

Four components of d on almost complex manifolds*

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Abstract

The exterior differential d on complex-valued differential forms of complex manifolds decomposes into the Cauchy-Riemann operator and its complex conjugate. Meanwhile on almost complex manifolds, the exterior d in general has two extra components, thus decomposes into four operators. In this talk, I will introduce these operators and discuss the structure of the (graded) associative algebra generated by these four components of d , subject to relations deduced from d squaring to zero. Then I will compare this algebra to the corresponding one in the complex (i.e. integrable) case, we shall see they are very different strictly speaking but similar in a weak sense (quasi-isomorphic). This is based on joint work with Shamuel Aueyung and Jin-Cheng Guu[AGH22].

1 Decomposition of d on almost complex manifolds

Given a complex manifold M of complex dimension n , locally it admits complex coordinates $z^i = x^i + \sqrt{-1}y^i$, $i = 1, \dots, n$. Its tangent space locally is spanned by $\partial/\partial x^i, \partial/\partial y^i$ and its complexified tangent space is spanned by $\partial/\partial z^i$ and $\partial/\partial \bar{z}^i$. Let $T^{1,0}M$ denote the span of $\partial/\partial z^i$'s and $T^{0,1}M$ the span of $\partial/\partial \bar{z}^i$'s. Thus we have

$$TM \otimes \mathbb{C} = T^{1,0}M \oplus T^{0,1}M$$

The two subspaces $T^{1,0}$ and $T^{0,1}$ of complexified tangent bundle can be characterized without coordinates as follows. Consider the linear map $J : TM \rightarrow TM$ defined by

$$J(\partial/\partial x^i) = \partial/\partial y^i, \quad J(\partial/\partial y^i) = -\partial/\partial x^i$$

It turns out J is globally defined and satisfies $J^2 = -1$. Then $T^{1,0}$ and $T^{0,1}$ are eigenspaces of J with eigenvalue $\pm\sqrt{-1}$.

Dually, the complexified cotangent space admits a decomposition

$$T^*M \otimes \mathbb{C} = T_{0,1}^*M \oplus T_{1,0}^*M$$

where locally $T_{1,0}^*$ and $T_{0,1}^*$ are spanned by dz^i 's and $d\bar{z}^i$'s respectively. The above decomposition yields a decomposition

$$\Lambda^k T^*M \otimes \mathbb{C} = \bigoplus_{p+q=k} \Lambda^p T_{1,0}^*M \otimes \Lambda^q T_{0,1}^*M$$

and consequently the Hodge decomposition¹

$$\Omega^k(M) = \bigoplus_{p+q=k} \Omega^{p,q}(M)$$

where Ω^k means group of complex valued differential k -forms and $\Omega^{p,q}$ the group of smooth sections of $\Lambda^p T_{1,0}^*M \otimes \Lambda^q T_{0,1}^*M$, usually referred to as the (p, q) -forms.

Note that the Hodge decomposition of differential forms *does not* rely on the complex coordinates on M , but rather only depends on the linear map $J : TM \rightarrow TM$. We call a (real) manifold M equipped with such a linear map an almost complex manifold and such J is called its almost complex structure.

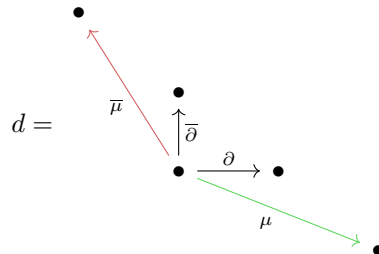
*This is an informal note, citations are not carefully put

