

THE DIRAC OPERATOR FOR HIGGS BUNDLES

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*Blaine Lawson's 70th birthday
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HIGGS BUNDLES

- Riemann surface Σ
- Vector bundle E + hermitian metric
- Higgs field $\Phi \in H^0(\Sigma, \text{End } E \otimes K)$
- equations $F_A + [\Phi, \Phi^*] = 0$

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TWO- AND THREE-DIMENSIONAL INSTANTONS

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The four-dimensional Yang-Mills Lagrangian implies corresponding structures in lower dimensions. Instantons, characterized by a zero energy-momentum tensor as well as finite action, emerge as the solutions of coupled first order equations. For the Abelian case all such solutions are determined by the non-linear Poisson-Boltzmann equation.

term have acquired special values. One can obtain first order equations which imply the field equations, and these are deduced from eqs. (6):

$$F_{ij}^a = \pm e \epsilon_{ij} \epsilon^{abc} \psi^b \phi^c ,$$
$$D_i \psi^a \pm \epsilon_{ij} D_j \phi^a = 0 . \quad (11)$$

Unfortunately, in this case the model is not interesting because the action is always zero.

LOCALLY...

- anti-self-dual connection on principal bundle over \mathbf{R}^4
- invariant under $(x_1, x_2, x_3, x_4) \mapsto (x_1, x_2, x_3 + t_1, x_4 + t_2)$

$$\nabla_3 = \frac{\partial}{\partial x_3} + \phi_1(x_1, x_2) \quad \nabla_4 = \frac{\partial}{\partial x_4} + \phi_2(x_1, x_2)$$

$$[\nabla_1, \nabla_2] + [\phi_1, \phi_2] = 0 = [\nabla_1, \phi_2] + [\nabla_2, \phi_1] = 0 = [\nabla_1, \phi_1] + [\phi_2, \nabla_2]$$

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$$F + [\Phi, \Phi^*] = 0$$

$$\bar{\partial}_A \Phi = 0$$

CHALLENGE

Solve these equations like the ADHM construction of instantons

SPECIAL CASES

- $\Phi = 0$:

find the flat unitary connection on a stable vector bundle

- $E = K^{1/2} \oplus K^{-1/2}$, $\Phi(u, v) = (0, u)$:

find the hyperbolic metric on a Riemann surface

INSTANTONS

- ASD connection on E over S^4
- Dirac operator $D : C^\infty(S_+ \otimes E) \rightarrow C^\infty(S_- \otimes E)$
- $D^*D = \nabla^*\nabla + \frac{1}{4}R \Rightarrow \ker D = 0$
- $\dim \ker D^* = k$

- $(x_1, x_2, x_3, x_4) \in \mathbf{R}^4$, G -connection ∇
- $D = -\nabla_1 + i\nabla_2 + j\nabla_3 + k\nabla_4$ quaternions

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- $D = -\nabla_1 + \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \nabla_2 + \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \nabla_3 + \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \nabla_4$

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- D^*D real iff $[\nabla_1, \nabla_2] + [\nabla_3, \nabla_4] = 0$ etc.
- anti-self-dual Yang-Mills

- $D^*D = -\nabla_1^2 - \nabla_2^2 - \nabla_3^2 - \nabla_4^2$
- \mathcal{L}^2 solutions ψ_1, \dots, ψ_k of $D^*\psi = 0$
- matrices T_1, \dots, T_4

$$\langle T_i \psi_\alpha, \psi_\beta \rangle = \int_{\mathbf{R}^4} (x_i \psi_\alpha, \psi_\beta)$$

- \sim ADHM construction

DIRAC OPERATOR FOR HIGGS BUNDLES

- dimensional reduction

- ASD connection $A_1 dx_1 + A_2 dx_2 + \phi_1 dx_3 + \phi_2 dx_4$

- $$D^* = \nabla_1 + \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} \nabla_2 + \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \phi_1 + \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} \phi_2$$

- Dirac operator
$$D^* = \begin{pmatrix} \bar{\partial}_A & \Phi \\ \Phi^* & \partial_A \end{pmatrix}$$

globally...

