

On positivity of a class of conformal covariant operators

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- ▶ On (M^d, g) , k an integer, $(P_{2k}^d)_g$ is a class of differential operator of order $2k$ with leading symbol $(-\Delta_g)^k$; with the conformal covariant property that under conformal change of metric $g_w = e^{2w}g$, we have

$$(P_{2k}^d)_{g_w}(\phi) = e^{-\frac{d+2}{2}w}(P_{2k}^d)_g(e^{\frac{d-2}{2}w}\phi)$$

for all smooth functions ϕ defined on M .

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- ▶ When $k = 1$, P_2^d is the conformal Laplace or Yamabe operator:

$$(P_2^d)_g = -\Delta_g + \frac{d-2}{4(d-1)}R_g$$

where R_g is the scalar curvature of the metric g .

- ▶ When $d > 2k$, $(P_{2k}^d)_g(1) := \frac{d-2}{2}(Q_{2k}^d)_g$, e.g. when $k = 1$, $d > 2$, $Q_2^d = \frac{1}{2(d-1)}R$ while when $k = 1$, $d = 2$, Q_2^2 is defined to be the Gaussian curvature.

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- ▶ When $d = 2k$, Branson's curvature $Q_{2k} = Q_{2k}^{2k}$ is also defined. When (M^d, g) is locally conformally flat, $C_d \chi(M) = \int (Q_d)_g dv_g$; in general, $\int (Q_d)_g dv_g$ is a conformal invariant.

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- ▶ On the flat case, \mathbb{R}_+^{n+1} , $g_+ = \frac{dy^2 + dx^2}{y^2}$, where $y > 0, x \in \mathbb{R}^n$, for all $\gamma > 0$, we have

$$P_{2\gamma} = (-\Delta_x)^\gamma, Q_{2\gamma} = 0.$$

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- ▶ Note when $\gamma = 1$, R (scalar curvature) > 0 implies P_2 (the conformal Laplace operators) > 0 .
- ▶ **Theorem** (Gursky ('99)) On (M^4, g) closed manifolds, $Y(M, g) > 0$ and $\int (Q_4)_g dv_g > 0$ implies that $(P_4)_g > 0$. In particular, when $R_g > 0$, $(Q_4)_g > 0$ then $(P_4)_g > 0$.

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- ▶ Formulas of P_4^d , Q_4^d by Paneitz ('83) and GJMS,

$$(P_4^d)_g = (-\Delta_g)^2 + \text{div}(a_d \text{Ric}_g + b_d R^g)D + \frac{(d-4)}{2} Q_{4,g}^d$$

$$Q_{4,g}^d = c_1(-\Delta_g R + c_2 R_g^2 - c_3 |\text{traceless Ric}_g|^2)$$

where c_1, c_2, c_3 are positive dimensional constants.

- ▶ Gursky-Malchiodi ('13) ($d \geq 4$),
On (M^d, g) closed manifolds, if $R_g > 0$, and $(Q_4^d)_g > 0$, then $(P_4^d)_g > 0$ when $d \geq 4$.
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- ▶ **Hang-Yang** ('14) ($d = 3$) When $R_g > 0$ and $(Q_4^3)_g > 0$, Green's function of $(P_4^3)_g$ is negative.

Main Theorems

- ▶ Positivity of $P_{2\gamma}$ as boundary operators of conformal compact Einstein manifolds On (X^{n+1}, M^n, g_+) Poincare Einstein manifold, where $M^n = \partial X$, the conformal infinity boundary. We (J. Case - Chang) have two results:

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- ▶ Theorem 1: When $0 < \gamma < 1$, $d \geq 2$, and $\Lambda_1(-\Delta_{g_+}) > \frac{n^2}{4} - \gamma^2$, $Q_{2\gamma}^n > 0$ implies $P_{2\gamma}^n > 0$.

Remark: By a work of J. Lee, if $R_{(M^n, g_0)} > 0$, where $g_0 = g|_M$, then $\Lambda_1(-\Delta_{g_+}) \geq \frac{n^2}{4}$.

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- ▶ Theorem 2: When $1 < \gamma < 2$, $n \geq 4$, $R_{(M, g_0)} > 0$ and $Q_{2\gamma}^n > 0$ implies $P_{2\gamma}^n > 0$. When $n = 3$, the same result holds when $1 < \gamma \leq \frac{3}{2}$.

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- ▶ Key step in the proof: The “right” choice of the conformal compactified Einstein metric in X^{n+1} .