¹this is not a standard terminology

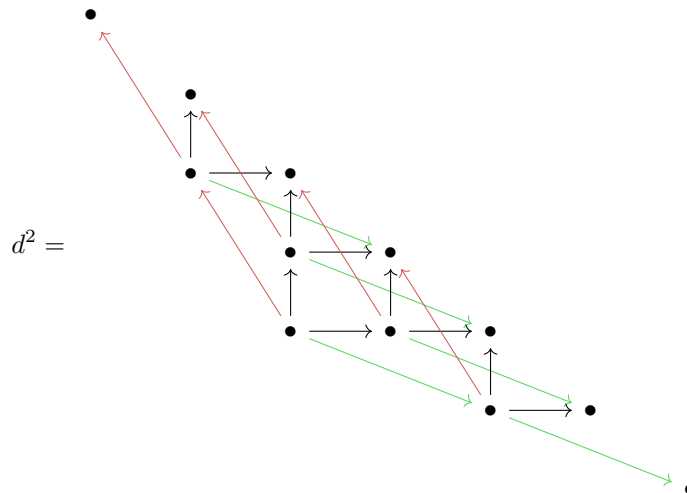
Now we would like to understand the chain complex $(\Omega^\bullet(M), d)$ with respect to the Hodge decomposition. First of all, $\Omega^\bullet(M)$ is a (graded) algebra when equipped with wedge product, called the de Rham algebra, on which d is a derivation. Since the de Rham algebra is generated by 0-forms (smooth functions) and 1-forms, it suffices to understand the action of d on 0-forms and 1-forms. It follows that d decomposes, with respect to the Hodge decomposition as

$$d = \bar{\mu} + \bar{\partial} + \partial + \mu$$

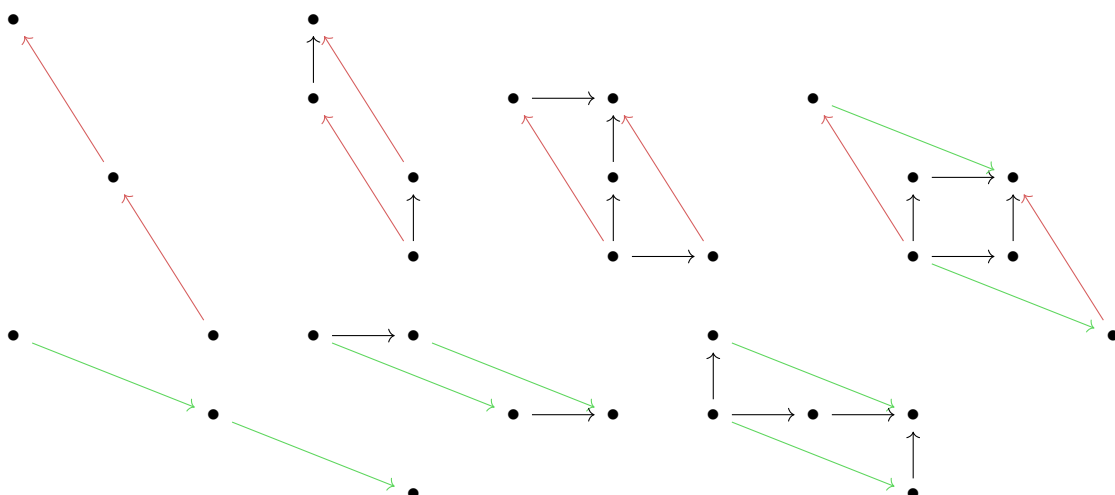
It is better to write this pictorially as



Then d^2 pictorially is



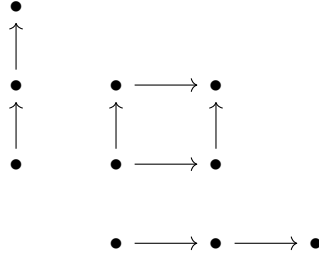
The equation $d^2 = 0$ is equivalent to the vanishing of the following diagrams



If our almost complex manifold (M, J) actually admits complex coordinates, i.e. complex/holomorphic, then it is well-known that d has only two components $\bar{\partial}, \partial$:

$$d = \begin{array}{c} \bullet \\ \uparrow \bar{\partial} \\ \bullet \xrightarrow{\partial} \bullet \end{array}$$

subject to relations given by the vanishing of



2 Algebra generated by the components of d

From the above discussion, the relations among the four components of d looks to be rather complicated. We would like to clarify what the relations “mean”. More precisely, consider the algebra

$$A = \frac{\text{free associative algebra generated by } \bar{\mu}, \bar{\partial}, \partial, \mu}{\text{relations given above by } d^2 = 0}$$

Question. What is the structure of the algebra A ?

This is a somewhat vague question, a more precise one can be: what is the dimension of A in each degree? Our answer to the above question is based on a sequence of observations.

