(\frac{P_x}{2} x^2 + \frac{Q_y}{2} y^2 + \frac{R_z}{2} z^2 \right) \cdot d\vec{r} = 0$$

What remains is

$$\oint_{C_i} (P_y y + P_z z) dx + (Q_x x + Q_z z) dy + (R_x x + R_y y) dz$$

Rearranging

$$= \oint_{C_i} (P_y y dx + Q_x x dy)$$

$$+ \oint_{C_i} (P_z z dx + R_x x dz)$$

$$+ \oint_{C_i} (Q_z z dy + R_y y dz)$$

Consider

Simple vector field!

$$\oint_{C_i} P_y y dx + Q_x x dy = \int_{C_i} (P_y y, Q_x x, 0) \cdot d\vec{r}$$

By the 2d Stokes formula, this is equal to

$$= \iint_{S_i} \text{curl} (P_y y, Q_x x, 0) \cdot \vec{n} dS$$

For the same reason

$$\oint_{C_i} (P_z z, 0, R_x x) \cdot d\vec{r} = \iint_{S_i} \text{curl} (P_z z, 0, R_x x) \cdot \vec{n} dS$$

$$\oint_{C_i} (0, Q_z z, R_y y) \cdot d\vec{r} = \iint_{S_i} \text{curl} (0, Q_z z, R_y y) \cdot \vec{n} dS$$

Thus, combining, for the little piece $\diamond S_i$ we proved:

$$\oint_{C_i} \vec{F} \cdot d\vec{r} = \iint_{S_i} \text{curl} \vec{F} \cdot \vec{n} dS$$

Logical scheme of "proof."

- ① derive theorem for simple case of a two-dimensional vector field by an application of Green's theorem.
- ② Reduce our problem to consideration to small pieces, just one "scale".
- ③ Within a scale, we approximate of functions P, Q, R by linear functions. We then see some terms disappear. What remains are three terms which represent 2d vector fields and use step ① to conclude.