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- ▶ Classical Setting: (work of Caffarelli-Silvestre)
- ▶ Well-known result: f smooth on \mathbb{R}^n

$$\Delta_{x,y} U(x,y) \text{ on } \mathbb{R}^{n+1}_+ = (x,y | x \in \mathbb{R}^n, y > 0,$$

$$U|_{\mathbb{R}^n}(x) = f(x)$$

then

$$-U_y(x,0) = (-\Delta)^{\frac{1}{2}} f(x).$$

- ▶ Theorem (Caffarelli-Silvestre '06)

$$0 < \gamma < 1, a = 1 - 2\gamma,$$

$$(*) \begin{cases} \operatorname{div}(y^a \nabla U) & = 0 \text{ on } \mathbb{R}_+^{n+1} \\ U|_{\mathbb{R}^n} & = f. \end{cases}$$

Then

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$$0 < \gamma < 1$$

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- ▶ Applications to free-boundary problems, study of non-local minimal surface etc.

- ▶ On Conformal Compact Einstein Setting, a class of conformal covariant operators $P_{2\gamma}$ exists for

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- ▶ $P_{2\gamma} = (-\Delta)^\gamma$ in special setting of $(\mathbb{R}_+^{n+1}, \mathbb{R}^n, g_+)$, where

$$g_+ = \frac{dy^2 + dx^2}{y^2}$$

while

$$\bar{g} = y^2 g_+.$$

is the compactified metric.

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- ▶ Thus

$$g_+ = \frac{dy^2 + g_y}{y^2} \text{ in } M \times (0, \epsilon).$$

Denote $\bar{g} = y^2 g_+$, then $\bar{g} = dy^2 + g_y$ on X , and $g_y|_{y=0} = g_0$ on M .

- On (X^{n+1}, M^n, g_+) , given $f \in C^\infty(M)$

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 $F = f + f_2 y^2 + f_4 y^4 \cdots$, where $f_{2i} \in C^\infty(M)$
- ▶ u is the solution of the Poisson equation with Dirichlet data f .

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When $s = \frac{n}{2} + \gamma$, $\gamma \notin \mathbb{Z}^+$,

define $P_{2\gamma} = \mathcal{S}(\frac{n}{2} + \gamma)$ is a non-local pseudo-differential operator with leading symbol $|\xi|^{2\gamma}$.

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- ▶ P_{2k} is the **GJMS operators**.

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- ▶ $(H^{n+1}/\Gamma, \Omega(\Gamma)/\Gamma, g_h)$, where Γ a Kleinian group.

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$$g_h = r^{-2}(|dr|^2 + (1 - \frac{r^2}{4})^2 g_{S^n}).$$

- ▶ $(H^{n+1}/\Gamma, \Omega(\Gamma)/\Gamma, g_h)$, where Γ a Kleinian group.
- ▶ Schwarzschild space.

Theorem (C- Gonzalez '11)

On (X^{n+1}, M^n, g_+) conformal compact Einstein setting, given a function $f \in C^\infty(M)$;

$$(*)_s \quad -\Delta_{g_+} u - s(n-s)u = 0 \text{ on } X$$



$$s = \frac{n}{2} + \gamma$$

$$(*)'_s \quad -\operatorname{div}_{\bar{g}}(\rho^a \nabla_{\bar{g}} U) + E(\rho, a)U = 0 \text{ on } X$$

$$U = \rho^{s-n} u \quad U| = f$$

where $\bar{g} = \rho^2 g_+$, $0 < \gamma \leq \frac{n}{2}$, ρ : any totally geodesic defining function.

We can express $P_{2\gamma} f$ in terms of boundary behavior of U .

- ▶ In general, the expression of $E(\rho, a)$ is complicated, but in the special case when $\gamma = \frac{1}{2}$, $a = 0$, $E(\rho, a) = C_n R_{\bar{g}}$, the equation becomes

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Suppose

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and v is Poisson operator on data $f \equiv 1$. **Note** if $v > 0$ on X^{n+1} , one can define $\rho^* = v^{\frac{1}{n-s}}$ then $E(\rho^*, a) = 0$, where $s = \frac{n}{2} + \gamma$ and $a = 1 - 2\gamma$.

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- ▶ It turns out when $0 < \gamma < 1$, $v > 0$ if and only if $\lambda_1(-\Delta_+) > \frac{n^2}{4} - \gamma^2$.