Observation 1. The relations among $\bar{\mu}, \bar{\partial}, \partial, \mu$ can be written using (graded) Lie brackets.

Indeed, $d^2 = 0$ is equivalent to $[d, d] = 2d^2 = 0$. The relations can be written as

$$\begin{aligned} [\bar{\mu}, \bar{\mu}] &= 0, & [\mu, \mu] &= 0; \\ [\bar{\mu}, \bar{\partial}] &= 0, & [\mu, \partial] &= 0; \\ [\bar{\mu}, \partial] &= -\frac{1}{2}[\bar{\partial}, \bar{\partial}], & [\mu, \bar{\partial}] &= -\frac{1}{2}[\partial, \partial]; \\ [\bar{\mu}, \mu] &= -[\bar{\partial}, \partial]. \end{aligned} \tag{1}$$

Theorem 1. $A = U\mathfrak{g}$ where \mathfrak{g} is the graded Lie algebra generated by $\bar{\mu}, \bar{\partial}, \partial, \mu$ and U means universal enveloping algebra.

Proof. Observe A is a primitively generated Hopf algebra in which $\bar{\mu}, \bar{\partial}, \partial, \mu$ are primitive. Then apply Milnor-Moore. \square

Therefore, determining the structure of A amounts to determining that of \mathfrak{g} . Note that \mathfrak{g} is much smaller than A in the following sense. If β_1, β_2, \dots are Betti numbers of \mathfrak{g} , that is $\beta_d = \dim \mathfrak{g}_d$ the degree d subspace of \mathfrak{g} , then those of A can be read off from its Poincaré series

$$\prod_i \frac{(1 + x^{2i+1})^{\beta_{2i+1}}}{(1 - x^{2i})^{\beta_{2i}}}$$

Observation 2. There are no relations between *only* $\bar{\partial}$ and ∂ in (1).

Theorem 2. The Lie subalgebra \mathfrak{h} of \mathfrak{g} generated by $\bar{\partial}$ and ∂ is free.

The proof of this theorem actually requires an understanding of the role played by $\bar{\mu}, \mu$ —they are derivations on \mathfrak{h} and \mathfrak{g} is a “semi-direct product” of \mathfrak{h} and $\bar{\mu}, \mu$.

Observation 3. (The adjoint actions of) $\bar{\mu}$ and μ take $\bar{\partial}, \partial$ into \mathfrak{h} .

Theorem 3. \mathfrak{h} is a Lie ideal of \mathfrak{g} . Moreover, $\mathfrak{g} = \mathfrak{h}$ in degrees ≥ 2 .

Proof. The second assertion follows from an easy induction. \square

At this point, we have a rather satisfying understanding of \mathfrak{g} through the exact sequence

$$0 \rightarrow \mathfrak{h} \rightarrow \mathfrak{g} \rightarrow \mathfrak{g}/\mathfrak{h} \rightarrow 0$$

Here \mathfrak{h} is the free Lie algebra generated by $\bar{\partial}, \partial$, and $\mathfrak{g}/\mathfrak{h}$ is an abelian Lie algebra generated by $\bar{\mu}, \mu$. The extension is through the derivation action of $\bar{\mu}, \mu$ on \mathfrak{h} .

Now let us compare the situation to the complex case. There is a corresponding algebra A_{hol} and Lie algebra \mathfrak{g}_{hol} where

$$\mathfrak{g}_{hol} = \frac{\text{free graded Lie algebra generated by } \bar{\partial}, \partial}{[\bar{\partial}, \bar{\partial}] = [\bar{\partial}, \partial] = [\partial, \partial] = 0}$$

That is to say, \mathfrak{g}_{hol} is the abelian Lie algebra generated by $\bar{\partial}, \partial$ and thus A_{hol} is the exterior algebra generated by $\bar{\partial}, \partial$.

The natural quotient map

$$\mathfrak{g} \rightarrow \mathfrak{g}_{hol}$$

has a gigantic kernel. Indeed, the difference is roughly (or precisely in degrees ≥ 2) the difference between free Lie algebra and abelian Lie algebra.

