



Free Bosons and Tau-Functions for Compact Riemann Surfaces and Closed Smooth Jordan Curves. Current Correlation Functions

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Abstract. We study families of quantum field theories of free bosons on a compact Riemann surface of genus g . For the case $g > 0$, these theories are parameterized by holomorphic line bundles of degree $g - 1$, and for the case $g = 0$ —by smooth closed Jordan curves on the complex plane. In both cases we define a notion of τ -function as a partition function of the theory and evaluate it explicitly. For the case $g > 0$ the τ -function is an analytic torsion, and for the case $g = 0$, the regularized energy of a certain natural pseudo-measure on the interior domain of a closed curve. For these cases we rigorously prove the Ward identities for the current correlation functions and determine them explicitly. For the case $g > 0$, these functions coincide with those obtained by using bosonization. For the case $g = 0$, the τ -function we have defined coincides with the τ -function introduced as a dispersionless limit of the Sato's τ -function for the two-dimensional Toda hierarchy. As a corollary of the Ward identities, we note some recent results on relations between conformal maps of exterior domains and τ -functions. For this case, we also define a Hermitian metric on the space of all contours of given area. As another corollary of the Ward identities, we prove that the introduced metric is Kähler and the logarithm of the τ -function is its Kähler potential.

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Introduction

The general notion of a Riemann surface was defined by Riemann. He proved that, in modern terminology, every compact Riemann surface is a Riemann surface of an algebraic function $f(x, y) = 0$. Continuing the work of Abel and Jacobi, Riemann also introduced the notion of a general theta-function and solved the Jacobi inversion problem. After Riemann's death, his approach was repeatedly criticized

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by Weierstrass for the lack of rigor: in the proof of Dirichlet principle Riemann assumed that Dirichlet functional has a minimum. Later, Riemann's main results were proved by Schwarz and C. Neumann without using this assumption and in 1901 Hilbert finally proved the Dirichlet principle itself. At the same time Weierstrass criticism stimulated the development of the algebraic theory of algebraic functions, i.e. a theory that does not use complex-analytic methods. This theory was formulated by Brill, M. Noether, Dedekind and H. Weber, quite in parallel with the theory of algebraic numbers. In 1924, Artin was the first to consider algebraic functions over fields different from the field \mathbb{C} of complex numbers. Artin's results were generalized by F. K. Schmidt who, in particular, developed the theory of algebraic functions over a finite field. This development finally led A. Weil, among other things, to the general formulation of the algebraic geometry over an arbitrary field (see [20] for an exposition and references).

This old story nowadays is again in the spotlight due to the advent of the string theory. It is amazing that dramatic applications of quantum fields and strings to various mathematical areas are quite similar in spirit to Riemann's original work. Thus in his proof of the existence of a harmonic function with prescribed singularities at given points, Riemann used the analogy with electrostatic theory, assuming as obvious that for any charge distribution there exists an electrostatic potential, or equivalently, there exists a flow of the ideal fluid with prescribed sources and sinks (see, e.g., [19]). Using modern terminology, Riemann was exploiting methods of classical field theory for mathematical purposes, 'probing' mathematical objects with physical theories and translating the physical 'output' back into the mathematical statements. The same exact idea is in the heart of today's applications, with the 'only' difference that classical fields are being replaced by quantum fields and strings. Succinctly, this idea can be described as follows. In classical field theory one studies critical points of action functionals (like Dirichlet functional) which satisfy partial differential equations. In quantum field theory one studies partition and correlation functions of quantum fields defined by Feynman path integrals. In this formalism the critical values of the action functional as well as 'higher invariants' — so-called quantum corrections, naturally appear in the perturbative expansion (see [15] for exposition aimed at mathematicians). When probing mathematical objects (e.g. topological, smooth or complex-analytic manifolds) by quantum field theories, the mathematical output is encoded in partition and correlation functions that are expanded in terms of the critical values and quantum corrections.

Symmetries of physical theories play a fundamental role. Continuous symmetries in classical field theories correspond to conservation laws, and in quantum field theories they correspond to the so-called Ward identities expressing correlation functions through the partition function. The exploitation of symmetry is a powerful tool for studying conformal field theories in two dimensions, as discovered by A. B. Belavin, A. M. Polyakov and A. B. Zamolodchikov [4]. In addition to important physical applications to the critical phenomena in statistical mechanics, conformal field theories have a rich mathematical structure with applications to representation

theory, complex analysis, uniformization of Riemann surfaces and complex algebraic geometry of moduli spaces.

Quantum field theories of free fermions and free bosons on compact Riemann surfaces are basic examples of conformal field theories. It was proved by L. Alvarez-Gaumé, J.-B. Bost, G. Moore, P. Nelson, and C. Vafa and by E. Verlinde and H. Verlinde (see [1, 3, 6, 44] and references therein) that these theories are equivalent. Namely, there exists a remarkable correspondence between bosons and fermions, called bosonization, that expresses partition and correlation functions for one theory in terms of the other theory. Since for both theories these functions can be written down explicitly in terms of theta-functions and related algebro-geometric objects, bosonization yields non-trivial identities between them. In particular, it gives another proof of the celebrated Fay's trisecant identity for theta-functions [12], which play a fundamental role in the KP hierarchy (see, e.g., [28]). Approach in [3, 44] was based on path integrals with mathematical proofs using analytic torsion and Quillen's type isometries of determinant line bundles for $\bar{\partial}$ -operators. Another approach to bosonization based on Sato's infinite Grassmannian manifold was developed by L. Alvarez-Gaumé, C. Gomez and C. Reina [2], C. Vafa [43] and by N. Kawamoto, Y. Namikawa, A. Tsuchiya and Y. Yamada [21]. Purely algebro-geometric approach was given by A. Raina [35, 36, 37]. In addition to the proof of Fay's trisecant identity, in papers [21, 36, 37] bosonization was also used to compute explicitly bosonic current correlation functions from correlation functions for fermion operators.

This story brings us to Part 1 of the paper, where we derive bosonic current correlation functions directly from the basic $U(1)$ -symmetry of the theory. We define τ -function as a partition function and, following [3], show that it is essentially an absolute value of Sato's τ -function computed in [21]. Using explicit description of the complex structure of the Jacobian variety of a compact Riemann surface and basic properties of the Abel-Jacobi map we prove $U(1)$ -gauge symmetry Ward identities. As the result, we completely determine all reduced normalized multi-point current correlation functions in terms of partial derivatives of $\log \tau$ with respect to the complex coordinates on the Jacobian. Obtained formulas are in perfect agreement with results in [21, 36, 37]. Though our exposition is based on path integrals, all integrals that actually appear are Gaussian and have pure mathematical definition. Therefore our approach is completely rigorous.

It is well-known that Ward identities for other quantum field theories also encode important mathematical information. An interesting example is given by the quantum Liouville theory — the theory that arises as a conformal anomaly in Polyakov's approach to the string theory [32]. It was conjectured by Polyakov [33] that the semi-classical limit of conformal Ward identities for the Liouville theory yields an expression for the Poincaré's accessory parameters in terms of the critical value of Liouville action functional. Together with P. G. Zograf we have proved this and other results on the relation between accessory parameters and the Weil–Petersson geometry of the Teichmüller space [48–50]. In [41] we summarized the

geometric approach to two-dimensional quantum gravity and interpreted results in [48–50] as conformal Ward identities for multi-point correlation functions of stress-energy tensor components.

From this perspective, Part 1 of this paper is just another application of the same idea. The only difference is that instead of the quantum Liouville theory we consider simpler quantum theory of free bosons on a Riemann surface and instead of correlation functions with stress-energy tensor components we consider current correlation functions, which are also much simpler. Also, instead of the ‘nonlinear’ Teichmüller theory — the deformation theory of complex structures on a given compact Riemann surface used in [48–50], we are using a ‘linear theory’ of the Jacobian variety — the deformation theory of degree zero holomorphic line bundles over the Riemann surface.

In Part 2, we exploit the same simple idea for the case of free bosons on bounded simply-connected domains in the complex plane \mathbb{C} . Using results of Part 1 as a motivation, for every smooth closed Jordan curve C in \mathbb{C} we formulate a quantum field theory on the Riemann sphere \mathbb{P}^1 that depends on the interior domain Ω of a contour C . We define the τ -function of a contour C as a normalized partition function of the theory and compute it explicitly. It turns out to be the exponential of the regularized energy of a certain pseudo-measure on the domain Ω and, quite remarkably, coincides with the Mineev-Weistein-Wiegmann-Zabrodin τ -function [24, 29, 45], defined as a dispersionless limit of the Sato’s τ -function for the two-dimensional Toda hierarchy. We prove the Ward identities which express through $\log \tau$ the difference between current correlation functions of free bosons parameterized by C and current correlation functions of free bosons on the exterior domain $\mathbb{P}^1 \setminus \Omega$ satisfying the Dirichlet boundary condition. Correlation functions for free bosons on $\mathbb{P}^1 \setminus \Omega$ with the Dirichlet boundary condition can be written explicitly in terms of the conformal map G of the domain $\mathbb{P}^1 \setminus \Omega$ onto the exterior of the unit disk. Thus the Ward identities give another proof of remarkable relations between the conformal map G and the τ -function, discovered by M. Mineev-Weistein, P. B. Wiegmann and A. Zabrodin [29] in the content of the physical problem of 2D interface dynamics and studied extensively in [24, 45]. We also introduce a Hermitian metric H on the space of all contours of given area and prove as a corollary of the Ward identities that this metric is Kähler and the logarithm of the τ -function is a Kähler potential of the metric H .

The existence of the conformal map G is a consequence of the Riemann mapping theorem that every simply-connected domain in the complex plane \mathbb{C} whose boundary consists of more than two points can be conformally mapped onto the unit disk. There is a Riemann’s proof of this theorem which is based on the classical field theory and defines the conformal mapping through the complex potential of a certain charge distribution/fluid flow in the domain [19]. Corresponding quantum field theory naturally introduces the τ -function of contour C . Just as conformal Ward identities for the quantum Liouville theory imply that the critical value of the Liouville action is a generating function for the accessory parameters [48–50], cur-

rent Ward identities for the theory of free bosons on \mathbb{P}^1 imply that the τ -function of smooth bounded contours is a generating function for the conformal maps. The role of Teichmüller theory in this case is played by the deformation theory of contours as it has been stated by I. M. Krichever [26].

Thus, both parts of the paper are related not just by exploiting the same idea of Ward identities for current correlation functions, but also by using Riemann’s two major achievements in complex analysis as a framework for quantum field theories.

The above examples illustrate the importance of probing one-dimensional complex manifolds: compact Riemann surfaces and simply-connected domains in the complex plane, by quantum field theories. This idea can be pushed further by considering adelic formulation of quantum field theories on algebraic curves over an arbitrary field of constants [42, 46, 47] where Ward identities result in reciprocity laws. There is also a possibility of defining quantum field theories on the fields of algebraic numbers and it seems that the proper setting should be ‘quantum field theories on regular one-dimensional schemes’.

In this paper we are systematically using physical terminology. This is done for the only purpose: to show the true origin of the ideas and methods, and not for the purpose of proofs. Mathematically oriented reader can completely ignore these terms, which are *emphasized* in the main text on their first appearance, and consider them simply as words.* All results in the paper are rigorously proved (with usual space limitations) and no knowledge of quantum physics is required for what follows.

Here is a more detailed description of the content of the paper. Section 1 of Part 1 reviews some necessary mathematical facts: Jacobians, properties of the Abel–Jacobi map and theta-functions in 1.1, Green’s functions 1.2. In addition to standard definition and properties of the Green’s function of a Laplace operator on an compact Riemann surface X of genus g , in Lemma 1.1 we included a proof of the Fay’s formula relating two classical kernels on X . In Section 2 we define the theory of free $U(1)$ -bosons on X with classical fields** $g \in C^\infty(X, U(1))$ and field currents $\mathbf{J} := g^{-1}dg/2\pi\sqrt{-1}$, parameterized by the holomorphic line bundle L over X of degree $g - 1$ with $h^0(L) = 0$. We introduce the action functional $S(\mathbf{J})$ of the theory as a sum of the action functional for the standard free bosons on X and the topological term depending on L . We follow [3] with minor technical improvement in the invariant definition of the topological term. We define τ -function — the partition function as

$$\tau = \langle \mathbf{1} \rangle_L := \int_{C^\infty(X, U(1))/U(1)} [\mathcal{D}\mathbf{J}] e^{-2\pi S_L(\mathbf{J})}$$

and in Proposition 2.3 evaluate it explicitly in terms of the Riemann theta-function. The computation is a slight simplification of the original proof [3]. Namely, we apply

*‘Guaranteed to raise a smile’ from physicists.