Proof of Theorem 1

Lemma: When $0 < \gamma < 1$, ($a = 1 - 2\gamma$) and

$$\rho^* = y + d_\gamma Q_{2\gamma} y^{1+2\gamma} + O(y^3) > 0.$$

Given f , $u =$ solution of Poisson equation with data f ,
 $U := (\rho^*)^{\gamma - \frac{n}{2}} u$. Then for some positive c_γ , we have

$$P_{2\gamma} f(x) = c_\gamma \lim_{y \rightarrow 0} (\rho^*)^a \frac{\partial U}{\partial n}(x) + \frac{n - 2\gamma}{2} Q_{2\gamma} f(x),$$

for $x \in \partial X$.

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- ▶ Proof:

With the choice of ρ^* , $g^* = (\rho^*)^2 g_+$, we have

$U := (\rho^*)^{\gamma - \frac{n}{2}} u$ and $U|_M = f$, U satisfies the PDE

$$-\operatorname{div}_{g^*}((\rho^*)^a \nabla_{g^*} U) = 0.$$

Hence

$$\int_X (-\operatorname{div}_{g^*}((\rho^*)^a \nabla_{g^*} U) U) = 0$$

and

$$\int_M (\rho^*)^a \frac{\partial U}{\partial n} U = \int_X (\rho^*)^a |\nabla U|^2 \geq 0.$$

We then apply the Lemma to finish the proof.

Outline of the Proof of Theorem 2

- ▶ Theorem 2: When $1 < \gamma < 2$, $d \geq 4$, $R_{(M^d, g_0)} > 0$ and $Q_{2\gamma}^d > 0$ implies $P_{2\gamma}^d > 0$. When $d = 3$, the same result holds when $1 < \gamma \leq \frac{3}{2}$.

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- ▶ Step 4: Apply the extension theorem and some proof similar to that of Gursky-Malchiodi to establish the theorem.

Extension Theorem in flat case when $\gamma > 1$

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Extension Theorem in flat case when $\gamma > 1$

- ▶ Recent work of R. Yang, here for the special case when $1 < \gamma < 2$.
- ▶ On $\mathbb{R}_+^{n+1} = \{(x, y) : x \in \mathbb{R}^n, y > 0\}$, Denote

$$\Delta_a U = y^{-a} \operatorname{div}(y^a \nabla U) = \Delta U + \frac{a}{y} \frac{\partial U}{\partial y}$$

Then

$$\Delta_a U = 0, \text{ with } U|_{\mathbb{R}^n} = f$$

where $a = 1 - 2\gamma$, iff

$$(\Delta_b)^2 U = 0, \text{ with } U|_{\mathbb{R}^n} = f, \text{ and } \lim_{y \rightarrow 0} y^b \frac{\partial U}{\partial y} = 0$$

with $b = 3 - 2\gamma$.

Extension Theorem in flat case when $\gamma > 1$

In this case



$$\int_{\mathbb{R}^n} (-\Delta_x f)^\gamma f dx = c_{n,\gamma} \int_{\mathbb{R}_+^{n+1}} (\Delta_b U)^2 y^b dx dy.$$

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$$(-\Delta_x f)^\gamma(x) = c_{n,\gamma} \lim_{y \rightarrow 0} y^b \frac{\partial}{\partial y} \Delta_b U(x, y).$$

We have the “renormalized energy”, e.g. when $\gamma = \frac{3}{2}$,
 $a = 1 - 2\gamma = -2$, $b = 3 - 2\gamma = 0$, then

$$\lim_{\epsilon \rightarrow 0} \left(- \int_{\mathbb{R}^n} \int_{y \geq \epsilon} |\nabla U|^2 y^{-2} dx dy + \frac{1}{\epsilon} \int_{y=\epsilon} |\nabla_x f|^2 dx \right) = \frac{1}{2} \int_{\mathbb{R}^n} \int_{y \geq 0}$$

Extension Theorem on Poincare Einstein setting

- ▶ Notation: Given a number $m \in \mathbb{R}$, ϕ a function defined on (X, g) , (F, h) a metric space of dimension m ; on the metric measure space $(X, g, e^{-\phi} dv_g)$, denote $P_{2k, \phi}^m$ the GJMS operators on the warped product space

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- ▶ In this notion, when $m = \infty$, $Ric_\phi^m = Ric + \nabla^2 \phi$ the Bakry-Emery Ricci tensor, Δ operator is replaced by $\Delta_\phi := \Delta - \nabla \phi \nabla$.