3 Cohomology

The Lie algebra \mathfrak{g} is equipped with a natural differential $[d, -]$. As we have pointed out, \mathfrak{g} and \mathfrak{g}_{hol} are very different, so the following fact is surprising to us:

Theorem 4. The quotient map $\mathfrak{g} \rightarrow \mathfrak{g}_{hol}$ is a quasi-isomorphism.

Let us unwind the statement. Notice that \mathfrak{g}_{hol} is concentrated in degree 1 and $[d, -]$ vanishes on \mathfrak{g}_{hol} as it is abelian, we see the cohomology of \mathfrak{g}_{hol} is itself, concentrated in degree 1. So the theorem actually says $H^k(\mathfrak{g}, [d, -])$ vanishes for $k \geq 2$. Since one can explicitly verify the quotient map yields an isomorphism on H^1 , the essence of this theorem is the vanishing of higher cohomologies.

The geometric meaning of the above theorem is, any differential operator of order ≥ 2 constructed from $\bar{\mu}, \bar{\partial}, \partial, \mu$, if commutes with d , is of the form $[d, D]$ for some operator D .

Proof sketch. From the exact sequence

$$0 \rightarrow \mathfrak{h} \rightarrow \mathfrak{g} \rightarrow \mathfrak{g}/\mathfrak{h} \rightarrow 0$$

The vanishing of higher (degree > 2) cohomologies of \mathfrak{g} is equivalent to the same vanishing for \mathfrak{h} . Then by the generalized Frölicher spectral sequence of Cirici-Wilson [CW21], it suffices to prove $H^k(\mathfrak{h}, [\bar{\mu}, -]) = 0$ for $k > 2$ and $H^2(\mathfrak{g}, [\bar{\mu}, -]) = 0$. Then one passes to universal enveloping algebra of \mathfrak{h} , i.e. the free associative algebra generated by $\bar{\partial}, \partial$ and apply induction. \square

Proof of vanishing $\bar{\mu}$ -cohomology for \mathfrak{h} . First of all, we note $\text{ad}_{\bar{\mu}} \bar{\partial} = 0$ and

$$\text{ad}_{\bar{\mu}}[\partial, \bar{\partial}] = [\text{ad}_{\bar{\mu}} \partial, \bar{\partial}] - [\partial, \text{ad}_{\bar{\mu}} \bar{\partial}] = -\frac{1}{2}[[\bar{\partial}, \bar{\partial}], \bar{\partial}] = 0.$$

We leave it to the reader to check when $k = 1, 2$, $H^k(\mathfrak{h}, \text{ad}_{\bar{\mu}})$ is one-dimensional and spanned by (the equivalence class of) $\bar{\partial}$ and $[\partial, \bar{\partial}]$ respectively. Next we observe $H(\mathfrak{h}, \text{ad}_{\bar{\mu}})$ is a Lie algebra and $H^1(\mathfrak{h}) \oplus H^2(\mathfrak{h})$ forms an *abelian* Lie subalgebra since $[\bar{\partial}, \bar{\partial}] = -2[\bar{\mu}, \partial]$ is $\text{ad}_{\bar{\mu}}$ -exact, and $[\bar{\partial}, [\partial, \bar{\partial}]] = 0$, $[[\partial, \bar{\partial}], [\partial, \bar{\partial}]] = 0$ by Jacobi identities.

Now consider the universal enveloping algebra $UH(\mathfrak{h}, \text{ad}_{\bar{\mu}})$ of $H(\mathfrak{h}, \text{ad}_{\bar{\mu}})$. It contains the universal enveloping algebra of the abelian Lie subalgebra $H^1(\mathfrak{h}) \oplus H^2(\mathfrak{h})$, which is the free graded commutative algebra $\Lambda(\bar{\partial}, [\partial, \bar{\partial}])$ generated by $\bar{\partial}$ and $[\partial, \bar{\partial}]$. Meanwhile, since the universal enveloping algebra functor commutes with cohomology, we have $UH(\mathfrak{h}, \text{ad}_{\bar{\mu}}) = H(U\mathfrak{h}, \text{ad}_{\bar{\mu}})$ where $\text{ad}_{\bar{\mu}}$ on $U\mathfrak{h}$ is the extended adjoint action. So we get

$$\Lambda(\bar{\partial}, [\partial, \bar{\partial}]) \subset H(U\mathfrak{h}, \text{ad}_{\bar{\mu}}).$$

By Poincaré-Birkhoff-Witt theorem, our proposition is equivalent to $\Lambda(\bar{\partial}, [\partial, \bar{\partial}]) = H(U\mathfrak{h}, \text{ad}_{\bar{\mu}})$. This equality clearly holds in degrees ≤ 2 . We plan to prove this by induction on degree, but we need to make some preparations.

