

10/03/2019

Gliders and stuff.

Talk at the board. Too bad for you cuz you didn't come to class. Probably you should at least read the [start of the text](#) about it.

Some yammering:

As you know, sometimes we can do integrals and get a nice answer, like

$$\begin{aligned} > \int x^2 dx \\ & \qquad \qquad \qquad \frac{x^3}{3} \end{aligned} \qquad (1)$$

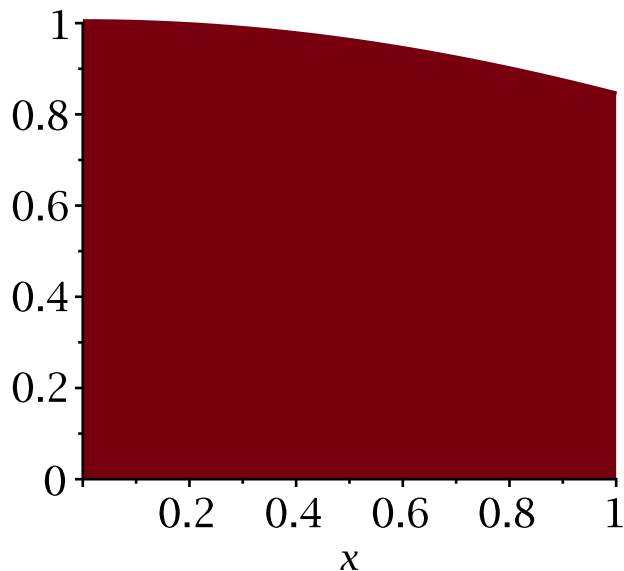
and sometimes the answers don't have an answer in terms of elementary functions, such as

$$\begin{aligned} > \int \frac{\sin(x)}{x} dx \\ & \qquad \qquad \qquad \text{Si}(x) \end{aligned} \qquad (2)$$

Above, [Si\(x\)](#) represents the "sine integral", which just means the function that represents $\int_0^x \frac{\sin(t)}{t} dt$, that is, the integral is just its own self.

Of course, we can UNDERSTAND this integral, for example, by thinking of it as an area under a curve.

> `plot($\frac{\sin(x)}{x}$, x = 0..1, filled = true, thickness = 3)`



Similarly, a differential equation might have a nice closed form solution, or not. Note that differential equations are just ways to solve integrals... that is, solving the differential equation

$$\frac{d}{dx} (g(x)) = x^2$$

is the same thing as saying "find the function of x whose derivative is x^2 ", that is, find $\int x^2 dx$

Of course, maple can do this just fine. Note that we can write $g'(x)$ as $\text{diff}(g(x), x)$

> $\text{dsolve}(\text{diff}(g(x), x) = x^2)$

$$g(x) = \frac{x^3}{3} + _C1 \quad (3)$$

If we have an initial condition, eg,

$$\left\{ \frac{d}{dx}(g(x)) = x^2, g(0) = 7 \right\}$$

we can have maple solve that too.

> $\text{dsolve}(\{ \text{diff}(g(x), x) = x^2, g(0) = 7 \})$

$$g(x) = \frac{x^3}{3} + 7 \quad (4)$$

Even if there are a few coupled equations, we might get lucky and there will be a nice solution:

> $\text{dsolve}(\{ \text{diff}(x(t), t) + y(t) = t \cdot x(t), \text{diff}(y(t), t) + (y(t))^3 = \sin(t) \})$

$$\left[\left\{ y(t) = \sin(t)^{1/3}, y(t) = -\frac{\sin(t)^{1/3}}{2} - \frac{\sqrt{3} \sin(t)^{1/3}}{2}, y(t) = -\frac{\sin(t)^{1/3}}{2} + \frac{\sqrt{3} \sin(t)^{1/3}}{2} \right\}, \left\{ x(t) = \frac{y(t)}{t} \right\} \right] \quad (5)$$

But this won't always work.

Now let's go back to talking about gliders. See the [start of the text](#) to see where these equations come from. But they describe how the velocity and angle of a glider affected by gravity, lift, and drag behave.

R is a "coefficient of friction".

> $\text{phug} := \left[\text{diff}(v(t), t) = -\sin(\text{theta}(t)) - R \cdot v(t)^2, \text{diff}(\text{theta}(t), t) = \frac{v(t)^2 - \cos(\text{theta}(t))}{v(t)} \right]$

$$\text{phug} := \left[\frac{d}{dt} v(t) = -\sin(\theta(t)) - R v(t)^2, \frac{d}{dt} \theta(t) = \frac{v(t)^2 - \cos(\theta(t))}{v(t)} \right] \quad (6)$$

We can ask for a solution, but I'm not happy with the "answer".

