MAT 303 Lecture 19

Central motion (continued)

Theorem 2 (conservation of angular momentum)
In polar coordinates we have $\vec{F} = (F_r, F_\theta)$
$F_\theta = 0$ for a central force
$\vec{a} = m\vec{r}'' = m\left((\vec{r}')^2 + r\vec{\theta}''\right)$
$\Rightarrow 0 = 2r\vec{r}'\vec{\theta}' + r^2\vec{\theta}'' = (r^2\vec{\theta}')'$
$\Rightarrow r^2\vec{\theta}' = h$ (conservation of angular momentum per unit mass)

Theorem 3 (equal areas are swept out in equal times):
$A = \int_{\theta_0}^{\theta_1} \frac{1}{2} r^2 \, d\theta$
$\frac{dA}{dt} = \frac{1}{2} r^2 \frac{d\theta}{dt}$
$\Rightarrow A' = \frac{1}{2} r^2 \theta' = \frac{1}{2} h$
$\Rightarrow A = \frac{1}{2} h t$ (area increases linearly)

Force fields

If you have a force at every point then a $\vec{F}(x)$

Consider $F_1 dx + F_2 dy + F_3 dz$

We can ask if this is $\vec{\nabla} f$ for some scalar function f; this would mean

$F_1 = \frac{\partial f}{\partial x}, \quad F_2 = \frac{\partial f}{\partial y}, \quad F_3 = \frac{\partial f}{\partial z}$

If such a U exists, then $-U$ is the potential of the field. F_1 the field is the known coordinates

Example 1: $F = -kx$ (spring)

$F_x = \frac{\partial}{\partial x}(\frac{1}{2}kx^2) = -kx \qquad F_y = \frac{\partial}{\partial y}(\frac{1}{2}kx^2) = 0$
$F_z = \frac{\partial}{\partial z}(\frac{1}{2}kx^2) = 0$

$U(x,y,z) = \frac{1}{2}k(x^2+y^2+z^2) = \frac{1}{2}kr^2$

Example 2: $F = \frac{k}{r^2}$ $\quad r = \sqrt{x^2+y^2+z^2}$

$F_x = \frac{x}{r}(\frac{k}{r^2}) = \frac{kx}{r^3} \qquad F_y = \frac{ky}{r^3} \qquad F_z = \frac{kz}{r^3}$

$U(x,y,z) = \frac{k}{\sqrt{x^2+y^2+z^2}} = \frac{k}{r}$

Theorem 4: In a conservative field, energy is conserved

$dU = F_1 dx + F_2 dy + F_3 dz$

$\Rightarrow \frac{\partial U}{\partial x} = mx'', \quad \frac{\partial U}{\partial y} = my', \quad \frac{\partial U}{\partial z} = mz''$

$mx''\frac{dx}{dt} + my'\frac{dy}{dt} + mz''\frac{dz}{dt} = \frac{dU}{dt}$

$\Rightarrow \frac{1}{2}mv_x^2 + \frac{1}{2}mv_y^2 + \frac{1}{2}mv_z^2 = U + E$

$\Rightarrow \frac{1}{2}m|\vec{v}|^2 - U = E$ (conserved) ∎

Motion when $F = -k^2 \vec{u}$ (inverse square law)

$F_x = mx'' = F\frac{x}{r} = -k^2 x \frac{1}{r}$

$F_y = my'' = -k^2 y \frac{1}{r}$

$\begin{cases} x'' = -k^2 x \\ y'' = -k^2 y \end{cases} \qquad \lambda^2 + k^2 = (\lambda + ki)(\lambda - ki)$

$x = A\cos(kt) + B\sin(kt)$
$y = C\cos(kt) + D\sin(kt)$

$\begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} \cos(kt) \\ \sin(kt) \end{bmatrix} = \begin{bmatrix} x \\ y \end{bmatrix}$

$\Rightarrow \begin{bmatrix} \cos(kt) \\ \sin(kt) \end{bmatrix} = \frac{1}{AD-BC} \begin{bmatrix} D & -B \\ -C & A \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix}$

$= \frac{1}{AD-BC} \begin{bmatrix} Dx - By \\ -Cx + Ay \end{bmatrix}$

$1 = \cos^2(kt) + \sin^2(kt)$
$= \frac{(Dx-By)^2 + (-Cx+Ay)^2}{(AD-BC)^2}$

$\Rightarrow (AD-BC)^2 =$
$(D^2+C^2)x^2 - 2(BD+AC)xy + (B^2+A^2)y^2$
$dx^2 + Bxy + Cy^2 = \delta$ (equation of ellipse)

Example: $\vec{F} = -\frac{k^2}{r^2}\vec{u}$ $\quad r = \sqrt{x^2+y^2}$ $\quad \vec{u} = \frac{1}{r}\begin{bmatrix} x \\ y \end{bmatrix}$

$x = A\cos(\frac{k}{r}) + B\sin(\frac{k}{r})$
$y = C\cos(\frac{k}{r}) + D\sin(\frac{k}{r})$
$x^2 + \frac{AC+BD}{AD-BC}xy + \frac{B^2+A^2}{AD-BC}y^2 = \frac{1}{AD-BC}$
$y^2 = \frac{1}{C^2+D^2}(\frac{1}{r} + D\sin(\frac{k}{r}))$

$\Rightarrow B = A \qquad$ 2) $D = C$

$\Rightarrow D - \frac{1}{C} = 0 \qquad$ 4) $C = \frac{1}{D}$

$x = 1\cos(\frac{k}{r}) + \frac{1}{D}\sin(\frac{k}{r}) \Rightarrow \left|x - \frac{1}{2}\right|^2 + \left(\frac{1}{2}\right)^2$
$y = D\sin(\frac{k}{r})$

Motion when $\vec{F} = -\frac{k^2}{r^2}\vec{u}$ (input is the force)

$F_x = \frac{x}{r}(\frac{k^2}{r^2}) = \frac{xk^2}{(x^2+y^2)^{3/2}} = mx'' = \frac{-k^2x}{(x^2+y^2)^{3/2}}$

We do not yet know anything!