- Σ compact, genus g

- $$D^* = \begin{pmatrix} \bar{\partial}_A & \Phi \\ \Phi^* & \partial_A \end{pmatrix} : \begin{pmatrix} E \otimes K \\ E \otimes \bar{K} \end{pmatrix} \rightarrow \begin{pmatrix} E \otimes K \bar{K} \\ E \otimes K \bar{K} \end{pmatrix}$$

Globally...

- Σ compact, genus g

- $$D^* = \begin{pmatrix} \bar{\partial}_A & \Phi \\ \Phi^* & \partial_A \end{pmatrix} : \begin{pmatrix} E \otimes K \\ E \otimes \bar{K} \end{pmatrix} \rightarrow \begin{pmatrix} E \otimes K \bar{K} \\ E \otimes K \bar{K} \end{pmatrix}$$

- $D^*D = -\nabla_1^2 - \nabla_2^2 - \phi_1^2 - \phi_2^2 \Rightarrow \ker D = 0$

- index theorem $\Rightarrow \dim \ker D^* = (2g - 2) \operatorname{rk} E$

- $(\psi_1, \psi_2) \in \mathbf{C}^\infty(EK \oplus E\bar{K})$
- conformally invariant inner product

$$\int_{\Sigma} \langle \psi_1, \psi_1 \rangle + \langle \psi_2, \psi_2 \rangle$$

- \Rightarrow inner product on $\ker D^*$

EXAMPLE

- E trivial line bundle on \mathbb{C} ; $\Phi = dz$
- $D^*(\psi_1, \psi_2) = 0$
- $\bar{\partial}\psi_1 + \psi_2 = 0, \psi_1 + \partial\psi_2 = 0$

$$\frac{\partial^2 \psi_1}{\partial z \partial \bar{z}} + \psi_1 = 0$$

- $\psi_2 = -\bar{\partial}\psi_1$

FOURIER TRANSFORM

M.Jardim, *Nahm transform for doubly-periodic instantons*, Commun.Math.Phys.. **225** (2002) 639–668

J.Bonsdorff *A Fourier transform for Higgs bundles*, Crelle **591** (2006) 21–48

J.Bonsdorff *Autodual Connection in the Fourier Transform of a Higgs Bundle*, Asian J. Math. **14** (2010), 153–174

S.Szabo, “Nahm transform for integrable connections on the Riemann sphere”, Société Mathématique de France (2007)

- Higgs bundle (E, Φ)
- L flat line bundle, θ holomorphic 1-form
- family $(E \otimes L, \Phi + \theta 1)$ parameters $(L, \theta) \in \mathcal{J}(\Sigma) \times H^0(\Sigma, K)$
- tensoring with a $U(1)$ Higgs bundle

- $\ker D_{L,\theta}^*$ defines a vector bundle over $T^*J = J \times H^0(\Sigma, K)$
- Higgs bundle equations \Rightarrow hermitian metric
- $\Rightarrow \mathcal{L}^2$ connection

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- Higgs bundle equations \Rightarrow hermitian metric
- $\Rightarrow \mathcal{L}^2$ connection
- the connection is hyperholomorphic

- Jacobian J flat Kähler torus
- T^*J flat hyperkähler manifold
- complex structures I, J, K
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hyperholomorphic = curvature type $(1, 1)$ wrt I, J, K

- ADHM: $D = A + qB$, D^*D real

- ADHM: $D = A + qB$, D^*D real
- $a, \theta \in H^0(\Sigma, K)$
- $$\begin{pmatrix} \bar{\partial}_A + \bar{a} & \Phi + \theta \\ \Phi^* + \bar{\theta} & \partial_A - a \end{pmatrix} = D^* + q$$
- $q \in H^0(\Sigma, K) \oplus \bar{H}^0(\Sigma, K)$ quaternionic vector space