**It should be always clear from the content whether g is a genus, $U(1)$ -valued bosonic field, or an inverse to a conformal map G as in Part 2.

the Poisson summation formula for the lattice \mathbb{Z}^{2g} instead of the lattice \mathbb{Z}^g as in [3] and use the transformation formula for the Riemann theta-function. In Section 3 we prove U(1)-gauged symmetry Ward identities for normalized reduced current correlation functions. Specifically, we introduce holomorphic and anti-holomorphic components of the bosonic field current as $J = -g^{-1}\partial g$ and $\bar{J} = -g^{-1}\bar{\partial}g$, and define multi-point current correlation functions as

$$\begin{aligned} & \langle J(P_1) \dots J(P_m) \bar{J}(Q_1) \dots \bar{J}(Q_n) \rangle \\ & := \int_{C^\infty(X, U(1))/U(1)} [\mathcal{D}\mathbf{J}] J(P_1) \dots J(P_m) \bar{J}(Q_1) \dots \bar{J}(Q_n) e^{-2\pi S_L(\mathbf{J})}. \end{aligned}$$

The U(1)-gauged Ward identities for normalized 1-point correlation function have the form

$$\frac{\partial \log \tau}{\partial z_i} = \int_{a_i} \langle \langle J(P) \rangle \rangle, \quad i = 1, \dots, g,$$

where z_i are complex coordinates of the point Z on the Jacobian $J(X)$ corresponding to the line bundle L , and directly follow from Riemann bilinear relations.

The Fay's formula is relevant for the computation of the reduced normalized 2-point correlation function. General expression for the multi-point correlation functions is given in Theorem 3.3 and is in perfect agreement with [21, 36, 37].

Basically, Part 1 serves as a motivation for Part 2. In Section 1 of Part 2 we recall necessary mathematical facts. In Section 1.1, following A. A. Kirillov [22], we review the infinite-dimensional space \mathcal{C} of all smooth closed Jordan curves on \mathbb{C} encircling the origin 0, and other spaces related to it. We also introduce a double \mathbb{P}_C^1 of the exterior domain $\mathbb{P}^1 \setminus \Omega$ of the contour C . In Section 1.2 we recall basic facts about classical Green's functions on \mathbb{P}^1 and on $\mathbb{P}^1 \setminus \Omega$, including Schiffer kernel and Bergman's reproducing kernel. Section 1.3 is devoted to detailed exposition of Krichever's deformation theory, outlined in [26]. We introduce the analogs of Faber polynomials for the conformal map G and define the harmonic moments of interior t_0, t_n of the contour C and harmonic moments of exterior v_0, v_n of C , $n \in \mathbb{N}$. We formulate and prove Theorem 1.5 which contains explicit description of holomorphic tangent and cotangent spaces to the infinite-dimensional complex manifold $\tilde{\mathcal{C}}_a$ of contours of fixed area $a > 0$. The main ingredient is the so-called Krichever's lemma [26] that expresses vector fields $\partial/\partial t_0, \partial/\partial t_n, \partial/\partial \bar{t}_n$ on \mathcal{C} as meromorphic $(1, 0)$ -forms on \mathbb{P}_C^1 . In addition to [26], we introduce natural Hermitian metric H on the spaces $\tilde{\mathcal{C}}_a$ using the Bergmann reproducing kernel. The metric H turns out to be Kähler (see Section 3 and description below).

In Section 2 we define theories of free bosons on \mathbb{P}^1 parameterized by smooth Jordan contours $C \in \mathcal{C}$. The action functional $S_C(\varphi)$ of the theory is defined as a sum of the action functional for the standard free bosons on \mathbb{P}^1 and of the analog

of the topological term depending on C . The partition function is defined as

$$\langle \mathbf{1} \rangle_C := \int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi] e^{-\frac{1}{\hbar} S_C(\varphi)},$$

and the τ -function of the contour C is given by

$$\tau := \frac{\langle \mathbf{1} \rangle_C}{\langle \mathbf{1} \rangle_\emptyset},$$

where \emptyset is the empty set that formally corresponds to the case when no contour is present. In Proposition 2.1, we compute the τ -function explicitly. It turns out to be the exponential of a regularized energy of the pseudo-measure given by the difference between the Lebesgue measure on the domain Ω and the delta-measure at 0 times the Euclidean area of Ω and coincides with the Wiegmann–Zabrodin τ -function [24, 45]. In Section 3 we introduce holomorphic and anti-holomorphic components $J = \partial\varphi, \bar{J} = \bar{\partial}\varphi$ of the field current $d\varphi$ and define multi-point current correlation functions

$$\begin{aligned} & \langle J(z_1) \dots J(z_m) \bar{J}(w_1) \dots \bar{J}(w_n) \rangle \\ & := \int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi] J(z_1) \dots J(z_m) \bar{J}(w_1) \dots \bar{J}(w_n) e^{-\frac{1}{\hbar} S_C(\varphi)} \end{aligned}$$

for free bosons on \mathbb{P}^1 parameterized by $C \in \mathcal{C}$. In a similar fashion we introduce current correlation functions for free bosons on $\mathbb{P}^1 \setminus \Omega$ with the Dirichlet boundary condition, denoted by $\langle \cdot \rangle_{\text{DBC}}$. In Theorem 3.1, we prove the Ward identity for the 1-point normalized correlation function

$$\frac{\partial \log \tau}{\partial t_n} = -\frac{1}{2\pi i} \int_C \langle \langle J(z) \rangle \rangle z^n dz, \quad n \in \mathbb{N}.$$

As a corollary we immediately get that the τ -function is a generating function for the harmonic moments of interior, as first proved in [24, 26, 45],

$$v_0 = \frac{\partial \log \tau}{\partial t_0}, \quad v_n = \frac{\partial \log \tau}{\partial t_n}, \quad n \in \mathbb{N}.$$

In Corollary 3.7 we derive another Wiegmann–Zabrodin result [45] — an ‘explicit formula’ for the conformal map G

$$\log G(z) = \log z - \frac{1}{2} \frac{\partial^2 \log \tau}{\partial t_0^2} - \sum_{n=1}^{\infty} \frac{z^{-n}}{n} \frac{\partial^2 \log \tau}{\partial t_0 \partial t_n}.$$

In Theorem 3.9, we prove the following Ward identities for normalized reduced

2-point correlation functions

$$\begin{aligned} \frac{\partial^2 \log \tau}{\partial t_m \partial t_n} &= \frac{1}{(2\pi i)^2} \int_C \int_C \left(\frac{G'(z)G'(w)}{(G(z) - G(w))^2} - \frac{1}{(z - w)^2} \right) z^m w^n dz dw \\ &= \frac{1}{(2\pi i)^2} \int_C \int_C z^m w^n (\langle \langle J(z)J(w) \rangle \rangle - \langle \langle J(z)J(w) \rangle \rangle_{\text{DBC}}) \end{aligned}$$

and

$$\begin{aligned} \frac{\partial^2 \log \tau}{\partial t_m \partial \bar{t}_n} &= -\frac{1}{(2\pi i)^2} \int_{C_+} \int_{C_+} \frac{G'(z)\bar{G}'(w)}{(1 - G(z)\bar{G}(w))^2} z^m \bar{w}^n dz d\bar{w} \\ &= -\frac{1}{(2\pi i)^2} \int_{C_+} \int_{C_+} z^m \bar{w}^n (\langle \langle J(z)\bar{J}(w) \rangle \rangle - \langle \langle J(z)\bar{J}(w) \rangle \rangle_{\text{DBC}}), \end{aligned}$$

where C_+ is an arbitrary contour around 0 containing the contour C inside.

The latter formula shows that the natural Hermitian metric H on infinite-dimensional complex manifolds $\tilde{\mathcal{C}}_a$ is Kähler and $\log \tau$ is its Kähler potential. This is a new result.

In numerous remarks throughout the paper we point out to interesting connections with other fields, classical complex analysis and theory of the univalent functions in particular. We plan to return to these questions, as well as to the discussion of the bosonization, in a separate paper.

PART 1. FREE BOSONS AND TAU-FUNCTIONS FOR COMPACT RIEMANN SURFACES

1. Mathematical Set-up

Here we recall, in a succinct form, necessary facts from complex (algebraic) geometry of compact Riemann surfaces (algebraic curves). The standard reference is [17]; see also [11, 30] as well as [3] for a ‘crash course’ for physicists.

Let X be a compact Riemann surface of genus g , which is always assumed to be connected and without the boundary. Denote by $K = K_X$ the canonical line bundle on X — holomorphic cotangent bundle of X , by $\text{Div}(X)$ — the group of divisors on X , and by $\text{Pic}(X)$ — the Picard group of isomorphism classes of holomorphic line bundles over X . The correspondence between line bundles and divisors provides canonical isomorphism $\text{Pic}(X) \simeq \text{Div}(X)/\text{PDiv}(X)$ — the group of divisors on X modulo the subgroup $\text{PDiv}(X)$ of principal divisors. Let $\check{H}^i(X, \mathcal{F})$ be the Čech cohomology groups with coefficients in a sheaf \mathcal{F} on X and set $h^i(\mathcal{F}) = \dim_{\mathbb{C}} \check{H}^i(X, \mathcal{F})$, $i = 0, 1$. As usual, we denote by L a holomorphic line bundle over X as well as the sheaf of germs of holomorphic sections of L . The Riemann–Roch

formula is

$$h^0(L) - h^0(K \otimes L^{-1}) = \deg L + 1 - g,$$

where $\deg L$ is a degree of the line bundle L .

1.1. JACOBIANS AND THETA-FUNCTIONS

The period map

$$\check{H}^0(X, K) \ni \omega \mapsto \int_c \omega \in \mathbb{C},$$

for all homology classes $[c] \in H_1(X, \mathbb{Z})$ of 1-cycles c on X , defines canonical inclusion of $H_1(X, \mathbb{Z})$ into $\check{H}^0(X, K)^\vee$ — the dual vector space to $\check{H}^0(X, K)$. The additive subgroup $H_1(X, \mathbb{Z})$ of $\check{H}^0(X, K)^\vee$ is a discrete subgroup of maximal rank $2g$ over \mathbb{R} . The Albanese variety $\text{Alb}(X)$ of a compact Riemann surface X is canonically defined as the following g -dimensional complex torus

$$\text{Alb}(X) := \check{H}^0(X, K)^\vee / H_1(X, \mathbb{Z}).$$

The Albanese variety $\text{Alb}(X)$ is a complex projective manifold. It carries translation-invariant Kähler metric, defined by the following Hermitian inner product at $\check{H}^0(X, K) = T_0^* \text{Alb}(X)$ — the holomorphic cotangent vector space to $\text{Alb}(X)$ at 0,

$$\langle \omega_1, \omega_2 \rangle := \frac{\sqrt{-1}}{2} \int_X \omega_1 \wedge \bar{\omega}_2, \quad \omega_1, \omega_2 \in \check{H}^0(X, K).$$

Let $\text{Pic}^0(X)$ be the group of degree 0 line bundles over X — the identity component of the Picard group $\text{Pic}(X)$. It follows from the standard exponential exact sequence and from the Dolbeault isomorphism

$$\text{Pic}^0(X) = \check{H}^1(X, \mathcal{O}) / \check{H}^1(X, \mathbb{Z}) = H_{\check{g}}^{0,1}(X) / H^1(X, \mathbb{Z}),$$

where \mathcal{O} is the structure sheaf on X and $H^1(X, \mathbb{Z}) := \text{Hom}(H_1(X, \mathbb{Z}), \mathbb{Z})$ is a lattice in the de Rham cohomology group $H^1(X, \mathbb{R})$ of cohomology classes of 1-forms with integral periods. By the Hodge theorem, $H^1(X, \mathbb{Z}) \simeq \mathcal{H}^1(X, \mathbb{Z})$ — the space of harmonic 1-forms on X with integral periods, and the mapping

$$\mathcal{H}^1(X, \mathbb{Z}) \ni \omega = \overline{\omega^{0,1}} + \omega^{0,1} \mapsto \omega^{0,1} \in H_{\check{g}}^{0,1}(X)$$

gives canonical embedding $H^1(X, \mathbb{Z}) \hookrightarrow H_{\check{g}}^{0,1}(X)$. The following canonical isomorphisms

$$\text{Div}^0(X) / \text{PDiv}(X) \simeq \text{Pic}^0(X) \simeq \text{Alb}(X)$$

follow from the Serre and Poincaré dualities.

A Riemann surface X is called Torelli marked if it is equipped with a symplectic basis $\{a_i, b_i\}_{i=1}^g$ for $H_1(X, \mathbb{Z})$. For a Torelli marked Riemann surface X , let $\{\omega_i\}_{i=1}^g$ be the basis for $\check{H}^0(X, K)$ of abelian differentials with normalized a -periods: $\int_{a_i} \omega_j = \delta_{ij}$, and let $\Omega = (\Omega_{ij})_{i,j=1}^g$ be the matrix of b -periods $\Omega_{ij} = \int_{b_i} \omega_j$. By Riemann bilinear relations, the matrix Ω is symmetric with positive definite imaginary part $\text{Im } \Omega = (\Omega - \bar{\Omega})/(2\sqrt{-1})$. Denote by $\Lambda = \mathbb{Z}^g \oplus \Omega\mathbb{Z}^g$ the period lattice of a Torelli marked Riemann surface X . It is a discrete subgroup of rank $2g$ in \mathbb{C}^g . The Jacobian variety of X is defined as the following g -dimensional complex torus $J(X) = \mathbb{C}^g/\Lambda$. There is a complex-analytic isomorphism $J(X) \simeq \text{Alb}(X)$ which is obtained by choosing the basis $\{\int_{a_i}\}_{i=1}^g$ for $\check{H}^0(X, K)^\vee$ of the a -periods. The complex coordinates $Z = {}^t(z_1, \dots, z_g)$ on $\text{Alb}(X)$ for this basis are the standard complex coordinates on \mathbb{C}^g and $H_1(X, \mathbb{Z}) \simeq \Lambda$. The invariant Kähler metric on $J(X)$ in these coordinates has the form:

$$ds^2 = \sum_{i,j=1}^g Y^{ij} dz_i \otimes d\bar{z}_j,$$

where $\{Y^{ij}\}_{i,j=1}^g = Y = (\Omega)^{-1}$, so that the corresponding Hermitian inner product on $\mathbb{C}^g = T_0 J(X)$ is given by

$$\langle U, V \rangle = \sum_{i,j=1}^g Y^{ij} u_i \bar{v}_j,$$

where $U = {}^t(u_1, \dots, u_g)$, $V = {}^t(v_1, \dots, v_g)$.