Do this in polar coordinates!

$F_r = -\frac{k^2}{r^2} = m(r'' - r(\theta')^2)$
$F_\theta = my' = m(2r'\theta' + r\theta'')$ ← r is constant, so $r'' = 0$

$F_\theta = 0$ (because central) $\Rightarrow r^2\theta' = h \Rightarrow \theta' = \frac{h}{r^2}$

$F_r = \frac{-k^2}{r^2} = m(r'' - r(\theta')^2) = m(r'' - \frac{h^2}{r^3})$

$\Rightarrow r'' - \frac{h^2}{r^3} = -\frac{k^2}{r^2}$

Goal: solve for r as a function of θ

$u = \frac{1}{r} \qquad \theta' = hu^2 \qquad r' = \frac{1}{u^2}$

Physicist proof:

$r'' = \frac{-2}{u^3}\frac{du}{d\theta}\theta' = \frac{-2}{u^3}\frac{du}{d\theta}hu^2 = -h\frac{du}{d\theta}$

$r'' - h\frac{du}{d\theta}\frac{d\theta}{dt} = -h^2u^3 \frac{du}{d\theta}$

Transform:

Denote $u(\theta) = B(\theta)$

$r'' = \frac{-2}{u^3}\frac{du}{d\theta}(B(\theta))\frac{d\theta}{dt} = \frac{-2}{u^3}\frac{dB}{d\theta}\frac{d\theta}{dt}$

$r'' - h\frac{dB}{d\theta}(\theta)\frac{d\theta}{dt} = -h^2u^3\frac{dB}{d\theta}(\theta)$

Recall:

$r'' - \frac{h^2}{r^3} = \frac{-k^2}{r^2} \qquad u = \frac{1}{r}$

$-h^2u^3\frac{dB}{d\theta} - h^2u^3 = -Ku^2$

$\frac{d^2B}{d\theta^2}(\theta) + B(\theta) = \frac{K}{h^2}$

Method: to solve $B''(\theta)$ with an affine B

$B'' + B = \frac{K}{h^2}$ (and note how $\frac{d^2}{d\theta^2}$ acts)

Homogeneous eqn: $\lambda^2 + 1 = 0 \Rightarrow \lambda = \pm i$

$B = A\cos(\theta - \delta)$ calls to brainstorm

$B = \frac{K}{h^2}$ solves the inhomogeneous problem (by inspection)

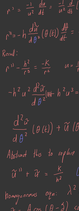
$B = A\cos(\theta - \delta) + \frac{K}{h^2}$

Let $A = \frac{K}{h^2}e$, then

$B(\theta) = \frac{K}{h^2}(1 + e\cos(\theta - \delta)) \Rightarrow r(\theta) = \frac{k^2}{K}\frac{1}{1+e\cos(\theta-\delta)}$

Choose coordinates such that $\delta = 0$

Conic sections are:



$e = \frac{c}{a} \qquad e = \frac{k^2}{Kd} \Rightarrow d = \frac{1}{e}\frac{k^2}{K}$

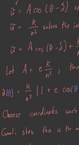
$d_1 = r\cos\theta + d \Rightarrow \frac{kr^2}{ke} = \frac{r}{e} + r\cos\theta$

$\Rightarrow \frac{kr}{K} = r(1 + e\cos(\theta))$

$\Rightarrow r = \frac{k}{K}\frac{1}{1+e\cos(\theta)}$

$e < 1 \Rightarrow$ ellipse ← depends on $e \geq 0$ exists
$e > 1 \Rightarrow$ parabola
$e = 1 \Rightarrow$ hyperbola

Cosmology Problem



λ = the linear density of a segment of light ray

Problem: To the x axis, the bottom of Earth $= -\lambda\Delta s$

To the y axis, you'd have $T_1\cos\theta_1 - T_2\cos\theta_2 = 0$

To the y axis, you'd have $T_1\sin\theta_1 - T_2\sin\theta_2 - \lambda\Delta s = 0$

$\Rightarrow \Delta(T\cos\theta) = 0$

$\Delta(T\sin\theta) = \lambda\Delta s$

$\Rightarrow \frac{d}{ds}(T\sin\theta) = 0 \qquad \Rightarrow T\sin\theta = H$

$\frac{d}{ds}(T\sin\theta) = \lambda\frac{ds}{ds}$

$\frac{d}{ds}\left(\frac{T\sin\theta}{T\cos\theta}\right) = \frac{1}{H}\frac{dy}{ds} = k\frac{dy}{ds} \qquad k = \frac{\lambda}{H}$

$\ln\theta = \frac{dy}{ds} \qquad ds = \sqrt{dx^2+dy^2}$

$\Rightarrow \frac{dy}{dx} = k\sqrt{1+(\frac{dy}{dx})^2}$

$\frac{d}{dx}y' = k\sqrt{1+(y')^2}$

Let $w = y'$

$\Rightarrow \begin{cases} \frac{dw}{dx} = k\sqrt{1+w^2} \\ \frac{dy}{dx} = w \end{cases} \Rightarrow k\,dx = \frac{dw}{\sqrt{1+w^2}}$