> $\text{dsolve}(\text{phug})$

$$\left[\left\{ \theta(t) = _a \text{ where } \left[\left(\frac{d}{d_a} _b(_a) \right) _b(_a) \right] \right\} \right] \quad (7)$$

$$\begin{aligned}
& + \frac{1}{2(-b(a) + \sqrt{b(a)^2 + 4 \cos(a)})} ((-b(a)^2 + 4 \cos(a))^{3/2} R \\
& + b(a)^2 \sqrt{b(a)^2 + 4 \cos(a)} R - 2 b(a)^3 R - 8 b(a) \cos(a) R \\
& + 4 \sin(a) \sqrt{b(a)^2 + 4 \cos(a)} + 4 b(a) \sin(a)) = 0 \}, \{ -a = \theta(t), \\
& b(a) = \frac{d}{dt} \theta(t) \}, \left\{ t = \int \frac{1}{-b(a)} da + C1, \theta(t) = -a \right\} \left. \right\}, \left\{ v(t) = \frac{\frac{d}{dt} \theta(t)}{2} \right. \\
& \left. - \frac{\sqrt{\left(\frac{d}{dt} \theta(t)\right)^2 + 4 \cos(\theta(t))}}{2}, v(t) = \frac{\frac{d}{dt} \theta(t)}{2} \right. \\
& \left. + \frac{\sqrt{\left(\frac{d}{dt} \theta(t)\right)^2 + 4 \cos(\theta(t))}}{2} \right\} \left. \right]
\end{aligned}$$

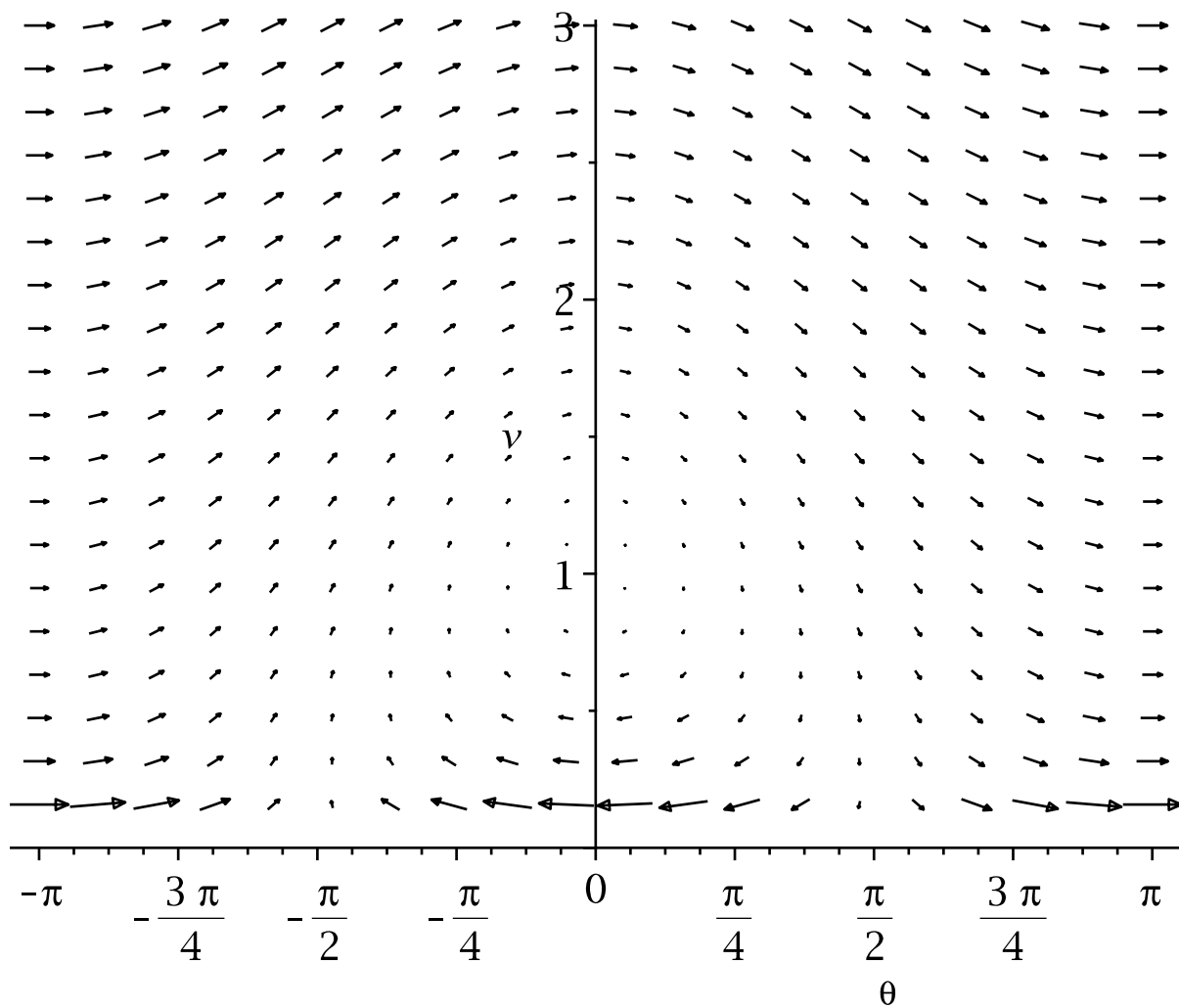
As I said, this is horrible, and doesn't actually tell us much since the solutions depend on derivatives, etc.... bleah!