For simplicity of notation, denote $B = U\mathfrak{h}$, which is the free tensor algebra on $\bar{\partial}, \partial$. Under the isomorphism

$$\phi : B_{k-1} \oplus B_{k-1} \cong B_k, (x, y) \mapsto \partial x + \bar{\partial} y,$$

the differential $\text{ad}_{\bar{\mu}}|_{B_k}$ can be written as the matrix

$$\text{ad}_{\bar{\mu}}|_{B_k} \cong \begin{pmatrix} -\text{ad}_{\bar{\mu}}|_{B_{k-1}} & 0 \\ -\bar{\partial}|_{B_{k-1}} & -\text{ad}_{\bar{\mu}}|_{B_{k-1}} \end{pmatrix} \quad (2)$$

by using the relations (1). To see this, we compute for $x, y \in B_{k-1}$

$$\begin{aligned} [\bar{\mu}, \partial x + \bar{\partial} y] &= \bar{\mu} \partial x - (-1)^k \partial x \bar{\mu} + \bar{\mu} \bar{\partial} y - (-1)^k \bar{\partial} y \bar{\mu} \\ &= -\bar{\partial}^2 x - \partial \bar{\mu} x - (-1)^k \partial x \bar{\mu} - \bar{\partial} \bar{\mu} y - (-1)^k \bar{\partial} y \bar{\mu} \\ &= -\partial (\bar{\mu} x - (-1)^{k-1} x \bar{\mu}) - \bar{\partial} (\bar{\partial} x + \bar{\mu} y - (-1)^{k-1} y \bar{\mu}) \\ &= -\partial [\bar{\mu}, x] - \bar{\partial} (\bar{\partial} x + [\bar{\mu}, y]). \end{aligned}$$

In particular, by setting $x = 0$, we see $\text{ad}_{\bar{\mu}}$ skew commutes with $\bar{\partial}$. This means both $\pm \bar{\partial}$ are morphisms of cochain complexes $\pm \bar{\partial} : B_{\bullet} \rightarrow B_{\bullet}[1]$, where $B_{\bullet} = (B, \text{ad}_{\bar{\mu}})$. Moreover, (2) shows the mapping cone of $-\bar{\partial}$ is isomorphic to $B_{\bullet}[2]$ by ϕ . Then the inclusion of $B_{\bullet}[1]$ into the mapping cone of $-\bar{\partial}$ is identified with $\bar{\partial} : B_{\bullet}[1] \rightarrow B_{\bullet}[2]$, and the projection from the mapping cone of $-\bar{\partial}$ onto $B_{\bullet}[1]$ is identified with $\delta : B_{\bullet}[2] \rightarrow B_{\bullet}[1]$ which takes $\partial x + \bar{\partial} y$ to x . It follows we have an exact triangle

$$B_{\bullet} \xrightarrow{-\bar{\partial}} B_{\bullet}[1] \xrightarrow{\bar{\partial}} B_{\bullet}[2] \xrightarrow{\delta} B_{\bullet}[1].$$

This exact triangle induces a long exact sequence in cohomology

$$\dots \rightarrow H^{k-2}(B_{\bullet}) \xrightarrow{-\bar{\partial}} H^{k-1}(B_{\bullet}) \xrightarrow{\bar{\partial}} H^k(B_{\bullet}) \xrightarrow{\delta} H^{k-1}(B_{\bullet}) \xrightarrow{-\bar{\partial}} H^k(B_{\bullet}) \rightarrow \dots$$