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- ▶ (1) On $(X^{n+1}, \partial X, g_+)$, C.C.E. with $Ric_{g_+} = -n$, When $s = \frac{n}{2} + \gamma$, $g = \rho^2 g_+$, the $(*)_s$ equation

$$-\Delta_{g_+} u - s(n - s)u = 0, \text{ on } X$$

can be re-written as $(*)''$

$$P_{2,\phi}^m U = 0 \text{ on } X$$

where (F^m, h) is chosen to be the (sphere) with $Ric_h = (m - 1)h$, $U = \rho^{s-n}u$ and $g = \rho^2 g_+$, $m = 1 - 2\gamma$ and $e^{-\phi} = \rho^m$.

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▶ With this notion, we have the extension theorem on Poincare Einstein manifolds.

Step 3, the right choice of compactified metric

- ▶ Recall if $v = v_s$ satisfies the Poisson equation with Dirichlet data $f \equiv 1$, under the condition $R_{\partial X, g_0} > 0$, we have $v > 0$ on X . Denote $\rho^* = v^{\frac{1}{n-s}}$, and $g = g^* = (\rho^*)^2 g_+$, $s = \frac{n}{2} + \gamma$.

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 - ▶ (b) When $1 < \gamma < 2$, $(Q_4)_{\phi_2}^{m_2} = 0$, where $m_2 = 1 - 2\gamma$, $e^{-\phi_2} = (\rho^*)^{m_2}$

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 - ▶ (c) When $\gamma > 1$, $R_{g^*}|_{\partial X} = c_{\gamma} R(\partial X, g_0) > 0$, where $c_{\gamma} > 0$.

Step 3, the right choice of compactified metric

- ▶ Another crucial property:

Lemma: Under the assumption $R_{\partial X, g_0} > 0$, for all $s \geq \frac{n}{2} + 1$, $R_{g^*} > 0$ on X .

Proof: Due to property (b) above, we have the PDE for $R = R(g^*)$,

$$\Delta_{\phi_2} R = c_1 R^2 - c_2 |E|^2,$$

where $c_1 = c_1(s)$, $c_2 = c_2(s)$ are positive constants, and E the traceless Ricci.

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- ▶ It turns out when $s = n + 1$, $c_1 = 0$, the metric has been studied before by J. Lee '95, where the equation by maximal principle together with property (c) gives $R_{g^*} > 0$.
- ▶ We can now run a “continuity” argument on the parameter s starting at $s = (n + 1)$, together property (c), apply strong maximal principle to conclude $R_{g^*} > 0$ on X for all $s \geq \frac{n}{2} + 1$. 

Proof of Theorem 2

When $1 < \gamma < 2$, we will show that when $R_{(\partial X, g_0)} > 0$ and $Q_{2\gamma} > 0$, implies $P_{2\gamma} > 0$.

Proof:

Given f defined on ∂X , by Extension theorem

$$\int_{\partial X} (P_{2\gamma} f) f dv_{g_0} = \frac{n-2\gamma}{2} \int_{\partial X} (Q_{2\gamma} f) f dv_{g_0} + c_\gamma \text{ Energy term of } (P_4)_{\phi_2}^{m_2}.$$

We apply the fact $R_{g^*} > 0$, together with an argument similar to that of [Grusky-Malchiodi](#) to prove the 4-th order energy term is non-negative, and which together with $Q_{2\gamma} > 0$ establishes the result.

Some discussion

► Theorem

(*J. Qing - Guillarmou '10*) On (X^{n+1}, M^n, g^+) C.C.E. manifolds with $n + 1 > 3$, $Y(M^n, g_0) > 0$ iff the first real scattering pole $\leq \frac{n}{2} - 1$.

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- Equivalent Statement: Under the assumption $Y(M, g_0) > 0$, $P_{2\gamma} \geq 0$ for all $0 < \gamma < 1$.

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► The result generalizes an earlier work of Schoen-Yau.

$X = H^{n+1}/\Gamma,$	Γ a Kleinian group
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► *Schoen-Yau*: If M is of positive scalar curvature, then $\delta(\Gamma) \doteq$ Hausdroff dim of $S^n \setminus \Omega(\Gamma)$, then $\delta(\Gamma) \leq \frac{n}{2} - 1$.

► Work of *Sullivan – Patterson, P. Perry* etc.

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- ▶ Work of [Gonzalez-Qing '12](#) studied the $Q_{2\gamma}$ equation and related positive mass problem when $0 < \gamma < 1$. When $\gamma = \frac{1}{2}$, $Q_1 = cH$, the mean curvature. In general, is there a geometric description of the fractional Q curvature?