

Dynamics of Foliations by Riemann Surfaces.

Example

In $\mathbb{C}^k$

$$z'_j(z) = P_j(z(z)) \quad 1 \leq j \leq k$$

P_j polynomials.

Singular pts in $\mathbb{C}^k$ ($f=0$)

Locally away from Sing.



The foliation can be extended
to $\mathbb{P}^k$.

It is the image of the

foliation of $\mathbb{C}^{k+1}$ given by

$$F(z) = \sum_{j=0}^k F_j(z) \frac{\partial}{\partial z_j} \quad \deg F_j = d$$

F_j homogeneous polynomial.

Foliations are parametrized

$$F = [F_0 : \dots : F_k]. \quad \deg F_j = d_j$$

Assume $\bigcap F_j^{-1}(0) = 0$.

$$\text{Singular points} \quad \# \text{Sing } F = \frac{d^{k+1} - 1}{d - 1}$$

Can also consider atlas.

$$(z, t) \rightarrow (z', t') = (h(z, t), t'(t))$$

Fold open Zariski set in some $\mathbb{P}^N$

Always sing pt.

$$DZ(0) = \{\lambda_1, \dots, \lambda_k\}.$$

Sing hyperbolic $\mathbb{R} \nsubseteq \mathbb{R} \quad \forall (i, j)$
if λ_i / λ_j

Clifford.

3)

Global behavior
of leaves?

Ilyashenko. Hudai-Veronov.
when Line at infinity invariant
in $\mathbb{P}^2$
generically leaves are dense

JOVANOVIC.

Generic foliations in $\mathbb{P}^k$
have no algebraic leaves.

$Hyp^d(\mathbb{P}^2)$. Foliations with
only hyperbolic singularities and
without algebraic leaves
is a real Zariski dense set
in $Fold(\mathbb{P}^2)$.

So analysis cannot start at
an alg. leave.

1) Leaves ?

Lins - Neto , Glatzyuk

If the singularities
are non degenerate ,
Then leaves are
covered by unit disc.

$$\Delta \xrightarrow{\downarrow \phi} \mathbb{D}$$

Gomez-Mout , Candel

If singularities are
hyperbolic
Poincaré metric is
continuous.

$$\omega_p(z)$$

$$\eta(z) = \frac{\|v\|}{\|v\|_p} \rightarrow 0 \text{ at sing.}$$

5) Discrete dyn. syst
Invariant measure



Foliations or laminations

?



measure invariant
by holonomy.

$$T = \int_B [v_\alpha] d\mu(\alpha).$$

μ inv. by holonomy $(=) T_{\text{closed}}$

Sullivan.

In general they don't
exist.

Harmonic measures

$$T = \int [V_\alpha] h_\alpha d\mu(\alpha) \quad \left. \vphantom{\int} \right\} \begin{array}{l} \partial\bar{\partial}\text{-closed} \\ \text{current.} \end{array}$$

$$m_T = T \wedge \beta.$$

$h_\alpha > 0$ harmonic on V_α .

Lucy Garnett: in Riemannian case (without sing.) by sliding sets of μ -measure zero ind. by sliding.

(X, E, L) E closed set.

"small"

$\chi|_E$ Lamination by Riemann

Surfaces.

Then there is a $(1,1) \geq 0$ current $\partial\bar{\partial}$ closed, directed by lamination.

Berdtsen - S.

$$\int T \wedge T \quad \int T_g T$$

Main pt

$$\int T \wedge T = 0 \quad \text{if } T \geq 0 \\ T \text{ direct. by } \tilde{F}. \\ i\partial\bar{\partial}T = 0.$$

Hodge-Riemann. But
cohomology space Hilbert-space

what happens in arbitrary dim?

Birkhoff Thm. for foliations?

Geomet. Diffusion with heat equation requires bounded geometry of leaves.
we have sing.

Din. Yang-Mills.

- Laminations by R.S. transversally G^2 in a complex manifold M .
- Holomorphic foliations with finite number of sing in M .
- Abstract. lamination. ex: suspensions.
- Poincaré metric WP. measurable.

$T \geq 0$ $\partial\bar{\partial}T = 0$ directed by Foliation.

Thm. F foliation with finite number of linearizable sing.

Then:

$$\omega_p \sim \frac{1}{\|z\|^2 (\log \|z\|)^2} \quad \text{near sing (z=0)}$$

Then.

$$\int T \wedge \omega_p < \infty$$

If

$$\partial \bar{T} = \tau \wedge T$$

$$\underline{\underline{\mu_T = i \tau \wedge \bar{\tau} \wedge T}} \quad \int \mu_T < \infty$$

$$\partial T = \int \frac{\partial h_\alpha}{h_\alpha} \wedge h_\alpha [V_\alpha] d\mu(\alpha)$$

Th. (X, L, E) compact Lowinat.