Now we can inductively prove $H(B_{\bullet}) = \Lambda(\bar{\partial}, [\partial, \bar{\partial}])$. We note $\Lambda(\bar{\partial}, [\partial, \bar{\partial}])$ is one-dimensional in each degree and spanned by powers of $[\partial, \bar{\partial}]$ and $\bar{\partial}$ times powers of $[\partial, \bar{\partial}]$. Assume the desired equality is proved in degrees $< k$. Observe that $\bar{\partial}$ vanishes on $(\Lambda(\bar{\partial}, [\partial, \bar{\partial}]))^{\text{odd}}$, so if k is even then from the above long exact sequence we have $\delta : H^k(B_{\bullet}) \rightarrow H^{k-1}(B_{\bullet})$ is an isomorphism. On the other hand, if k is odd, then $\bar{\partial} : H^{k-1}(B_{\bullet}) \rightarrow H^k(B_{\bullet})$ is monic since $\bar{\partial}$ takes $(\Lambda(\bar{\partial}, [\partial, \bar{\partial}]))^{\text{even}}$ injectively into $(\Lambda(\bar{\partial}, [\partial, \bar{\partial}]))^{\text{odd}} \subset H^{\text{odd}}(B_{\bullet})$ and thus the above long exact sequence implies $\bar{\partial} : H^{k-1}(B_{\bullet}) \rightarrow H^k(B_{\bullet})$ is an isomorphism. So in either case we have $H^k(B_{\bullet}) \cong H^{k-1}(B_{\bullet})$, and in particular $H^k(B_{\bullet})$ is one-dimensional. But $H^k(B_{\bullet}) \supset (\Lambda(\bar{\partial}, [\partial, \bar{\partial}]))^k$, therefore $H^k(B_{\bullet})$ must be equal to $(\Lambda(\bar{\partial}, [\partial, \bar{\partial}]))^k$. This finishes the inductive step and thus completes the proof. \square

Proof of vanishing of H^2 for \mathfrak{g} . One can compute $H^2(\mathfrak{g}, [\bar{\mu}, -]) = 0$ and apply the spectral sequence of Cirici-Wilson. \square

Observation 4. There are many inner differentials on \mathfrak{g} .

For example $[\bar{\mu}, -]$ and $[\mu, -]$ are inner differentials on \mathfrak{g} .

Question.² Find all inner differentials on \mathfrak{g} .

This amounts to solving the equation $[a, a] = 0$ for $a \in \mathfrak{g}_1$. Writing $a = x\bar{\mu} + y\bar{\partial} + z\partial + w\mu$ we have

$$\begin{aligned} xz - y^2 &= 0, \\ yw - z^2 &= 0, \\ xw - yz &= 0. \end{aligned}$$

Denote the set of solutions $MC(\mathfrak{g}) = \{a \in \mathfrak{g}_1 : [a, a] = 0\}$.

²This question is also addressed in [TT20].

Theorem 5. $\mathbb{P}MC(\mathfrak{g})$ is the twisted cubic in \mathbb{P}^3 , parametrized by

$$\mathbb{P}^1 \rightarrow \mathbb{P}MC(\mathfrak{g}), [s:t] \mapsto d_{s,t} = s^3\bar{\mu} + s^2t\bar{\partial} + st^2\partial + t^3\mu$$

Proposition 6. For $(s, t) \neq (0, 0)$, $H^1(\mathfrak{g}, [d_{s,t}, -])$ is spanned by $\frac{\partial}{\partial s}d_{s,t}$ and $\frac{\partial}{\partial t}d_{s,t}$.

Proof. $[d_{s,t}, d_{s,t}] = 0$ implies $0 = \frac{\partial}{\partial s}[d_{s,t}, d_{s,t}] = 2[d_{s,t}, \frac{\partial}{\partial s}d_{s,t}]$. Similar for t . The fact that they span can be proved by explicit computation or deduced from deformation theory (cf. Goldman-Milson). \square

Corollary 7. For $st \neq 0$, $H^1(\mathfrak{g}, [d_{s,t}, -])$ is spanned by $d_{s,t}$ and

$$d_{s,t}^J = \sqrt{-1}(3s^3\bar{\mu} + s^2t\bar{\partial} - st^2\partial - 3t^3\mu).$$

The operator $d^J = \sqrt{-1}(3\bar{\mu} + \bar{\partial} - \partial - 3\mu)$ is in fact the commutator $[d, J]$ where J is extended as a derivation on the de Rham algebra.