- complex structure I

$$\begin{array}{ccc} \Omega^{0,p}(E) & \xrightarrow{\Phi} & \Omega^{0,p}(EK) \\ \downarrow \bar{\partial} & & \downarrow \bar{\partial} \\ \Omega^{0,p+1}(E) & \xrightarrow{\Phi} & \Omega^{0,p+1}(EK) \end{array}$$

- total differential $\bar{\partial} \pm \Phi$
- Hodge theory: $\ker D^* \cong$ hypercohomology \mathbf{H}^1

- complex structure J

- $\nabla = \partial_A + \bar{\partial}_A + \Phi + \Phi^*$ flat $GL(n, \mathbf{C})$ connection

$$\Omega^0(E) \xrightarrow{\nabla} \Omega^1(E) \xrightarrow{d\nabla} \Omega^2(E)$$

- Hodge theory: $\ker D^* \cong$ de Rham cohomology $H^1(\mathcal{E})$

FOURIER TRANSFORM

- complex structure $I: (L, \theta)$
- hypercohomology \mathbf{H}^1 of

$$\mathcal{O}(EL) \xrightarrow{\Phi + \theta} \mathcal{O}(ELK)$$

- complex structure $J: \mathcal{L} \in H^1(\Sigma, \mathbf{C}^*)$
- $H^1(\Sigma, \mathcal{E}\mathcal{L})$

- hyperholomorphic \Rightarrow holomorphic bundle on twistor space of $J \times H^0(\Sigma, K)$

- $[\Theta] \in H^1(J, T^*)$: extension $T^* \rightarrow V \rightarrow \mathcal{O}$
- $P(V) \setminus P(T^*)$ affine bundle over J
- $H^1(\Sigma, \mathbf{C}^*) \rightarrow H^1(\Sigma, \mathcal{O}^*) = J$
- $\zeta \in \mathbf{C}, \zeta[\Theta] \Rightarrow$ bundle (V_ζ, ζ) over $J \times \mathbf{C}$
- $\zeta = 0$ T^*J : product $J \times H^0(\Sigma, K)$

- $P(V) \setminus P(T^*)$ affine bundle $\cong H^1(\Sigma, \mathbb{C}^*)$

$$H^1(\Sigma, \mathbb{C}^*) \rightarrow J(\Sigma)$$

flat connection ∇

holomorphic structure $\nabla^{0,1}$

- $P(V) \setminus P(T^*)$ affine bundle $\cong H^1(\Sigma, \mathbb{C}^*)$

$$H^1(\Sigma, \mathbb{C}^*) \rightarrow J(\Sigma)$$

flat connection ∇ holomorphic structure $\nabla^{0,1}$

$$H^1(\Sigma, \mathbb{C}^*) \rightarrow J(\bar{\Sigma})$$

flat connection ∇ holomorphic structure on $\bar{\Sigma}$ $\nabla^{1,0}$

- twistor space: identify over $\mathbb{C}^* \subset \mathbb{C}$

CONCLUSION

- Higgs bundle (E, Φ)
- Fourier transform = rank $\text{rk } E(2g - 2)$ bundle \hat{E} on $\mathcal{J} \times H^0(\Sigma, K)$ with hyperholomorphic connection
- \sim holomorphic $\text{rk } E(2g - 2)$ bundle on twistor space Z

QUESTIONS

1. Does this vector bundle encode the Higgs bundle and its metric?

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2. If so, what is the inverse transform?
3. Find the transform of the canonical Higgs bundle.