Let's try something else. We can make a plot of the vector field, where at each point (θ, v) , we put an arrow to indicate the direction the solution is going, and its length tells us how fast it is changing there.

Let's just do that for $R=0$, and we can ignore t .

> *with(plots) :*

> *fieldplot* $\left(\left[v - \frac{\cos(\theta)}{v}, -\sin(\theta) \right], \theta = -\text{Pi}..\text{Pi}, v = 0..3, \right.$
arrows = slim, tickmarks = [piticks, default] $\left. \right)$



Let's write this as a differential equation, which is a pain to type all the time.

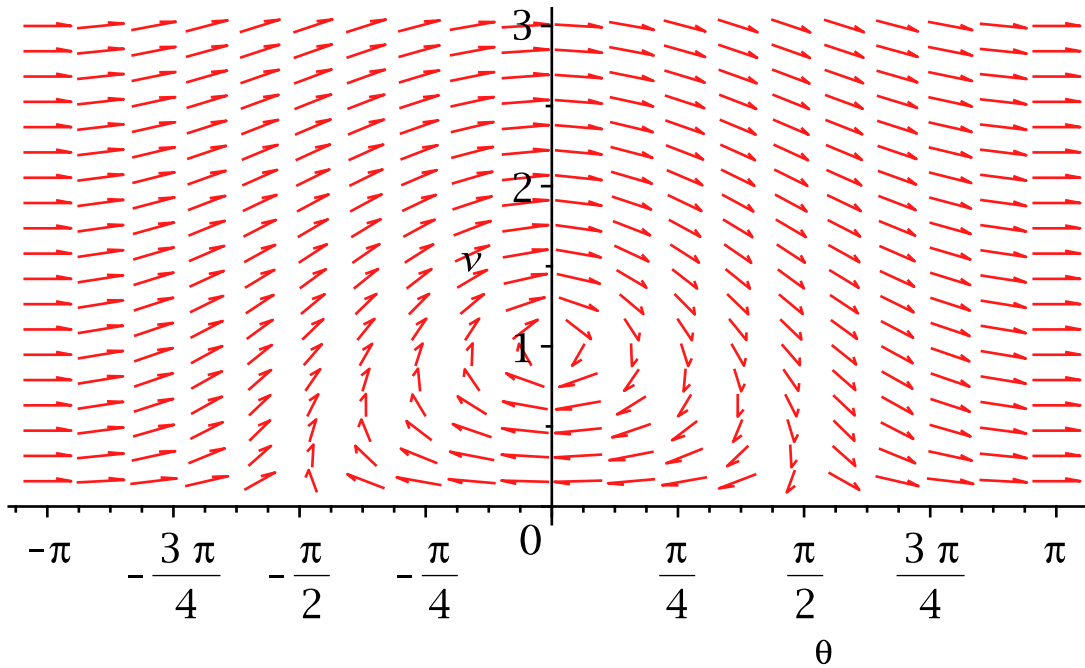
$$\begin{aligned}
 > \text{phug} := \left[\text{diff}(v(t), t) = -\sin(\text{theta}(t)), \text{diff}(\text{theta}(t), t) = \frac{v(t)^2 - \cos(\text{theta}(t))}{v(t)} \right] \\
 & \text{phug} := \left[\frac{d}{dt} v(t) = -\sin(\theta(t)), \frac{d}{dt} \theta(t) = \frac{v(t)^2 - \cos(\theta(t))}{v(t)} \right] \quad (8)
 \end{aligned}$$

Let's make a direction field-- a vector field where all the vectors have been rescaled to have length one. It is actually easier to see the solutions on the direction field.

```

> with(DEtools):
> DEplot(phug,
[theta(t), v(t)], t = 0..1, theta = -Pi..Pi, v = 0..3,
tickmarks = [piticks, default], scaling = constrained);

```



It is quite apparent, if you just look a little, that solutions near $\theta=0, v=1$ "circle around it".

This corresponds to a regular up and down motion of the plane (θ increases and decreases) and an oscillating velocity.

For solutions with a higher velocity, θ is monotonically increasing, while the velocity rises and falls a little (so the plane is looping).

Let's ask maple to plot two of these: one starting at $\theta=0, v=1/2$, and the other starting at $\theta=-\pi$ (that is, flying the other way, upside-down) and $v=2$.

```
> DEplot( phug,
```

```
  [theta(t), v(t)], t = 0..4,
```

```
  [ [theta(0) = 0, v(0) = 1/2], [theta(0) = -Pi, v(0) = 2] ],
```

```
  theta = -Pi..Pi, v = 0..3, tickmarks = [piticks, default], scaling = constrained,
```

```
  linecolor = blue );
```