The Abel–Jacobi map $\mu: \text{Div}^0(X) \rightarrow J(X)$ for a Torelli marked Riemann surface X is defined by

$$\mu(D) = {}^t \left(\sum_{i=1}^n \int_{Q_i}^{P_i} \omega_1, \dots, \sum_{i=1}^n \int_{Q_i}^{P_i} \omega_g \right) \in J(X),$$

where $D = \sum_{i=1}^n (P_i - Q_i) \in \text{Div}^0(X)$. By Abel’s theorem, the Abel–Jacobi map establishes a complex-analytic isomorphism $\text{Div}^0(X)/\text{PDiv}(X) \simeq J(X)$.

Two other maps, the holonomy map

$$\text{hol}: \text{Div}^0(X) \rightarrow \text{Hom}(\pi_1(X), \text{U}(1)) = \text{Hom}(H_1(X, \mathbb{Z}), \text{U}(1)),$$

and the inclusion map

$$r: H^1(X, \mathbb{R})/H^1(X, \mathbb{Z}) \rightarrow J(X),$$

are defined as follows. For $D = \sum_{i=1}^l n_i P_i \in \text{Div}^0(X)$ let ω_D be the unique meromorphic differential on X with only simple poles at P_i with residues n_i and with pure imaginary periods:

$$\int_c \omega_D \in \sqrt{-1} \mathbb{R} \text{ for all } [c] \in H_1(X, \mathbb{Z}).$$

The differential of the third kind ω_D can be also interpreted in terms of the holomorphic line bundle $L = [D]$ associated with the divisor D as a unique connection-current ω_D with the property that its curvature-current $(\sqrt{-1}/2\pi)\bar{\partial}\omega_D$ is Poincaré dual to the 0-cycle D . Set

$$\Pi_a = {}^t \left(\int_{a_1} \omega_D, \dots, \int_{a_g} \omega_D \right), \quad \Pi_b = {}^t \left(\int_{b_1} \omega_D, \dots, \int_{b_g} \omega_D \right) \in \mathbb{R}^g,$$

and define the holonomy map $\text{hol}(D) = \rho \in \text{Hom}(\pi_1(X), \text{U}(1))$ by

$$\rho(a_i) = e^{2\pi\sqrt{-1}\Pi_{ai}}, \quad \rho(b_i) = e^{2\pi\sqrt{-1}\Pi_{bi}}, \quad i = 1, \dots, g.$$

To define the inclusion map ι consider the isomorphism

$$H^1(X, \mathbb{R})/H^1(X, \mathbb{Z}) \simeq \text{Hom}(\pi_1(X), \text{U}(1)),$$

given by the exponential of the period map

$$\begin{aligned} H^1(X, \mathbb{R})/H^1(X, \mathbb{Z}) \ni [\omega] &\mapsto e^{2\pi\sqrt{-1}\int_c \omega} \\ &:= e^{2\pi\sqrt{-1}\Pi_c} \in \text{Hom}(\pi_1(X), \text{U}(1)), \end{aligned}$$

and set

$$\iota([\omega]) = -\Omega\Pi_a + \Pi_b \bmod \Lambda.$$

The reciprocity law between differentials of the first and third kinds (see, e.g., [25]) gives

$$\mu(D) = \iota(\text{hol}(D)).$$

These relations between the Abel–Jacobi map, the holonomy map and the inclusion map can be summarized in the following commutative diagram

$$\begin{array}{ccccc} \text{Pic}^0(X) & \xrightarrow{\simeq} & \text{Alb}(X) & \xrightarrow{\simeq} & \text{Div}^0(X)/\text{PDiv}(X) \\ \downarrow \text{hol} & & \downarrow v & & \downarrow \mu \\ \text{Hom}(\pi_1(X), \text{U}(1)) & \xrightarrow{\simeq} & H^1(X, \mathbb{R})/H^1(X, \mathbb{Z}) & \xrightarrow{\iota} & J(X) \end{array}$$

where the isomorphism v is defined by the commutativity of the diagram.

Let $\theta(Z|\Omega)$ be the Riemann theta-function of a Torelli marked Riemann surface X , defined by the following absolutely convergent series,

$$\theta(Z|\Omega) := \sum_{n \in \mathbb{Z}^g} e^{\pi\sqrt{-1}((\Omega n, n) + 2(n, Z))}, \quad Z \in \mathbb{C}^g,$$

where $(A, B) = {}^t AB$ for $A, B \in \mathbb{C}^g$, and denote by $\theta[\xi](Z|\Omega)$ the theta-function with characteristics $\xi = {}^t(\xi_a, \xi_b) \in \mathbb{R}^{2g}$

$$\theta[\xi](Z|\Omega) := e^{\pi\sqrt{-1}((\Omega\xi_a, \xi_a) + 2(\xi_a, Z + \xi_b))} \theta(Z + \Omega\xi_a + \xi_b | \Omega).$$

The Riemann theta-function satisfies the modular transformation formula

$$\theta(-\Omega^{-1}Z \mid -\Omega^{-1}) = \left(\det \left(\frac{\Omega}{\sqrt{-1}} \right) \right)^{1/2} e^{\pi\sqrt{-1}(\Omega^{-1}Z, Z)} \theta(Z \mid \Omega).$$

Let $\Theta \subset J(X)$ be theta-divisor — the zero locus of Riemann theta-function on $J(X)$. It can be defined geometrically as a divisor of the unique (up to a translation) holomorphic line bundle over $J(X)$ whose first Chern class is given by the intersection form in $H_1(J(X), \mathbb{Z}) = \Lambda \simeq H_1(X, \mathbb{Z})$. The theta-divisor is even, $\Theta = -\Theta$, and depends on the marking. Denote by $\mathcal{E} = \{L \in \text{Pic}^{g-1}(X) \mid h^0(L) > 0\}$ the so-called canonical theta-divisor. It follows from the Riemann–Roch formula that $L \in \mathcal{E}$ if and only if $K \otimes L^{-1} \in \mathcal{E}$. Next, choose a base point $P_0 \in X$ and denote by $\mu_{g-1}: \text{Div}^{g-1}(X) \rightarrow J(X)$ the corresponding Abel map: $\mu_{g-1}(D) = \mu(D - (g-1)P_0)$ for $D \in \text{Div}^{g-1}(X)$. Let $W_{g-1} = \mu_{g-1}(\mathcal{E}) = \mu_{g-1}(X^{(g-1)}) \subset \text{Pic}^{g-1}(X)$, where $X^{(g-1)}$ is $g-1$ -fold symmetric product of X and we have used the isomorphism $\text{Pic}^{g-1}(X) = \text{Div}^{g-1}(X)/\text{PDiv}(X)$. According to the Riemann theorem, there exists $\kappa \in J(X)$ such that $\Theta = W_{g-1} + \kappa$. A vector of Riemann constants κ depends on a Torelli marking of the Riemann surface X and on the choice of a base-point, and has the property that there exists $\mathcal{K} \in \text{Pic}^{g-1}(X)$ satisfying $\mathcal{K} \otimes \mathcal{K} = K$ such that $\kappa = -\mu_{g-1}(\mathcal{K})$. Using the spin structure \mathcal{K} we identify $\text{Pic}^{g-1}(X) \simeq \text{Pic}^0(X)$ by $L \mapsto L \otimes \mathcal{K}^{-1}$.

1.2. GREEN'S FUNCTIONS

For any Hermitian (i.e. conformal) metric ds^2 on a Riemann surface X let $2\Delta_0 = - * \bar{\partial} * \bar{\partial}$ be the $\bar{\partial}$ -Laplacian acting on functions on X . Here $*$ is the Hodge star-operator for the metric ds^2 , and $\bar{\partial}$ is the $(0, 1)$ -component of the de Rham differential d on X , $d = \partial + \bar{\partial}$. In local coordinates $ds^2 = \rho(z)|dz|^2$ and $\Delta_0 = -\rho(z)^{-1} \partial^2 / \partial z \partial \bar{z}$. The Green's function G of the Laplacian Δ_0 is canonically defined by the following properties (see, e.g., [27]).

G1. $G \in C^\infty(X \times X \setminus \Delta, \mathbb{R})$, where Δ is the diagonal in $X \times X$.

G2. For every $P \in X$ there exists a neighborhood $U \ni P$ such that the function

$$\tilde{G}(P, Q) := G(P, Q) + \frac{1}{\pi} \log |z(P) - z(Q)|^2$$

is smooth in $U \times U$, where z is a local coordinate at U .

G3. For every $Q \in X$ the function $G_Q(P) := G(P, Q)$ on $X \setminus \{Q\}$ satisfies

$$\Delta_0 G_Q = -1 \text{Area}(X),$$

where $\text{Area}(X) = \int_X * 1$ is the area of X .

G4. $\int_X * G_Q = 0$ for every $Q \in X$.

In distributional sense, conditions **G1–G3** can be summarized in the following single equation

$$\Delta_0 G(P, Q) = \delta(P - Q) - \frac{1}{\text{Area}(X)},$$

where Laplacian Δ_0 acts on the first argument of G .

It follows from **G1–G4** that the Green’s function is symmetric: $G(P, Q) = G(Q, P)$.

Next, consider the following tensor on $X \times X$

$$S := -\pi \partial \bar{\partial}' G = -\pi \frac{\partial^2 G}{\partial z \partial \bar{w}} dz \otimes d\bar{w},$$

where ∂ and $\bar{\partial}'$ act on the first and second arguments of G correspondingly. It follows from **G1–G3** that S is a symmetric bidifferential of the second kind on $X \times X$ with biresidue 1. Specifically, S is a symmetric section of the line bundle $K_X \boxtimes K_X$ over $X \times X$, holomorphic on $X \times X \setminus \Delta$ and having a pole of order 2 at Δ with residue 1:

$$S(P, Q) = \frac{dz(P) \otimes dz(Q)}{(z(P) - z(Q))^2} + O(1) \text{ as } P \rightarrow Q.$$

Succinctly, $S \in \check{H}^0(X \times X, K_X \boxtimes K_X(2\Delta))^{S_2}$ and Bires $|_{\Delta} S = 1$. As it follows from **G1–G3** and the Stokes theorem, bidifferential S has the property

$$\text{v.p.} \int_X S_Q \bar{\omega} = 0 \quad \text{for all } Q \in X \quad \text{and} \quad \omega \in \check{H}^0(X, K).$$

Here $S_Q = \iota_Q^* S \in \check{H}^0(X, K(2Q))$ is a pull-back of S by the map $\iota_Q: X \hookrightarrow X \times X$ defined by $\iota_Q(P) = (P, Q)$, and the integral is understood in the principal value sense. This property uniquely characterizes S as a classical Schiffer kernel. The Schiffer kernel is defined canonically and does not depend on the choice of a conformal metric ds^2 on X .

Another symmetric bidifferential B of the second kind with biresidue 1 is canonically associated with a Torelli marked Riemann surface X and is uniquely characterized by the property that for all $Q \in X$ the differential of the second kind $B_Q = \iota_Q^* B$ on X has zero a -periods. Remaining b -periods of B satisfy the relations

$$\int_{b_i} B_Q = 2\pi\sqrt{-1} \omega_i(Q), \quad i = 1, \dots, g,$$

which follow from the reciprocity law between differentials of the first and second kinds (see, e.g., [17, 25]).

The Schiffer kernel and the B -kernel are related as follows [13].