COMPLEX STRUCTURE *I*

HYPERCOHOMOLOGY

- $\mathcal{O}(E) \xrightarrow{\Phi} \mathcal{O}(EK)$

two spectral sequences:

- cohomology of $\Phi : H^p(\Sigma, E) \rightarrow H^p(\Sigma, EK)$:

$$0 \rightarrow H^0(E) \rightarrow H^0(EK) \rightarrow \mathbf{H}^1(E) \rightarrow H^1(E) \rightarrow H^1(EK) \rightarrow 0$$

- cohomology of $\mathcal{O}(U, E) \xrightarrow{\Phi} \mathcal{O}(U, EK)$

$$\mathbf{H}^1(E) \cong H^0(\Sigma, \text{coker } \Phi)$$

- hypercohomology $\Phi + \theta 1 : \mathcal{O}(E \otimes L) \rightarrow \mathcal{O}(E \otimes LK)$
- $x_0\Phi + x_1\theta 1 : \mathcal{O}(E \otimes L) \rightarrow \mathcal{O}(E \otimes LK)$
- bundle extends to $J \times P(H^0(\Sigma, K) \oplus 1) = J \times P^g$

J.Bonsdorff *A Fourier transform for Higgs bundles*, Crelle **591**
(2006) 21–48

- $\text{ch}(\hat{E}) = \text{rk } E(g - 1 + (g - 1)e^x + \theta(1 - e^x))$

$$x \in H^2(\mathbb{P}^g, \mathbf{Z}), \theta \in H^2(\mathbf{J}, \mathbf{Z})$$

- $\dim H^p(\mathbf{J} \times \mathbb{P}^g, \hat{E}) = \text{rk } E \binom{g-1}{p-1}$

$1 \leq p \leq g$ zero otherwise

EXAMPLE

- $\Phi = 0$: stable bundle E over Σ , $\deg E = 0$
- $H^0(\Sigma, E) = 0$, $\dim H^1(\Sigma, E) = \text{rk } E(g - 1)$
- $\mathcal{P} \rightarrow \Sigma \times J$ Poincaré bundle
- \Rightarrow Picard bundle $R^1(E) \sim H^1(\Sigma, EL)$ over J
- $0 \rightarrow H^0(\Sigma, ELK) \rightarrow \mathbf{H}^1(\Sigma, EL) \rightarrow H^1(\Sigma, EL) \rightarrow 0$

- $p : \mathbf{J} \times \mathbf{P}^g \rightarrow \mathbf{J}$

- $0 \rightarrow p^*R^0(E) \rightarrow \hat{E} \rightarrow p^*R^1(E)(1) \rightarrow 0$

- $\Rightarrow \hat{E} \cong p^*R^0(E) \oplus p^*R^1(E)(1)$

- hypercohomology $\Phi + \theta 1 : \mathcal{O}(E \otimes L) \rightarrow \mathcal{O}(E \otimes LK)$
- $x_0\Phi + x_1\theta 1 : \mathcal{O}(E \otimes L) \rightarrow \mathcal{O}(E \otimes LK)$
- bundle extends to $J \times P(H^0(\Sigma, K) \oplus 1) = J \times P^g$

- Fourier-Mukai: equivalence of derived categories
- recover (E, Φ) as a holomorphic object

THE HYPERHOLOMORPHIC CONNECTION

- connection on a bundle over $J \times \mathbb{C}^g$
- What is its asymptotic behaviour?
- Why does the holomorphic structure extend to $J \times \mathbb{P}^g$?

M.Jardim, *Construction of doubly-periodic instantons*, CMP **216** (2001) 1–15

M.Jardim, *Nahm transform for doubly-periodic instantons*, CMP **225** (2002) 639–668

- instanton on $T^2 \times \mathbf{R}^2$,

$$\int_{T^2 \times \mathbf{R}^2} |F|^2 < \infty$$

- \Rightarrow Higgs bundle on \widehat{T}^2 with two singularities

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- instanton on $T^2 \times \mathbf{R}^2$,

$$\int_{T^2 \times \mathbf{R}^2} |F|^2 < \infty$$

- \Rightarrow Higgs bundle on $\widehat{T^2}$ with two singularities
- I -holomorphic structure extends from $T^2 \times \mathbf{C}$ to $T^2 \times \mathbf{P}^1$

- Dirac operator D^*

$$D^* = \begin{pmatrix} \bar{\partial}_A & \Phi \\ \Phi^* & \partial_A \end{pmatrix} + t \begin{pmatrix} 0 & \theta \\ \bar{\theta} & 0 \end{pmatrix}$$

- What is the asymptotic behaviour as $t \rightarrow \infty$ of solutions?