Corollary 8. The quotient map $\mathfrak{g} \rightarrow \mathfrak{g}_{hol}$ is a quasi-isomorphism with respect to $[d_{s,t}, -]$ for $st \neq 0$.

Proof. $d_{s,t}$ is conjugate to d by multiplication by $s^{p+2q}t^{2p+q}$ on bidegree (p, q) subspace. \square

4 Interpretations and comments

4.1 Deformation theory

It is expected from deformation theory that the first cohomology $H^1(\mathfrak{g}, [d_{s,t}, -])$ is the Zariski tangent space of $MC(\mathfrak{g})$ at $d_{s,t}$. Also the quasi-isomorphism should yield a bijection between $MC(\mathfrak{g}, [d, -])$ and $MC(\mathfrak{g}_{hol}, [d, -])$. But the former is isomorphic to $MC(\mathfrak{g})$ and the latter is isomorphic to \mathbb{A}^2 .

4.2 Rational homotopy theory

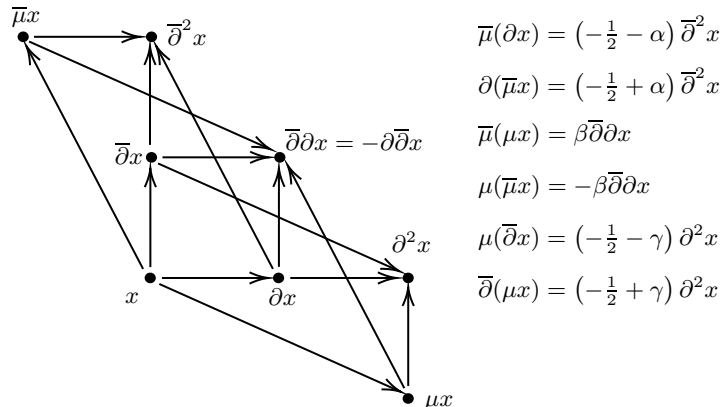
The subalgebra \mathfrak{h} is isomorphic to the homotopy Lie algebra of $\mathbb{C}\mathbb{P}^1 \vee \mathbb{C}\mathbb{P}^1$ and the subalgebra of \mathfrak{g}_{hol} generated by $\bar{\partial}, \partial$, which is \mathfrak{g}_{hol} itself, is isomorphic to the homotopy Lie algebra of $\mathbb{C}\mathbb{P}^\infty \times \mathbb{C}\mathbb{P}^\infty$. So the difference between \mathfrak{g} and \mathfrak{g}_{hol} is captured by higher dimensional cells of $\mathbb{C}\mathbb{P}^\infty \times \mathbb{C}\mathbb{P}^\infty$. For example, the 4-dimensional cell corresponds to $[\partial, \bar{\partial}]$. The quotient

$$\mathfrak{g}/([\partial, \bar{\partial}] = 0)$$

is 6-dimensional whose subalgebra generated by $\partial, \bar{\partial}$ is isomorphic to the homotopy Lie algebra of $\mathbb{C}\mathbb{P}^1 \times \mathbb{C}\mathbb{P}^1$.

4.3 Representation theory

The ultimate goal is through the understanding of A and its representations, we can understand almost complex manifolds better. But the representation theory for A or equivalently for \mathfrak{g} looks to be very complicated.



This is a faithful representation of $\mathfrak{g}/[\partial, \bar{\partial}]$. We remark that $\mathfrak{g}/[\partial, \bar{\partial}]$ contains two (or three) copies of graded Heisenberg Lie algebra.

One might wish to design certain cohomology groups that are able to distinguish different representations of \mathfrak{g} . My suggestion is to consider the role of d^J and define

$$H_{BC}^\bullet = \frac{\ker d}{\text{im } dd^J}$$

and its dual notion

$$H_A^\bullet = \frac{\ker dd^J}{\text{im } d}$$

The catch of this definition is that H_{BC} is graded, real, but not multiplicative because $(d^J)^2 \neq 0$ in general. One might say the failure of H_{BC} to be multiplicative measures the failure of J to be integrable.

We note there are other definitions of Bott-Chern and Aeppli cohomology for almost complex manifolds. See [CPS21] for example.

One probably should justify either of the definition by relating those groups to geometry like Bott and Chern did, or to quantities from representations of \mathfrak{g} .

References

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