LEMMA 1.1 (The Fay’s formula).

$$B(P, Q) = S(P, Q) + \pi \sum_{i,j=1}^g Y^{ij} \omega_i(P) \omega_j(Q).$$

Proof. It is sufficient to verify that the right-hand side of the Fay's formula has zero a -periods. Writing $S = d'(-\pi\partial G) + \tilde{S}$, where $\tilde{S} = \pi\partial\bar{\partial}G$ is a regular bidifferential on $X \times X$ holomorphic with respect to P and anti-holomorphic with respect to Q , we get by Stokes theorem that the periods of S with respect to the variable Q are the same as periods of \tilde{S} . Next, it follows from properties **G1–G3** and the Stokes theorem that

$$\begin{aligned} \int_X \tilde{S}_P \omega &= -\pi\omega(P) \quad \text{for all } P \in X \text{ and } \omega \in \check{H}^0(X, K), \\ \int_X \tilde{S}_P \partial f &= 0 \quad \text{for all } P \in X \text{ and } f \in C^\infty(X). \end{aligned}$$

Here $\tilde{S}_P = j_P^* \tilde{S} \in \check{H}^0(X, \bar{K})$ is a pull-back of the bidifferential \tilde{S} by the map $j_P: X \hookrightarrow X \times X$, defined by $j_P(Q) = (P, Q)$. These properties uniquely characterize $K := -\tilde{S}/\pi$ as a kernel of the Hodge projection operator $P: L_2^{1,0}(X) \rightarrow H_{\bar{\partial}}^{1,0}(X)$ onto the subspace of harmonic $(1, 0)$ -forms on X , so that

$$K(P, Q) = \sum_{i,j=1}^g Y^{ij} \omega_i(P) \bar{\omega}_j(Q).$$

Therefore

$$\int_{a_j} \tilde{S}_P = -\pi \sum_{i=1}^g Y^{ij} \omega_i(P), \quad j = 1, \dots, g,$$

and the a -periods of the right-hand side of the Fay's formula are indeed 0. \square

Remark 1.2. Using classical terminology (see, e.g., [25]), the kernel K is the Bergman reproducing kernel for the space of $(1, 0)$ -forms on the Riemann surface X .

2. Bosonic Action Functional and Partition Function

For a classical $U(1)$ -valued bosonic field $g \in C^\infty(X, U(1))$ define corresponding field current as the following 1-form on X

$$\mathbf{J} := \frac{1}{2\pi\sqrt{-1}} g^{-1} dg.$$

The 1-form \mathbf{J} is closed, real-valued, and has integral periods: $[\mathbf{J}] \in H^1(X, \mathbb{Z})$. Denoting by $\mathcal{J}(X)$ the set of all 1-forms on X with these properties, we have canonical isomorphism

$$C^\infty(X, U(1))/U(1) \simeq \mathcal{J}(X),$$

where the inverse map is given by the integration: $g(P) = e^{2\pi\sqrt{-1} \int_{p_0}^P \mathbf{J}}$. Consider the

following functional on the space of field currents $\mathcal{J}(X)$

$$S(\mathbf{J}) := \frac{1}{4} \int_X \mathbf{J} \wedge * \mathbf{J},$$

where $*$ is the Hodge star-operator (which for 1-forms on X does not depend on the choice of a conformal metric on X). According to the Hodge decomposition, every $\mathbf{J} \in \mathcal{J}(X)$ can be uniquely written as $\mathbf{J} = d\varphi_0 + h$, where $\varphi_0 \in C^\infty(X, \mathbb{R})$ and $h \in \mathcal{H}^1(X, \mathbb{Z})$ is harmonic 1-form, $dh = d * h = 0$, with integral periods. Therefore

$$\mathcal{J}(X) = C^\infty(X, \mathbb{R})/\mathbb{R} \times \mathcal{H}^1(X, \mathbb{Z}),$$

so that

$$S(\mathbf{J}) = \frac{1}{4} \int_X d\varphi_0 \wedge * d\varphi_0 + \frac{1}{4} \int_X h \wedge * h := S_0(\varphi_0) + S_{\text{inst}}(h).$$

Here the term S_0 is interpreted as an *action functional* of the standard theory of *free bosons* on X , and the term S_{inst} — as a contribution from *instantons*.

Remark 2.1. The field current \mathbf{J} can be also written as $\mathbf{J} = d\varphi$, where φ is an additive multi-valued real function on X with integral periods, i.e. φ is a single-valued function on the universal cover \tilde{X} of X satisfying $\varphi \circ \gamma - \varphi \in \mathbb{Z}$ for all $\gamma \in \pi_1(X)$, where $\pi_1(X)$ acts on \tilde{X} by deck transformations. The classical field φ with these properties is the simplest example of an *instanton configuration* with the *instanton numbers* given by the periods of the 1-form $\mathbf{J} = d\varphi$.

According to the *bosonization procedure* [3, 21], quantum theories of free bosons on X are parameterized by the set $\text{Pic}^{g-1}(X) \setminus \mathcal{E}$ of generalized spin structures without *zero modes*. For every $L \in \text{Pic}^{g-1}(X) \setminus \mathcal{E}$, following [3], define the *topological term* of the bosonic action functional by

$$S_{\text{top}}(\mathbf{J}) := \sqrt{-1} \int_X \mathbf{J} \wedge \theta_L + \frac{\sqrt{-1}}{2} \epsilon(h) = \sqrt{-1} \int_X h \wedge \theta_L + \frac{\sqrt{-1}}{2} \epsilon(h).$$

Here (see Section 2.1)

$$[\theta_L] := \text{hol}(L_0) \in H^1(X, \mathbb{R})/H^1(X, \mathbb{Z}) \simeq \text{Hom}(\pi_1(X), \text{U}(1)),$$

where $L_0 = L \otimes \mathcal{K}^{-1}$, and $\epsilon(h)$ is a parity of the spin structure associated with $\iota(h)/2$:

$$\epsilon(h) := (l, m) \bmod 2 \quad \text{where } \iota(h) = -\Omega l + m \in \Lambda, \quad l, m \in \mathbb{Z}^g.$$

Since $h \in H^1(X, \mathbb{Z})$, the exponential $\exp\{2\pi S_{\text{top}}(\mathbf{J})\}$ does not depend on the choice of representative θ_L .

The total action functional of the theory of free bosons on X with a topological term parameterized by $L \in \text{Pic}^{g-1}(X) \setminus \mathcal{E}$ is given by

$$S_L(\mathbf{J}) := S(\mathbf{J}) + S_{\text{top}}(\mathbf{J}) = S_0(\varphi_0) + S_{\text{inst}}(h) + S_{\text{top}}(h),$$

and is, in general, complex-valued. The *partition function* of the theory is defined by the following functional integral

$$\langle \mathbf{1} \rangle_L := \int_{\mathcal{J}(X)} [\mathcal{D}\mathbf{J}] e^{-2\pi S_L(\mathbf{J})}.$$

Mathematically rigorous definition is the following. Using the Hodge decomposition, set

$$\langle \mathbf{1} \rangle_L := \mathbf{Z}_0 \mathbf{Z}_{\text{inst}}.$$

Here

$$\mathbf{Z}_0 := \int_{C^\infty(X, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi_0] e^{-2\pi S_0(\varphi_0)}$$

is a *fluctuation part* of $\langle \mathbf{1} \rangle_L$ — the partition function of the quantum field theory of free bosons on X given by the standard Gaussian integral, and

$$\mathbf{Z}_{\text{inst}} := \sum_{h \in \mathcal{H}^1(X, \mathbb{Z})} e^{-\frac{\pi}{2} \int_X h \wedge *h - 2\pi \sqrt{-1} \int_X h \wedge \theta_L - \pi \sqrt{-1} \epsilon(h)}$$

is an *instanton part* of $\langle \mathbf{1} \rangle_L$ with topological term. The instanton part \mathbf{Z}_{inst} is well-defined: it is given by the absolutely convergent series over the lattice $\mathcal{H}^1(X, \mathbb{Z})$ of rank $2g$ in \mathbb{R}^{2g} . In order to define the Gaussian integral for \mathbf{Z}_0 choose a conformal metric ds^2 on X so that the functional S_0 becomes a quadratic form of the Laplacian Δ_0 . According to [34], metric ds^2 defines a ‘Riemannian metric’ on the infinite-dimensional Frechét manifold $C^\infty(X, \mathbb{R})$

$$\|v\|^2 := \int_X *|v|^2, v \in T_{\varphi_0} C^\infty(X, \mathbb{R}),$$

and the integration measure $[\mathcal{D}\varphi_0]$ on $C^\infty(X, \mathbb{R})$ is defined as the volume form of this metric. Mathematically, this is equivalent (see, e.g., [39]) to a definition

$$\mathbf{Z}_0 = \int_{C^\infty(X, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi_0] e^{-2\pi \int_X \Delta_0 \varphi_0 * \varphi_0} := \left(\frac{\text{Area}(X)}{\det_\zeta \Delta_0} \right)^{1/2}.$$

Here $\det_\zeta \Delta_0$ is a functional determinant of the Laplacian Δ_0 of the metric ds^2 defined through the zeta-function of Δ_0 . The area term — contribution from *zero modes* — reflects the integration over the coset $C^\infty(X, \mathbb{R})/\mathbb{R}$. This completes the rigorous definition of the partition function $\langle \mathbf{1} \rangle_L$.

Remark 2.2. Though the functional S_0 does not depend on the choice of the conformal metric ds^2 , this metric is required for defining the integration measure $[\mathcal{D}\varphi_0]$, and the determinant $\det_\zeta \Delta_0$ depends on ds^2 . This is Polyakov’s *conformal anomaly* [32], which is due to the fact that any ‘regularization procedure’ defining $\det \Delta_0$ breaks conformal invariance.

It is remarkable that the instanton part \mathbf{Z}_{inst} can be computed explicitly in terms of the Riemann theta-function with characteristics $\xi = {}^t(-\Pi_a, \Pi_b) \in \mathbb{R}^{2g}$, where

$$e^{2\pi\sqrt{-1}\Pi_{a_i}} = \text{hol}(L_0)(a_i), \quad e^{2\pi\sqrt{-1}\Pi_{b_i}} = \text{hol}(L_0)(b_i), \quad i = 1, \dots, g.$$

Namely, the following statement holds [3].

PROPOSITION 2.3.

$$\mathbf{Z}_{\text{inst}} = \frac{2^{g/2}}{\sqrt{\det Y}} e^{\xi(Z - \bar{Z}, Z - \bar{Z})} |\theta(Z | \Omega)|^2 = \frac{2^{g/2}}{\sqrt{\det Y}} |\theta[\xi](0 | \Omega)|^2,$$

where $Z := \iota(\theta_L) = \Omega \xi_a + \xi_b = -\Omega \Pi_a + \Pi_b \in \mathbb{C}^g$.

Proof. It is a straightforward application (cf. [3]) of the Poisson summation formula

$$\sum_{n \in \mathbb{Z}^{2g}} f(n) = \sum_{n \in \mathbb{Z}^{2g}} \hat{f}(n),$$

where $f \in \mathcal{S}(\mathbb{R}^{2g})$ is a function of the Schwartz class and

$$\hat{f}(p) := \int_{\mathbb{R}^{2g}} e^{-2\pi\sqrt{-1}\langle x, p \rangle} f(x) d^{2g}x$$

is its Fourier transform, where $(,)$ is the standard Euclidean inner product in \mathbb{R}^{2g} . Let $\{\alpha_i, \beta_i\}_{i=1}^g$ be the basis for the lattice $\mathcal{H}^1(X, \mathbb{Z})$ dual to a symplectic basis $\{a_i, b_i\}_{i=1}^g$ in $H_1(X, \mathbb{Z})$, so that for $h \in \mathcal{H}^1(X, \mathbb{Z})$

$$h = \sum_{i=1}^g (l_i \alpha_i + m_i \beta_i) \text{ and } \theta_L = \sum_{i=1}^g (\xi_{a_i} \alpha_i + \xi_{b_i} \beta_i).$$

Using $\epsilon(h) = (l, m) \bmod 2$ and the formulas

$$\int_X h \wedge *h = \langle \lambda, \lambda \rangle, \quad 2\sqrt{-1} \int_X h \wedge \theta_L = \langle Z, \lambda \rangle - \langle \lambda, Z \rangle,$$

where $\lambda := -\Omega l + m \in \Lambda$, the instanton part \mathbf{Z}_{inst} can be represented as the following theta-series

$$\mathbf{Z}_{\text{inst}} = \sum_{n \in \mathbb{Z}^{2g}} e^{-\frac{\pi}{2}(Qn, n) - \pi(A, n)}.$$

Here $n = \begin{pmatrix} l \\ m \end{pmatrix} \in \mathbb{Z}^{2g}$, Q is the following $2g \times 2g$ matrix

$$Q = \begin{pmatrix} \bar{\Omega} Y \Omega & -\bar{\Omega} Y \\ -Y \Omega & Y \end{pmatrix}, \quad \text{and} \quad A = \begin{pmatrix} \Omega Y \bar{Z} - \bar{\Omega} Y Z \\ Y(Z - \bar{Z}) \end{pmatrix} \in \mathbb{C}^{2g}.$$

We apply the Poisson summation formula to the function

$$f(x) = e^{-\frac{\pi}{2}(Qx, x) - \pi(A, x)}.$$

The inverse matrix Q^{-1} is readily computed

$$Q^{-1} = \frac{\sqrt{-1}}{2} \begin{pmatrix} \Omega^{-1} & I \\ I & \bar{\Omega} \end{pmatrix}$$

(this is the place where the extra term $\epsilon(h)$ in the topological action is crucial, cf. [3]), so that

$$\det Q = 2^g \det Y \det(-\sqrt{-1}\Omega).$$

The Gaussian integral for \hat{f} is computed explicitly

$$\hat{f}(x) = \left(\frac{2^g}{\det(-\sqrt{-1}Y\Omega)} \right)^{1/2} e^{\frac{g}{2}(Q^{-1}A,A) - 2\pi(Q^{-1}(x+\sqrt{-1}A),x)},$$

so that

$$\begin{aligned} \sum_{n \in \mathbb{Z}^{2g}} \hat{f}(n) &= \left(\frac{2^g}{\det(-\sqrt{-1}Y\Omega)} \right)^{1/2} e^{\frac{g}{2}(Z-\bar{Z},Z-\bar{Z}) - 2\sqrt{-1}(\Omega^{-1}Z,Z)} \times \\ &\times \bar{\theta}(Z|\Omega)\theta(-\Omega^{-1}Z | -\Omega^{-1}). \end{aligned}$$

Using the modular transformation formula for the Riemann theta-function completes the proof. □

Remark 2.4. Note that though the action functional $S_L(\mathbf{J})$ is not real-valued, Proposition 2.3 shows that the partition function $\langle \mathbf{1} \rangle_L$ is real and positive. Since under the involution $L \mapsto K \otimes L^{-1}$, or equivalently, $Z \mapsto -Z$, $S_L(\mathbf{J}) \mapsto \overline{S_L(\mathbf{J})}$ — the complex conjugate, the partition function $\langle \mathbf{1} \rangle_L$ has a \mathbb{Z}_2 -symmetry $Z \mapsto -Z$, in accordance with $\theta(Z|\Omega) = \theta(-Z|\Omega)$.