- Weitzenböck formula

$$DD^* = \nabla^* \nabla + A + tB + t^2|\theta|^2$$

- $D^*\psi = 0 \quad \int_{\Sigma} |\psi|^2 = 1$

- $\int_{\Sigma} \langle A\psi, \psi \rangle + t \int_{\Sigma} \langle B\psi, \psi \rangle + t^2 \int_{\Sigma} |\theta|^2 |\psi|^2 \leq 0$

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- $\int_{\Sigma} \langle A\psi, \psi \rangle + t \int_{\Sigma} \langle B\psi, \psi \rangle + t^2 \int_{\Sigma} |\theta|^2 |\psi|^2 \leq 0$

- $t^2 \int_{\Sigma} |\theta|^2 |\psi|^2 \leq tM$

- $U \subset \Sigma \quad \int_U |\theta|^2 |\psi|^2 \leq t^{-1} M$

- choose U disjoint from the zeros of θ

- $m \int_U |\psi|^2 \leq \int_U |\theta|^2 |\psi|^2 \leq t^{-1} M$

- $\int_U |\psi|^2 \rightarrow 0$

- solutions localize around the $2g - 2$ zeros of $\theta \in H^0(\Sigma, K)$
- $\mathcal{O}(E) \xrightarrow{\Phi} \mathcal{O}(EK)$
- hypercohomology $\mathbf{H}^1 \cong \text{cokernel of } \Phi \dots$
- supported on zeros of $\det \Phi \in H^0(\Sigma, K^n)$

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- supported on zeros of $\det \Phi \in H^0(\Sigma, K^n)$
- $\det(\Phi + t\theta) = t^n \theta^n + \dots + \det \Phi$
- zeros of $\theta \sim$ basis for \hat{E} on $\mathbb{P}(T^*)$

SPECIAL CASES

- $\Phi = 0$:

find the flat unitary connection on a stable vector bundle

- $E = K^{1/2} \oplus K^{-1/2}, \Phi(u, v) = (0, u)$:

find the hyperbolic metric on a Riemann surface

- S^1 -action $(A, \Phi) \mapsto (A, e^{i\theta}\Phi)$
- S^1 -action on T^*J
- $\Phi = 0$ fixed $\Rightarrow \hat{E}$ invariant

- $E = K^{1/2} \oplus K^{-1/2}, \Phi(u, v) = (0, u)$

- $(e^{i\theta/2}u, e^{-i\theta/2}v) : \Phi \mapsto e^{i\theta}\Phi$

- $\Rightarrow \hat{E}$ invariant

THM. (Feix, Kaledin). Let M be a real analytic Kähler manifold, then there is a unique S^1 -invariant hyperkähler metric on a neighbourhood of the zero section extending the Kähler metric.

B. Feix, *Hyperkähler metrics on cotangent bundles*, *J. Reine Angew. Math.* **532** (2001), 33-46.

THM. (Feix). Let A be a connection in a vector bundle over M whose curvature is of type $(1, 1)$, then this extends to a unique S^1 -invariant hyperholomorphic connection on a neighbourhood of the zero section.

B. Feix, *Hypercomplex manifolds and hyperholomorphic bundles*, *Math. Proc. Camb. Phil. Soc.* **133** (2002), 443-457

CANONICAL HIGGS BUNDLE

- $u_1 \in H^0(\Sigma, K^{3/2}L)$ $v_2 \in H^1(\Sigma, K^{-1/2}L)$
- $\deg K^{3/2}L = 3g - 3$
- Picard bundle V , $P(V) \cong S^{3g-3}\Sigma$
- \mathcal{L}^2 metric wrt hyperbolic metric.