Sing. E . is a complex manifold

M.

T extremal positive current

direct by L . $\partial\bar{\partial}$ -closed.

without mass or Parabolic
leaves.

$$m_p = T \wedge \omega_p \quad \phi_a(0) = a$$

Define $m_{a,R} = \frac{1}{M_R} (\phi_a)_* \left(\log^+ \frac{a}{|z|} \omega_p \right)$

$(M_R \sim R)$ Then

$$m_{a,R} \rightarrow m_p \quad \text{m.p. a.e.}$$

We can use another metric

$\omega_p \rightsquigarrow \text{metric. } \omega$

$$\frac{1}{c_n} (\phi_a)_* \left(\log^+ \frac{a}{|z|} \right) \rightarrow \frac{T}{11\pi n} \quad \text{T.a.e.}$$

Log

$$\int_{\mathbb{R}^D} |\phi'_a(z)|^2 \log^+ \frac{a}{|z|} \, dz \wedge d\bar{z} \simeq \log \frac{1}{1-\bar{z}}$$

$a \rightarrow 1$.

$$B_R u(a) = \frac{1}{M_R} \int \log \frac{r}{|z|} (\Phi_a)_* (u \omega_P)$$

$$= \langle m_{a,R}, u \rangle$$

Lemma

$u \in L^1(m_P)$ implies

$$\int (B_R u) T \wedge \omega_P = \int u T \wedge \omega_P.$$

Fubini!

L^P ergodicity.

$$(\Delta_P v) \omega_P := i \partial \bar{\partial} v$$

$$u = \Delta_P v$$

$$B_R u(a) = B_R (\Delta_P v)(a) = \langle m_{a,R}, (\Delta_P v) \omega_P \rangle$$

$$= \langle i \partial \bar{\partial} T_{a,R}, v \rangle \longrightarrow 0.$$

Pointwise ergodic thm.

$$\widetilde{B}u(a) := \overline{\lim_{R \rightarrow \infty}} |B_R u(a)|.$$

Lemma If $u \in L^1(\mu_P)$.

$$\mu_P \{ \widetilde{B}u > \varepsilon \} \leq C \varepsilon^{-1} \|u\|_{L^1(\mu_P)}.$$

The comparison.

$$\left| B_R u(a) - \frac{2\eta}{M_R} \int_0^{\frac{M_R}{2\eta}} S_t u(a) dt \right| \leq C R^{-\frac{1}{2}} \sqrt{\log R}$$

Heat equation for $T \geq 0$ (1.1) $i\partial\bar{\partial}T = 0$
 directed by a lamination (singular).

consider a set of bounded functions

1. bounded. 2. continuous. 3. with $\partial\bar{\partial}T = 0$.

$$T = \int_{\mathbb{H}} [V_{\alpha}] d\mu_{\alpha}$$

$$\partial T = \tau \wedge T$$

choose a (1,1) form $\beta \geq 0$.
 (ex $\omega_P, i\tau \wedge \bar{\tau}$)

$$\int T \wedge \beta < \infty.$$

Dirichlet Space w. respect to T .

$$\|u\|^2 = \int |u|^2 T \wedge \beta + \int i\partial u \wedge \bar{\partial} u \wedge T.$$

on smooth functions and complete.

Def: $(\Delta_P u) T \wedge \beta := i\partial\bar{\partial} u \wedge T.$

$$\tilde{\Delta}_P := \Delta_P + \frac{1}{2} \nabla_P$$

$$(\nabla_P u) T \wedge \beta = i(\partial u \wedge \tau - \bar{\partial} u \wedge \bar{\tau}) \wedge T.$$

$$m_P := T \wedge \beta$$

harmonic measure

2)

Th: $-\tilde{\Delta}_p$ ($\Delta_p - \Delta_p$) are max.
monotone on $L^2(u_p)$.

Then (Hille Yosida) they generate
a semi-group of contractions
on $L^2(u_p)$. $S(t)$

u_p is $S(t)$ invariant.

contract. in $L^p(u_p)$ i.e.
 $i\partial\bar{\partial}u^2 = 2iu\partial\bar{\partial}u + 2i\partial u\bar{\partial}u$.

Solutions are only in L^2 .

$$\begin{cases} \frac{\partial u(x,t)}{\partial t} = \Delta_p u(x,t) & \text{leaves} \\ u(x,0) = u_0. \end{cases}$$

Pr. $S(t)u \geq 0$ when $u \geq 0$

(Stampachia! trick).