Remark 2.5. As it follows from the Proposition 2.3 and [3,14], the partition function $\langle \mathbf{1} \rangle_L$ coincides with the analytic torsion $T(L)$ for the unitary line bundle $L_0 = L \otimes \mathcal{K}^{-1}$ — the zeta-function determinant of the $\bar{\partial}$ -Laplacian acting on sections of the line bundle L . This is the essence of the spin-1/2 bosonization formula [1, 6].

Let $\tau(\mathbf{t}, \mathbf{U})$, $\mathbf{t} = (t_1, t_2, \dots)$, be Sato's τ -function on the universal Grassmannian manifold, defined as a solution of Hirota's bilinear equations (see, e.g., [21]). It is well known [2, 21, 43] that if a point \mathbf{U} in the universal Grassmannian manifold is the image of a Torelli marked Riemann surface X with the line bundle $L \in \text{Pic}^{g-1}(X) \setminus \mathcal{E}$ under the Krichever map, then

$$\tau(\mathbf{t}, \mathbf{U})|_{\mathbf{t}=0} = \tau(Z) = \theta[\xi](0|\Omega),$$

$$Z = \Omega \xi_a + \xi_b.$$

We define the τ -function $\tau(Z)$ of a Torelli marked Riemann surface X with the line bundle $L \in \text{Pic}^{g-1}(X) \setminus \mathcal{E}$ as the partition function $\langle \mathbf{1} \rangle_L$. According to the

Proposition 2.3,

$$\tau(Z) := \langle \mathbf{1} \rangle_L = \frac{2^{g/2}}{\sqrt{\det Y}} |\tau(Z)|^2.$$

3. U(1)-gauge Symmetry Ward Identities

Introduce holomorphic and anti-holomorphic components of the bosonic field current \mathbf{J} through the decomposition

$$\mathbf{J} = \frac{\sqrt{-1}}{2\pi} (J + \bar{J}),$$

so that $J = -g^{-1}\partial g$ and $\bar{J} = -g^{-1}\bar{\partial}g$. By definition, *multi-point correlation functions* of current components J and \bar{J} are given by the following functional integral

$$\begin{aligned} & \langle J(P_1) \dots J(P_m) \bar{J}(Q_1) \dots \bar{J}(Q_n) \rangle \\ & := \int_{\mathcal{J}(X)} [\mathcal{D}\mathbf{J}] J(P_1) \dots J(P_m) \bar{J}(Q_1) \dots \bar{J}(Q_n) e^{-2\pi S(\mathbf{J})}. \end{aligned}$$

We start with the computation of the *normalized* 1-point correlation function of the holomorphic component J , defined by

$$\langle\langle J(P) \rangle\rangle := \frac{\langle J(P) \rangle}{\langle \mathbf{1} \rangle_L}.$$

First observe that for the standard quantum field theory of free bosons on X with the action functional $S_0(\varphi_0)$ all multi-point current correlation functions with odd number of components $\partial\varphi_0$ and $\bar{\partial}\varphi_0$ vanish. This is a standard fact from the theory of Gaussian integration (for mathematical treatment see, e.g., [16]). In particular, $\langle \partial\varphi_0(P) \rangle = 0$ and from the Hodge decomposition $J = -2\pi\sqrt{-1}(\partial\varphi_0 + h^{1,0})$ it follows that

$$\langle\langle J(P) \rangle\rangle = -2\pi\sqrt{-1}\langle\langle h^{1,0}(P) \rangle\rangle.$$

Second, consider the partial derivative $\partial\tau/\partial z_i$ of the τ -function. In this case differentiation under the functional integral sign is legitimate since the instanton part \mathbf{Z}_{inst} — the only part that actually depends on Z — is given by the absolutely convergent series. Using the explicit form of the topological term $S_{\text{top}}(\mathbf{J})$ and the equation

$$-\frac{\partial\langle Z, \lambda \rangle}{\partial z_i} = -\sum_{j=1}^g Y^{ij} \bar{\lambda}_j = -2\sqrt{-1} \int_{a_i} h^{1,0},$$

we obtain

$$\frac{\partial \log \tau}{\partial z_i} = \int_{a_i} \langle \langle J(P) \rangle \rangle, \quad i = 1, \dots, g.$$

This is a U(1)-*gauge symmetry* Ward identity for the 1-point correlation function $\langle \langle J(P) \rangle \rangle$. Using the formula

$$\frac{\partial \log \tau}{\partial z_i} = \sum_{j=1}^g \pi Y^{ij}(z_j - \bar{z}_j) + \frac{\partial \log \theta(Z | \Omega)}{\partial z_i} = \frac{\partial \log \theta[\xi](U | \Omega)}{\partial u_i} \Big|_{U=0},$$

where $U = (u_1, \dots, u_g)$, and taking into account that $\langle \langle J(P) \rangle \rangle$ is an abelian differential on X , we finally obtain

$$\langle \langle J(P) \rangle \rangle = \sum_{i=1}^g \frac{\partial \log \theta[\xi](U | \Omega)}{\partial u_i} \Big|_{U=0} \omega_i(P).$$

This expression is in perfect agreement with spin-1/2 bosonization formulas in [21, 36, 37].

Remark 3.1. In this derivation it was assumed that 1-point correlation function $\langle \langle J(P) \rangle \rangle$ is holomorphic on X . It is one of the basic principles of conformal field theory that multi-point correlation functions of holomorphic and anti-holomorphic current components are multi-differentials on X , holomorphic and anti-holomorphic with respect to the corresponding variables. In our case the action functional $S(\mathbf{J})$ is quadratic and the statement immediately follows from the representation for $\langle h^{1,0}(P) \rangle$, since h is harmonic.

Similarly,

$$\frac{\partial \langle \lambda, Z \rangle}{\partial \bar{z}_i} = \sum_{j=1}^g Y^{ij} \lambda_j = -2\sqrt{-1} \int_{a_i} h^{0,1},$$

and we obtain

$$\frac{\partial \log \tau}{\partial \bar{z}_i} = \int_{a_i} \langle \langle \bar{J}(P) \rangle \rangle, \quad i = 1, \dots, g$$

— a U(1)-*gauge symmetry* Ward identity for the 1-point correlation function $\langle \langle \bar{J}(P) \rangle \rangle$. As before,

$$\langle \langle \bar{J}(P) \rangle \rangle = \sum_{i=1}^g \frac{\partial \log \overline{\theta[\xi]}(U | \Omega)}{\partial \bar{u}_i} \Big|_{U=0} \bar{\omega}_i(P).$$

Remark 3.2. There is no contradiction between the equations $\bar{J}(P) = -\overline{J(P)}$ and $\langle \langle \bar{J}(P) \rangle \rangle = \overline{\langle \langle J(P) \rangle \rangle}$ because the action functional $S_L(\mathbf{J})$ is not real-valued. Indeed, introducing explicit Z -dependence into the action functional and correlation

functions, we get

$$\begin{aligned}\langle\langle\bar{J}(P)\rangle\rangle_Z &= \int_{\mathcal{J}(X)} [\mathcal{D}\mathbf{J}]\bar{J}(P) e^{-2\pi S(\mathbf{J};Z)} \\ &= - \int_{\mathcal{J}(X)} [\mathcal{D}\mathbf{J}]\overline{J(P)} e^{-2\pi\overline{S(\mathbf{J};-Z)}} \\ &= -\overline{\langle\langle J(P)\rangle\rangle_{-Z}} = \overline{\langle\langle J(P)\rangle\rangle_Z},\end{aligned}$$

since $\partial \log \theta(Z|\Omega)/\partial z_i$ is odd.

Now consider the normalized *reduced* 2-point correlation function

$$\langle\langle J(P)J(Q)\rangle\rangle := \frac{\langle J(P)J(Q)\rangle}{(\mathbf{1})_L} - \langle\langle J(P)\rangle\rangle\langle\langle J(Q)\rangle\rangle.$$

First observe that

$$\langle\langle J(P)J(Q)\rangle\rangle = -4\pi^2(\langle\langle \partial\varphi_0(P)\partial\varphi_0(Q)\rangle\rangle + \langle\langle h^{1,0}(P)h^{1,0}(Q)\rangle\rangle).$$

Next, use another standard fact from the theory of Gaussian integration [16]:

$$\langle\langle \partial\varphi_0(P)\partial\varphi_0(Q)\rangle\rangle = \frac{1}{4\pi} \partial\partial' G(P, Q) = -\frac{1}{4\pi^2} S(P, Q),$$

and consider the second partial derivative $\partial^2 \log \tau / \partial z_i \partial z_j$. Arguing as before, we obtain

$$\frac{\partial^2 \log \tau}{\partial z_i \partial z_j} = \int_{a_i} \int_{a_j} (\langle\langle J(P)J(Q)\rangle\rangle - S(P, Q)), \quad i, j = 1, \dots, g.$$

This is a U(1)-gauge symmetry Ward identity for the 2-point correlation function $\langle\langle J(P)J(Q)\rangle\rangle$. It follows from the derivation and discussion in Section 2.2 that the integrand is a holomorphic symmetric bidifferential on $X \times X$. Therefore $\langle\langle J(P)J(Q)\rangle\rangle$ is a symmetric bidifferential of the second kind on $X \times X$ with biresidue 1 on the diagonal. Using the equation

$$\frac{\partial^2 \log \tau}{\partial z_i \partial z_j} = \pi Y^{ij} + \frac{\partial^2 \log \theta(Z|\Omega)}{\partial z_i \partial z_j} = \pi Y^{ij} + \frac{\partial^2 \log \theta[\xi](U|\Omega)}{\partial u_i \partial u_j} \Big|_{U=0}$$

and the Fay's formula, we finally obtain

$$\langle\langle J(P)J(Q)\rangle\rangle = B(P, Q) + \sum_{i,j=1}^g \frac{\partial^2 \log \theta[\xi](U|\Omega)}{\partial u_i \partial u_j} \Big|_{U=0} \omega_i(P)\omega_j(Q).$$

Again, this expression is in perfect agreement with formulas in [21, 36, 37].

Next, consider 2-point correlation function $\langle\langle J(P)\bar{J}(Q)\rangle\rangle$ with holomorphic and anti-holomorphic components. Arguing as before we obtain

$$\frac{\partial^2 \log \tau}{\partial z_i \partial \bar{z}_j} = \int_{a_i} \int_{a_j} (\langle\langle J(P)\bar{J}(Q)\rangle\rangle - K(P, Q)),$$

where

$$\langle\langle \partial\varphi_0(P)\bar{\partial}\varphi_0(Q) \rangle\rangle = \frac{1}{4\pi} \partial\bar{\partial}G(P, Q) = \frac{1}{4\pi^2} K(P, Q),$$

and the Bergman reproducing kernel K was introduced in Section 2.2. Using explicit representation for K and the formula

$$\frac{\partial^2 \log \tau}{\partial z_i \partial \bar{z}_j} = -\pi Y^{ij}$$

we get

$$\langle\langle J(P)\bar{J}(Q) \rangle\rangle = 0.$$

To summarize, introduce multi-point reduced normalized correlation functions, defined inductively as follows (see, e.g., [41])

$$\begin{aligned} & \langle\langle J(P_1) \dots J(P_m)\bar{J}(Q_1) \dots \bar{J}(Q_n) \rangle\rangle \\ & := \frac{\langle\langle J(P_1) \dots J(P_m)\bar{J}(Q_1) \dots \bar{J}(Q_n) \rangle\rangle}{\langle \mathbf{1} \rangle_L} - \\ & \quad - \sum_{l=2}^{m+n} \sum_{I=I_l} \sum_{\bigcup \dots \bigcup I_l} \langle\langle J(I_1) \rangle\rangle \dots \langle\langle J(I_l) \rangle\rangle. \end{aligned}$$

Here $I = \{P_1, \dots, P_m; Q_1, \dots, Q_n\}$ and summation goes over all partitions of the set I into the non-empty subsets I_1, \dots, I_l ; for every subset $I_k = \{P_{i_1}, \dots, P_{i_{m_k}}; Q_{j_1}, \dots, Q_{j_{n_k}}\}$ of I

$$\langle\langle J(I_k) \rangle\rangle := \langle\langle J(P_{i_1}) \dots J(P_{i_{m_k}})\bar{J}(Q_{j_1}) \dots \bar{J}(Q_{j_{n_k}}) \rangle\rangle.$$

We have the following result (cf. [21, 37]):

THEOREM 3.3.

$$\begin{aligned} & \langle\langle J(P_1) \dots J(P_m)\bar{J}(Q_1) \dots \bar{J}(Q_n) \rangle\rangle \\ & = \sum_{i_1=1}^g \dots \sum_{i_m=1}^g \sum_{j_1=1}^g \dots \sum_{j_n=1}^g \frac{\partial^{m+n} \log |\theta(U|\Omega)|^2}{\partial u_{i_1} \dots \partial u_{i_m} \partial \bar{u}_{j_1} \dots \partial \bar{u}_{j_n}} \Big|_{U=0} \times \\ & \quad \times \omega_{i_1}(P_1) \dots \omega_{i_m}(P_m) \bar{\omega}_{j_1}(Q_1) \dots \bar{\omega}_{j_n}(Q_n) + \\ & \quad + \delta_{m,2} \delta_{n,0} B(P_1, P_2) + \delta_{m,0} \delta_{n,2} \overline{B(Q_1, Q_2)}. \end{aligned}$$

Proof. Directly follows from the definition of normalized reduced correlations functions and the arguments used above. Note that when $mn > 0$ all reduced correlation functions vanish. \square

PART 2. FREE BOSONS AND TAU-FUNCTIONS FOR CLOSED SMOOTH JORDAN CURVES

1. Mathematical Set-up

Let C be a smooth closed Jordan curve in the complex plane \mathbb{C} encircling the origin 0. Denote by Ω the interior of a contour C — a simply connected domain in \mathbb{C} containing 0 and bounded by C , and denote by \mathcal{C} the set of all such contours C .

1.1. RIEMANN MAPPING THEOREM

Let $D = \{w \in \mathbb{C} \mid |w| < 1\}$ be the unit disk in the complex plane \mathbb{C} . By Riemann mapping theorem, there exists a conformal isomorphism $F: \Omega \rightarrow D$, uniquely characterized by the conditions $F(0) = 0$ and $F'(0) > 0$. By Carathéodory's correspondence of the boundaries principle F extends to the regular map between the closure $\bar{\Omega}$ of the domain Ω and the closed unit disk \bar{D} , and $F|_C$ is a C^∞ -isomorphism between the contour $C = \partial\Omega$ and the unit circle $S^1 = \partial D$. The inverse map $f = F^{-1}: D \rightarrow \Omega$ is a univalent function on D , smooth up to the boundary. The value $r = F'(0) = f'(0)$ is called the conformal radius of C with respect to 0.

Denote by \mathcal{C}_1 the set of all contours of conformal radius 1 and by \mathcal{F}_1 — the set of all univalent functions on D which are smooth up to the boundary and normalized by the conditions $f(0) = 0$ and $f'(0) = 1$. Let \mathcal{A}_1 be the affine space of all holomorphic functions on D which are smooth up to the boundary and have the same normalization at 0 as functions in \mathcal{F}_1 . The space \mathcal{A}_1 has a structure of an infinite-dimensional complex manifold with Frechét topology given by the uniform convergence with all derivatives in the closed unit disk \bar{D} . Taylor coefficients of the power series expansion at $z = 0$ are natural global coordinates on \mathcal{A}_1 . The set \mathcal{F}_1 is open in \mathcal{A}_1 and, therefore, is an infinite-dimensional complex manifold. By Riemann mapping theorem there is a bijection $\mathcal{C}_1 \simeq \mathcal{F}_1$, which turns \mathcal{C}_1 into an infinite-dimensional complex manifold as well.

Let $\text{Diff}_+(S^1)$ be a group of all orientation preserving diffeomorphisms of the unit circle S^1 , with S^1 interpreted as a rotation subgroup of $\text{Diff}_+(S^1)$. It was observed by A.A. Kirillov [22] that there is a remarkable bijection between \mathcal{C}_1 and the infinite-dimensional homogeneous manifold $\text{Diff}_+(S^1)/S^1$. It turns $\text{Diff}_+(S^1)/S^1$ into an infinite-dimensional complex manifold and gives rise to canonical isomorphisms $\mathcal{C}_1 \simeq \mathcal{F}_1 \simeq \text{Diff}_+(S^1)/S^1$.

In addition to the conformal map F , Kirillov's construction of the correspondence $C \mapsto \gamma$ uses a conformal map G of the exterior domain $\mathbb{C} \setminus \Omega$ of the contour C onto the exterior domain $\mathbb{C} \setminus D$ of the unit circle S^1 , normalized by the conditions $G(\infty) = \infty$ and $G'(\infty) > 0$.^{*} Namely, if $g = G^{-1}$ is the inverse map to G , then $\gamma := F \circ g \in \text{Diff}_+(S^1)$. To prove that $C \mapsto \gamma$ is onto and to construct the inverse map, A. A. Kirillov [22] used the Riemann theorem that all complex structures

^{*}The value $d = 1/G'(\infty)$ is called the transfinite diameter of C .

on the two-dimensional sphere S^2 are isomorphic to the standard complex structure of the Riemann sphere \mathbb{P}^1 .

The manifold $\text{Diff}_+(S^1)/S^1$ admits an interpretation as a set of all complex structures on the space $\Omega \mathbb{R}^d$ of based loops in \mathbb{R}^d and as such plays a fundamental role in the string theory [7, 8].

Remark 1.1. Similarly [23, 31] it can be shown that another homogeneous space $\text{Diff}_+(S^1)/\text{Möb}(S^1)$ is also an infinite-dimensional complex manifold and $\text{Diff}_+(S^1)/S^1$ is a holomorphic disk fiber space over $\text{Diff}_+(S^1)/\text{Möb}(S^1)$.

There is a natural inclusion

$$\text{Diff}_+(S^1)/\text{Möb}(S^1) \hookrightarrow T(1),$$

where $T(1)$ is a classical Bers universal Teichmüller space

$$T(1) := \text{Homeo}_{qs}(S^1)/\text{Möb}(S^1)$$

(see [5] for definitions and details). The space $T(1)$ has a natural structure of an infinite-dimensional complex manifold and it was shown by S. Nag and A. Verjovsky [31] that the inclusion map is holomorphic (moreover, $\text{Diff}_+(S^1)/\text{Möb}(S^1)$ is one leaf of a holomorphic foliation on $T(1)$). It was also shown in [31] that the pull-back by the inclusion map of the canonical Weil–Petersson metric on $T(1)$ to $\text{Diff}_+(S^1)/\text{Möb}(S^1)$ coincides with the Kähler metric introduced by A. A. Kirillov and D. V. Yuriev [22,23] via the orbit method.

Let \mathbb{P}_C^1 be the double of the open Riemann surface $\mathbb{P}^1 \setminus \Omega$, defined by gluing together $\mathbb{P}^1 \setminus \Omega$ and its copy $\overline{\mathbb{P}^1 \setminus \Omega}$ with the opposite complex structure along their common boundary C . As a smooth manifold \mathbb{P}_C^1 is diffeomorphic to the two-dimensional sphere S^2 . The complex structure on \mathbb{P}_C^1 is defined by the structure sheaf $\mathcal{O} = \{\mathcal{O}_U\}$ – the sheaf of germs of holomorphic functions on \mathbb{P}_C^1 . Namely, for any connected open set $U \subset \mathbb{P}_C^1$ define the complex vector space \mathcal{O}_U as follows.

- (a) If $U \subset \mathbb{P}^1 \setminus \Omega$ then \mathcal{O}_U is a complex vector space of all holomorphic functions on U with respect to the standard complex structure on \mathbb{P}^1 .
- (b) If $U \subset \overline{\mathbb{P}^1 \setminus \Omega}$ then \mathcal{O}_U is a complex vector space of all anti-holomorphic functions on U with respect to the standard complex structure on \mathbb{P}^1 .
- (c) If $U \cap C \neq \emptyset$ then, setting $U_+ = U \cap (\mathbb{P}^1 \setminus \Omega)$, $U_- = U \cap (\overline{\mathbb{P}^1 \setminus \Omega})$, define

$$\mathcal{O}_U = \{(f_+, f_-) \mid f_{\pm} \in \mathcal{O}_{U_{\pm}} \text{ and } f_+(z) = f_-(\bar{z}) \text{ for } z \in U \cap C\}.$$

Finally, for any open set $U \subset \mathbb{P}_C^1$ define \mathcal{O}_U as a direct sum of the corresponding complex vector spaces for the connected components of U . Clearly, complex vector spaces $\{\mathcal{O}_U\}$ form a sheaf \mathcal{O} which canonically defines the complex structure on \mathbb{P}_C^1 . By the Riemann theorem \mathbb{P}_C^1 is complex-isomorphic to the Riemann sphere \mathbb{P}^1 , i.e. there exists a global meromorphic coordinate ζ on \mathbb{P}_C^1 such that $\zeta(\infty) = \infty$, $\zeta(\overline{\infty}) = 0$. In terms of the conformal mapping G this coordinate is given

explicitly by

$$\zeta(z) = \begin{cases} G(z) & \text{if } z \in \mathbb{P}^1 \setminus \Omega, \\ 1/\overline{G(z)} & \text{if } z \in \overline{\mathbb{P}^1 \setminus \Omega}. \end{cases}$$

1.2. GREEN'S FUNCTIONS

Let G be the Green's function^{*} of the $\bar{\partial}$ -Laplacian Δ_0 of the conformal metric ds^2 on \mathbb{P}^1 acting on functions on \mathbb{P}^1 , and let S and B be corresponding bidifferentials (see Section 1.2 in Part 1). Since \mathbb{P}^1 has genus 0, the Fay's formula reads $S = B$ and we have explicitly

$$S(z, w) := -\pi \frac{\partial^2 G(z, w)}{\partial z \partial w} dz \otimes dw = \frac{dz \otimes dw}{(z - w)^2}, \quad z, w \in \mathbb{C}.$$

Similarly, since the subspace of harmonic $(1, 0)$ -forms on \mathbb{P}^1 is $\{0\}$, the Bergman reproduction kernel K for \mathbb{P}^1 vanishes

$$K(z, \bar{w}) := -\frac{\partial^2 G(z, w)}{\partial z \partial \bar{w}} dz \otimes d\bar{w} = 0.$$

Let G_{DBC} be the Green's function for the Laplacian Δ_0 on the exterior domain $\mathbb{P}^1 \setminus \Omega$ with the Dirichlet boundary condition $\varphi|_C = 0$. One has

$$G_{\text{DBC}}(z, w) = \frac{2}{\pi} G_{\text{cl}}(z, w),$$

where G_{cl} is the classical Green's function for the domain $\mathbb{P}^1 \setminus \Omega$, uniquely characterized by the following properties (see, e.g., [18]).

- CG1.** For every $w \in \mathbb{C} \setminus \Omega$ function $G_{\text{cl}}(z, w) + \log |z - w|$ is harmonic in $(\mathbb{C} \setminus \Omega) \setminus \{w\}$.
- CG2.** $G_{\text{cl}}(z, w) = 0$ for all $w \in \mathbb{P}^1 \setminus \Omega$ and $z \in C$.
- CG3.** $G_{\text{cl}}(\infty, w) = \log |w| + V + O(1/w)$ as $w \rightarrow \infty$, where V is the so-called Robin's constant of the domain Ω .