Corollary

$$\frac{1}{R} \int_0^R S(t)u_0 dt \rightarrow u_0^*$$

i.e

con in $L^p(u_p)$.

Th If T extenal (mixing) $-\tilde{\Delta}_p$

$$S(t)u_0 \rightarrow \langle \infty, T \wedge \omega_p \rangle$$

2) u_0 a.e. and in $L^2(u_p)$.

Poincaré Metric

Entropy.

- $\mathbb{C}^{2+s}$ lamination in complex manifold M .
- Foliation with linearizable singularities.

Leaves covered by unit disc.

Entropy.

Quantify how leaves get apart.

Ex. general smooth lamination
without sing.

Ghys - Langevin - Wolke.

Existence of entropy

$\approx$ Lipschitz exp.

Entropy.

(X, d) metric space totally

bounded. $Y \subset X$

$D = \{d_t\} \nearrow$ metrics. $\left[d_n^{\text{inf}} = \sup_{j \leq n} d(f^j(x), f^j(y)) \right]$

$N(Y, t, \varepsilon)$ = Min number of

balls $B_\varepsilon(x, \varepsilon)$ necessary

to cover Y .

$$h_Y^{\text{inf}} = \overline{\lim} \frac{1}{t} \log N(Y, t, \varepsilon).$$

Ex 1 $M \xrightarrow{f} M$ macro-micro

M comp. Kühlen.

$$X_i = M / f^i(I)$$

$$X \xrightarrow{f} X$$

f has finite entropy (T.C. Dickson)

Calculs en entropies.

(X, L, E) hyperbolic.

$$\tau_a: D \rightarrow L$$

$$\tau_a(0) = a.$$

$$J(a, b) := \inf_{\gamma \in R} \sup_{x \in D_R} d(\tau_x(\dot{\gamma}^0), \tau_x(\dot{\gamma}^1)).$$

R- hyperbolic time.

(R, ε) close

Leaves through a, b time R .

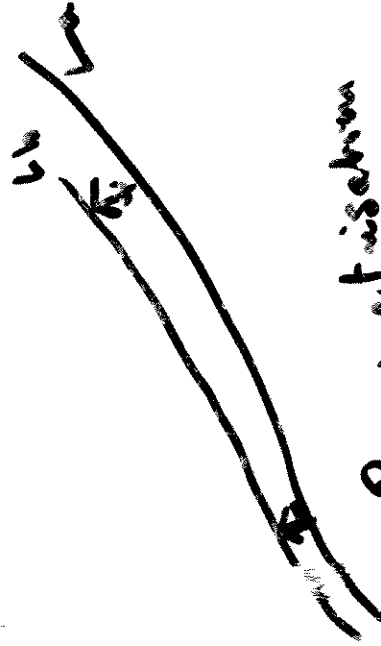
ε -close $\rightarrow$

$\mathcal{C}^{2+\delta}$ laminat. (without sing)

Thm

Poincaré metric is Hölder cont.
entropy is finite.

$$d_R \leq A^R d_0 + C e^{-R}.$$



Parametrisation

connection Laplace Beltrami

estimates.

M compact complex.
 F foliation linearizable sing.

Precise modulus of continuity

for $\eta = \frac{\|v\|_h}{\|v\|_p}$. Metric.

$$\sim \frac{1}{(\log \text{dist}(a,b))^{1+\epsilon}}$$

'Then' F foliat. Lin sing.

Entropy is finite.

Difficulty (R, ϵ) close near sing.

$$\approx \epsilon e^{-e^R}$$

but situation in homogeneous

dim $M = 2$. OK Fol. transitivity

conform. holonomy conform.

(X, L, E) Lamination
by R.S in a complex manifold
M. E finite set.

$$\Delta \xrightarrow{\phi} L$$

$$\tau_n^L := \frac{1}{c_n} \int_0^{\infty} \frac{d}{ds} \phi_*(\Delta_s) = \frac{1}{c_n} \phi_* \left(\log^+ \frac{2}{|s|} \right)$$

For example:

clusters points are positive
directed $\partial\bar{\partial}$ -closed currents

(1,1) τ_n is like a Birkhoff
sum on the "orbit"

Th (Favre-S) F foliation in $\mathbb{P}^2$
without algebraic leaves and
only hyperbolic singularities.
Then there is a unique $T_{\geq 0}$ $\partial\bar{\partial}$ -closed
directed current of mass 1.

In particular

$\tau_n^L \rightarrow T$ independently of

Dim 2. Sing. Intersection. Very

9) $\partial\bar{\partial}$ -closed currents.