The classical Green's function for the domain $\mathbb{P}^1 \setminus \Omega$ can be written explicitly as a pull-back by the conformal map G of the classical Green's function for the domain $\mathbb{P}^1 \setminus D$,

$$G_{\text{cl}}(z, w) = \log \left| \frac{1 - G(z)\overline{G(w)}}{G(z) - \overline{G(w)}} \right|.$$

According to Section 1.2 in Part 1, the Schiffer kernel for $\mathbb{P}^1 \setminus \Omega$ is defined as a

^{*}It should be always clear from the content whether G is a conformal map — a function of one variable, or the Green's function — a function of two variables.

symmetric bidifferential S of the second kind on $(\mathbb{P}^1 \setminus \Omega) \times (\mathbb{P}^1 \setminus \Omega)$ with the property

$$\text{v.p.} \int_{\mathbb{P}^1 \setminus \Omega} S(z, w) \wedge \overline{u(w)} = 0 \quad \text{for all } z \in \mathbb{P}^1 \setminus \Omega \text{ and } u \in \mathcal{H}^{1,0}(\mathbb{P}^1 \setminus \Omega),$$

where $\mathcal{H}^{1,0}(\mathbb{P}^1 \setminus \Omega)$ is the complex vector space of holomorphic $(1, 0)$ -forms on the domain $\mathbb{P}^1 \setminus \Omega$. In terms of the Green's function the Schiffer kernel is given by

$$S(z, w) = -\pi \frac{\partial^2 G_{\text{DBC}}(z, w)}{\partial z \partial w} dz \otimes dw,$$

and in terms of the conformal map G it has the form

$$S(z, w) = \frac{G'(z)G'(w)}{(G(z) - G(w))^2} dz \otimes dw.$$

Correspondingly, the Bergman reproducing kernel K for the domain $\mathbb{P}^1 \setminus \Omega$ is defined as a regular bidifferential on $(\mathbb{P}^1 \setminus \Omega) \times (\mathbb{P}^1 \setminus \Omega)$, holomorphic with respect to the first variable and anti-holomorphic with respect to the second variable, satisfying the property*

$$\frac{1}{2i} \int_{\mathbb{P}^1 \setminus \Omega} K(z, \bar{w}) \wedge u(w) = u(z) \quad \text{for all } z \in \mathbb{P}^1 \setminus \Omega \text{ and } u \in \mathcal{H}^{1,0}(\mathbb{P}^1 \setminus \Omega).$$

The Bergman reproducing kernel K can be also characterized as a kernel of the projection operator from $L_2^{1,0}(\mathbb{P}^1 \setminus \Omega)$ onto the subspace $\mathcal{H}^{1,0}(\mathbb{P}^1 \setminus \Omega)$, where the Hilbert space $L_2^{1,0}(\mathbb{P}^1 \setminus \Omega)$ is the space of $(1, 0)$ -forms on $\mathbb{P}^1 \setminus \Omega$ with the L_2 -norm

$$\|u\|^2 = \frac{1}{2} \int_{\mathbb{P}^1 \setminus \Omega} u \wedge *u = \frac{i}{2} \int_{\mathbb{P}^1 \setminus \Omega} u \wedge \bar{u}.$$

In terms of the orthonormal basis $\{u_n\}_{n \in \mathbb{N}}$ for the subspace $\mathcal{H}^{1,0}(\mathbb{P}^1 \setminus \Omega)$, the Bergman reproducing kernel can be written as

$$K(z, \bar{w}) = \sum_{n=1}^{\infty} u_n(z) \otimes \overline{u_n(w)},$$

and does not depend on the choice of the basis. In terms of the Green's function the Bergman kernel is given by

$$K(z, \bar{w}) = -\frac{\partial^2 G_{\text{DBC}}(z, w)}{\partial z \partial \bar{w}} dz \otimes d\bar{w},$$

*In this part we revert to the standard complex analyst's notation $i = \sqrt{-1}$.

and in terms of the conformal map G it has the form

$$K(z, \bar{w}) = \frac{G'(z)\overline{G'(w)}}{\pi(1 - G(z)\overline{G(w)})^2} dz \otimes d\bar{w} = \sum_{n=1}^{\infty} u_n(z) \otimes \overline{u_n(w)},$$

where

$$u_n(z) = \sqrt{\frac{n}{\pi}} G^{-n-1}(z) G'(z) dz, \quad n \in \mathbb{N}.$$

1.3. DEFORMATION THEORY

Let G be the conformal map $G: \mathbb{C} \setminus \Omega \rightarrow \mathbb{C} \setminus D$, normalized as in the previous section. Following [10] (see also [9]), we define the analog of Faber polynomials associated with G by the following Laurent expansion at $z = \infty$ in the region $|G(z)| > |w|$

$$\frac{zG'(z)}{G(z) - w} = \sum_{n=0}^{\infty} F_n(w) z^{-n},$$

obtained by substituting Laurent series for $G(z)$

$$G(z) = b_{-1}z + b_0 + \frac{b_1}{z} + \dots$$

into the geometric series for $(G(z) - w)^{-1}$. We call the degree n polynomials F_n Faber polynomials of the conformal map G . In terms of the inverse map $g = G^{-1}$, the Faber polynomials can be written explicitly $F_n(w) = [g^n(w)]_+$, where $[g^n]_+$ is a polynomial part of the Laurent series for g^n . Faber polynomials are uniquely characterized by the property

$$F_n(G(z)) = z^n + O(z^{-1}) \quad \text{as } z \rightarrow \infty.$$

All these facts can be obtained from the representation

$$F_n(w) = g^n(w) + \frac{1}{2\pi i} \int_{S^1} \frac{g^n(z)}{z - w} dz, \quad w \in \mathbb{C} \setminus D,$$

which follows from the Cauchy integral formula.

Remark 1.2. In complex analysis Faber polynomials P_n are usually introduced for the inverse map $g = G^{-1}: \mathbb{C} \setminus D \rightarrow \mathbb{C} \setminus \Omega$ through the expansion [9, 10]

$$\frac{zg'(z)}{g(z) - w} = \sum_{n=0}^{\infty} P_n(w) z^{-n},$$

where $|g(z)| > |w|$.

Following [26, 45] (see also references therein) introduce the harmonic moments of exterior and interior of the contour C by the following formulas

$$t_n = \frac{1}{2\pi i n} \int_C z^{-n} \bar{z} dz, \quad n \in \mathbb{N}, \quad t_0 = \frac{1}{2\pi i} \int_C \bar{z} dz = \frac{1}{\pi} A(\Omega),$$

and

$$v_n = \frac{1}{2\pi i} \int_C z^n \bar{z} dz, \quad n \in \mathbb{N}, \quad v_0 = \frac{2}{\pi} \int_{\Omega} \log |z| \bar{z} d^2 z.$$

Here $A(\Omega)$ is the area of the interior domain Ω with respect to the standard Lebesgue measure $d^2 z = |dz \wedge d\bar{z}|/2$ on \mathbb{C} . According to [26, 45] and references therein, parameters $\{t_0, t_n, \bar{t}_n\}_{n \in \mathbb{N}}$ are local coordinates for the space \mathcal{C} in some neighborhood of the contour C . Consequently, there is a foliation of \mathcal{C} with the leaves \tilde{C}_a of contours of fixed area $a > 0$. The leaves \tilde{C}_a can be considered (at least locally) as infinite-dimensional complex manifolds with complex coordinates $\{t_n\}_{n \in \mathbb{N}}$.

Remark 1.3. There is also a foliation of \mathcal{C} with the leaves C_r of contours of conformal radius $r > 0$. For fixed r the coordinate t_0 can be expressed in terms of t_n, \bar{t}_n provided that $\partial r / \partial t_0 \neq 0$ so that the leaves C_r can be considered (at least locally) as infinite-dimensional complex manifolds. Similarly, there is a foliation of \mathcal{C} with the leaves C^d of contours of transfinite diameter $d > 0$. In this paper we do not address a very interesting problem of global description of complex manifolds C_r and C^d .

The goal of the deformation theory is to describe the tangent vector space $T_C \mathcal{C}$ to the manifold \mathcal{C} at a contour C in terms of the data associated with C . By definition, a deformation of the contour C is a smooth family of contours $\{C_t\}_{t \in (-\varepsilon, \varepsilon)}$ such that $C_0 = C$. The smoothness property is a condition that there exists a parameterization of contours C_t which depends smoothly on t : $C_t = \{z(\sigma, t), \sigma \in \mathbb{R}/2\pi\mathbb{Z}\}$ for every $|t| < \varepsilon$, where $z(\sigma, t) \in C^\infty(\mathbb{R}/2\pi\mathbb{Z} \times (-\varepsilon, \varepsilon))$. Corresponding infinitesimal deformation is defined as the vector field $v = \dot{z}(\sigma) d/d\sigma$ along C , where dot stands for $d/dt|_{t=0}$,

$$\dot{z}(\sigma) := \left. \frac{\partial z(\sigma, t)}{\partial t} \right|_{t=0}.$$

By definition, a deformation C_t is called trivial if it consists of reparameterizations of the contour C , so that the vector field v is tangential to C . The deformation C_t is called infinitesimally trivial if the corresponding vector field v is tangential to C . The tangent vector space $T_C \mathcal{C}$ can be canonically identified with the quotient space of the real vector space of all vector fields along C by the subspace of vector fields tangential to C . Equivalently, $T_C \mathcal{C}$ is a real vector space of normal vector fields to C .

With every deformation C_t one can associate the following 1-form on C

$$\dot{\omega}_C := \frac{\partial \bar{z}}{\partial t} \Big|_{t=0} dz - \frac{\partial z}{\partial t} \Big|_{t=0} d\bar{z}.$$

The 1-form $\dot{\omega}_C$ is a restriction to the contour C of a d^{-1} of the Lie derivative of the standard 2-form $d\bar{z} \wedge dz$ on Ω . It satisfies the following ‘calculus formula’

$$\frac{d}{dt} \Big|_{t=0} \int_{C_t} f(z, \bar{z}, t) dz = \int_C (f \dot{dz} + \frac{\partial f}{\partial \bar{z}} \dot{\omega}_C)$$

for any smooth function f in the domain $(\bigcup_{t \in (-\varepsilon, \varepsilon)} C_t) \times (-\varepsilon, \varepsilon)$.

Remark 1.4. Classically, as it goes back to Volterra and Hadamard, deformations of a contour C are described by

$$z(\sigma, t) = z(\sigma) + t \delta n(\sigma) n(\sigma), \quad \sigma \in \mathbb{R}/2\pi\mathbb{Z},$$

where $z(\sigma)$ defines the contour C , $\delta n(\sigma) \in C^\infty(\mathbb{R}/2\pi\mathbb{Z}, \mathbb{R})$, and $n(\sigma)$ is the outer normal to C . For t sufficiently small this equation defines smooth closed Jordan curve C_t . The relation between two approaches is given by the elementary calculus

$$\delta n ds = \frac{1}{2i} \dot{\omega}_C,$$

where $ds := |z'(\sigma)| d\sigma$ is a 1-form on C .

The basic facts of the deformation theory of contours are summarized in the following theorem. Parts (i) and (ii) were proved in [26, 45] and we refer to them as ‘Krichever’s lemma’.

THEOREM 1.5. (i) *Any deformation C_t of the contour C which does not change harmonic moments of exterior $\{t_0, t_n, \bar{t}_n\}$ is infinitesimally trivial. The parameters $\{t_0 - t_0(C), t_n - t_n(C), \bar{t}_n - \bar{t}_n(C)\}$ are local coordinates on \mathcal{C} in some neighborhood of C .*

(ii) *The following 1-forms on C*

$$\dot{\omega}_C^{(n)} := \frac{\partial \bar{z}}{\partial t_n} \Big|_{t_n=t_n(C)} dz - \frac{\partial z}{\partial t_n} \Big|_{t_n=t_n(C)} d\bar{z}, \quad n \in \{0\} \cup \mathbb{N},$$

extends to meromorphic $(1, 0)$ -forms on the double \mathbb{P}_C^1 of the exterior domain $\mathbb{P}^1 \setminus \Omega$ with a single pole at ∞ of order $n + 1$ if $n \in \mathbb{N}$ and simple poles at ∞ and $\overline{\infty}$ with residues 1 and -1 if $n = 0$. Explicitly,

$$\dot{\omega}_C^{(n)} = d(F_n \circ G), \quad \dot{\omega}_C^{(0)} = d \log G$$

in the domain $\mathbb{P}^1 \setminus \Omega$ and

$$\dot{\omega}_C^{(n)} = d(F_n \circ 1/\bar{G}), \quad \dot{\omega}_C^{(0)} = d \log 1/\bar{G}$$

in the domain $\overline{\mathbb{P}^1 \setminus \Omega}$, where F_n is the n th Faber polynomial of G , $n \in \mathbb{N}$.

(iii) The 1-forms $\dot{\omega}_C^{(n)}$ satisfy the property

$$\frac{1}{2\pi i} \int_C \frac{z^{-m}}{m} \dot{\omega}_C^{(n)} = \delta_{mn},$$

and can be identified with the vector fields $\partial/\partial t_n$. For every $a > 0$ the holomorphic tangent vector space $T'_C \tilde{\mathcal{C}}_a$ to $\tilde{\mathcal{C}}_a$ at C is canonically identified with the complex vector space $\mathcal{M}^{1,0}(\mathbb{P}^1_C)$ of meromorphic $(1, 0)$ -forms on \mathbb{P}^1_C with a single pole at ∞ of order ≥ 2 .

(iv) The holomorphic cotangent vector space $T_C^* \tilde{\mathcal{C}}_a$ to $\tilde{\mathcal{C}}_a$ at C is naturally identified with the complex vector space $\mathcal{H}^{1,0}(\mathbb{P}^1 \setminus \Omega)$ of all holomorphic $(1, 0)$ -forms on $\mathbb{P}^1 \setminus \Omega$ which are smooth up to the boundary, with the pairing $(\cdot, \cdot)_C: T'_C \tilde{\mathcal{C}}_a \otimes T_C^* \tilde{\mathcal{C}}_a \rightarrow \mathbb{C}$ given by

$$(\omega, u)_C := \frac{1}{2\pi i} \int_C d^{-1} u \omega.$$

Differentials dt_n correspond to $(1, 0)$ -forms

$$dt_n(z) := d(z^{-n}/n) = -z^{-n-1} dz.$$

Proof. As in [26, 45], we start with the following Riemann–Hilbert problem. Find functions S_+ and S_- that are holomorphic in the domains Ω and $\mathbb{C} \setminus \Omega$ respectively, are smooth in the corresponding closed domains and on their common boundary C satisfy $S_+(z) - S_-(z) = \bar{z}$ for all $z \in C$. With the normalization $S_-(\infty) = 0$ this problem has a unique solution given by the Cauchy integral

$$S_{\pm}(z) = \frac{1}{2\pi i} \int_C \frac{\bar{w} dw}{w - z},$$

where $z \in \Omega$ for the $+$ sign and $z \in \mathbb{C} \setminus \Omega$ for the $-$ sign.

Now let C_t be a deformation of C satisfying conditions in part (i). By the calculus formula,

$$\dot{S}_{\pm}(z) = \frac{1}{2\pi i} \int_C \frac{\dot{\bar{w}} dw - \bar{w} d\dot{w}}{w - z},$$

so that on C $\dot{\omega}_C = \dot{S}_+ dz - \dot{S}_- dz$. Since in some neighborhood of 0 $S_+(z) = \sum_{n=1}^{\infty} n t_n z^{n-1}$, we see that if $dt_n/dt|_{t=0} = 0$ for all $n \in \mathbb{N}$ then $\dot{S}_+ = 0$ in Ω . This implies that on C , $\dot{\omega}_C = -\dot{S}_- dz|_C$, and $\dot{\omega}_C$ admits holomorphic continuation as a $(1, 0)$ -form on $\mathbb{C} \setminus \Omega$. Since $dt_0/dt|_{t=0} = 0$ and

$$S_-(z) = -\frac{t_0}{z} + O(z^{-2}) \text{ as } z \rightarrow \infty,$$

we get $\dot{S}_-(z) = O(z^{-2})$, so that $\dot{\omega}_C$ is also regular at ∞ . Since, by definition, the 1-form $\dot{\omega}_C$ is pure imaginary on C , by Riemann–Schwarz reflection principle it can be analytically continued to a holomorphic $(1, 0)$ -form on \mathbb{P}_C^1 . Since \mathbb{P}_C^1 has genus 0 we conclude that $\dot{\omega}_C = 0$. In particular, $\dot{\omega}_C = 0$ on C , which by Remark 1.4 is equivalent to the condition that vector field v corresponding to the deformation C_t is tangential to C . This proves part (i).

For the proof of part (ii), set $t = \operatorname{Re}(t_n - t_n(C))$. We have for $n \in \mathbb{N}$

$$\dot{S}_+ = nz^{n-1},$$

so that the 1-form $\dot{\omega}_C - dz^n$ admits holomorphic continuation to the domain $\mathbb{C} \setminus \Omega$ and is regular at ∞ . Similarly, for $n = 0$ the 1-form $\dot{\omega}_C$ is holomorphic on $\mathbb{C} \setminus \Omega$ with a simple pole at ∞ with residue 1. As before, we conclude that 1-form $\dot{\omega}_C$ admits a meromorphic continuation to \mathbb{P}_C^1 with only poles at ∞ and $\overline{\infty}$. For $n = 0$ they are simple poles with residues 1 and -1 , so that using the global coordinate ζ on \mathbb{P}_C^1 we get

$$\dot{\omega}_C^{(0)} = \frac{d\zeta}{\zeta}.$$

In particular,

$$\dot{\omega}_C^{(0)} = d \log G \quad \text{on } \mathbb{P}^1 \setminus \Omega \quad \text{and} \quad \dot{\omega}_C^{(0)} = d \log 1/\overline{G} \quad \text{on } \overline{\mathbb{P}^1 \setminus \Omega}.$$

For $n \in \mathbb{N}$ we get

$$\dot{\omega}_C = d(z^n + O(z^{-1})) \quad \text{as } z \rightarrow \infty \quad \text{and} \quad \dot{\omega}_C = -d(\bar{z}^n + O(z^{-1})) \quad \text{as } z \rightarrow \bar{\infty}.$$

Similarly, setting $t = \operatorname{Im}(t_n - t_n(C))$ we get for $n \in \mathbb{N}$

$$\dot{\omega}_C = id(z^n + O(z^{-1})) \quad \text{as } z \rightarrow \infty \quad \text{and} \quad \dot{\omega}_C = id(\bar{z}^n + O(z^{-1})) \quad \text{as } z \rightarrow \bar{\infty}.$$

From here we conclude that for $n \in \mathbb{N}$ the 1-forms $\dot{\omega}_C^{(n)}$ are meromorphic on \mathbb{P}_C^1 with the only pole at ∞ and

$$\dot{\omega}_C^{(n)} = d(z^n + O(z^{-1})) \quad \text{as } z \rightarrow \infty.$$

Similarly, the 1-forms

$$\dot{\omega}_C^{(\bar{n})} := \left. \frac{\partial \bar{z}}{\partial \bar{t}_n} \right|_{t=0} dz - \left. \frac{\partial z}{\partial \bar{t}_n} \right|_{t=0} d\bar{z}$$

are meromorphic on \mathbb{P}_C^1 with the only pole at $\bar{\infty}$ and

$$\dot{\omega}_C^{(\bar{n})} = -d(\bar{z}^n + O(\bar{z}^{-1})) \quad \text{as } z \rightarrow \bar{\infty}.$$

Using characteristic property of the Faber polynomials F_n we get that on $\mathbb{P}^1 \setminus \Omega$

$$\dot{\omega}_C^{(n)} = d(F_n \circ G) \quad \text{and} \quad \dot{\omega}_C^{(\bar{n})} = -d(\overline{F_n \circ 1/\overline{G}}),$$

and on $\overline{\mathbb{P}^1 \setminus \Omega}$

$$\dot{\omega}_C^{(m)} = d(F_n \circ 1/\overline{G}) \quad \text{and} \quad \dot{\omega}_C^{(\bar{m})} = -d(\overline{F_n \circ G}).$$

The proof of parts (iii)–(iv) is now straightforward. By definition of the harmonic moments t_n , the calculus formula and the Cauchy theorem we get

$$\frac{\partial t_n}{\partial t_m} = \frac{1}{2\pi i n} \int_C z^{-n} \dot{\omega}_C^{(m)} = \delta_{mn},$$

so that 1-forms $\dot{\omega}_C^{(m)}$ correspond to the vector fields $\partial/\partial t_m$ on C . In particular, the same formula shows that corresponding differentials dt_n — (1, 0)-forms on C , are given by the 1-forms $d(z^{-n}/n)$ on C and satisfy

$$\left(\frac{\partial}{\partial t_m}, dt_n \right)_C = \delta_{mn}. \quad \square$$

Remark 1.6. Let $\mathbf{d} = \mathbf{d}' + \mathbf{d}''$ be the decomposition of the de Rham differential on \tilde{C}_a into (1, 0) and (0, 1) components with respect to the complex structure defined by the harmonic moments t_n . According to part (iv) of the theorem, for any smooth function F on \tilde{C}_a the (1, 0)-form

$$\mathbf{d}'F = \sum_{n=1}^{\infty} \frac{\partial F}{\partial t_n} dt_n$$

at $C \in \tilde{C}_a$ can be identified with the holomorphic (1, 0)-form on $\mathbb{P}^1 \setminus \Omega$ defined by the following Laurent expansion at $z = \infty$

$$- \sum_{n=1}^{\infty} z^{-n-1} \frac{\partial F}{\partial t_n} dz.$$

The convergence of this series and holomorphy of the corresponding (1, 0)-form on $\mathbb{P}^1 \setminus \Omega$ follow from the smoothness of the function F — the existence of $\mathbf{d}F$. An example of smooth function F is given by

$$F = \frac{1}{2\pi i} \int_C h(z) \bar{z} dz,$$

where h is holomorphic on $\mathbb{P}^1 \setminus \Omega$ and is smooth up the boundary. In this case, $\mathbf{d}'F = h'(z) dz$. Though we do not address here the question of defining various functional classes on \tilde{C}_a , all said above holds for the class of real-analytic functions.

Here we introduce a natural Hermitian metric on complex manifolds \tilde{C}_a , which turns out to be Kähler, as we shall prove in Section 3. Namely, for every $C \in \tilde{C}_a$ consider the following inner product in the holomorphic tangent vector space $T'_C \tilde{C}_a$

$$H\left(\frac{\partial}{\partial t_m}, \frac{\partial}{\partial t_n}\right) = h^{m\bar{n}} := -\frac{1}{(2\pi i)^2} \int_{C_+} \int_{C_+} z^m \bar{w}^n K(z, \bar{w}),$$

where C_+ is an arbitrary contour containing C inside (not that K is singular as

$z = w \in C$). From the representation of the Bergman kernel in terms of the orthonormal system in Section 1.2 it easily follows that this inner product is positive-definite. Using this orthonormal system it is not difficult to show that components $h^{m\bar{n}}$ are smooth on \tilde{C}_a and thus define the Hermitian metric H . In terms of this metric we get the following Laurent expansion at $z = w = \infty$ for the Bergman reproducing kernel K ,

$$K(z, \bar{w}) = \sum_{m,n=1}^{\infty} h^{m\bar{n}} dt_m(z) \otimes \overline{dt_n(w)}.$$

Remark 1.7. It is instructive to compare the complex structure on C_1 introduced by A. A. Kirillov [22] with the complex structure on \tilde{C}_1 defined (at least locally) by the harmonic moments of exterior. According to [31] (see Remark 1.1) the former is a pull-back by the inclusion map of the Ahlfors-Bers complex structure on $T(1)$ and is defined using quadratic (or, equivalently, Beltrami) differentials, whereas the latter is defined using $(1, 0)$ -forms — the ordinary differentials. Correspondingly, the Hermitian metric on $\text{Diff}_+(S^1)/\text{Möb}(S^1)$ is given by the Petersson inner product of holomorphic quadratic (equivalently, Beltrami) differentials [5, 31], whereas the Hermitian metric on \tilde{C}_1 is defined via the canonical inner product of holomorphic $(1, 0)$ -forms on $\mathbb{P}^1 \setminus \Omega$.

2. Bosonic Action Functional and Partition Function

For a classical field $\varphi \in C^\infty(\mathbb{P}^1, \mathbb{R})$ consider the following action functional

$$S_0(\varphi) := \frac{1}{8} \int_{\mathbb{P}^1} d\varphi \wedge *d\varphi = \frac{i}{4} \int_{\mathbb{P}^1} \partial\varphi \wedge \bar{\partial}\varphi,$$

which describes the standard theory of free bosons on the Riemann sphere \mathbb{P}^1 . Corresponding partition function is defined by the functional integral

$$\langle \mathbf{1} \rangle_0 := \int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi] e^{-\frac{1}{\pi} S_0(\varphi)},$$

where integration goes over the coset $C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}$ and reflects the symmetry $\varphi \mapsto \varphi + c$. As in Part 1, mathematically rigorous definition requires a choice of a conformal metric ds^2 on \mathbb{P}^1 and leads to the result

$$\langle \mathbf{1} \rangle_0 = \int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi] e^{-\frac{1}{2\pi} \int_{\mathbb{P}^1} \Delta_0 \varphi * \varphi} = \left(\frac{\text{Area}(\mathbb{P}^1)}{\det_\zeta \Delta_0} \right)^{1/2}.$$

Here the area term is a contribution from zero modes — the one-dimensional kernel of the $\bar{\partial}$ -Laplacian Δ_0 of the metric ds^2 acting on functions on \mathbb{P}^1 .

For every $C \in \mathcal{C}$ introduce the following analog of the topological term (cf. with the discussion in Section 2 of Part 1):

$$S_{\text{top}}(\varphi) := \int_{\mathbb{C}} (A(\Omega)\delta_0 - \chi_{\Omega})\varphi d^2z = A(\Omega)\varphi(0) - \int_{\Omega} \varphi d^2z.$$

Here χ_{Ω} is a characteristic function of the domain Ω , and δ_0 is a Dirac delta-function at 0 with respect to the Lebesgue measure d^2z . The functional S_{top} has the property $S_{\text{top}}(\varphi + c) = S_{\text{top}}(\varphi)$.

The total bosonic action $S_C(\varphi) := S_0(\varphi) + S_{\text{top}}(\varphi)$ defines the theory of free bosons on \mathbb{P}^1 in the presence of a contour C and we consider a family of such field theories parameterized by \mathcal{C} .

For every $C \in \mathcal{C}$ define the partition function of the corresponding quantum field theory by the following functional integral

$$\langle \mathbf{1} \rangle_C := \int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi] e^{-\frac{1}{\pi} S_C(\varphi)}.$$

Mathematically rigorous definition is the following. Approximate, in the distributional sense, characteristic function χ_{Ω} and Dirac delta-function δ_0 by smooth functions $\chi_{\Omega}^{(\varepsilon)}$ and $\delta_0^{(\varepsilon)}$ with compact supports satisfying

$$\int_{\mathbb{C}} (A(\Omega)\delta_0^{(\varepsilon)} - \chi_{\Omega}^{(\varepsilon)}) d^2z = 0$$

and define

$$\langle \mathbf{1} \rangle_C := \lim_{\varepsilon \rightarrow 0} \exp \left\{ \frac{A^2(\Omega)}{\pi^2} \int_{\mathbb{C}} \int_{\mathbb{C}} \log |z - w| \delta_0^{(\varepsilon)}(z) \delta_0^{(\varepsilon)}(w) d^2z d^2w \right\} \int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi] e^{-\frac{1}{\pi} S_C^{(\varepsilon)}(\varphi)},$$

where

$$S_C^{(\varepsilon)}(\varphi) := S_0(\varphi) + \int_{\mathbb{C}} (A(\Omega)\delta_0^{(\varepsilon)} - \chi_{\Omega}^{(\varepsilon)})\varphi d^2z.$$

We introduce the τ -function $\tau = \tau(C)$ of the smooth Jordan contour C as the *normalized expectation value* of C , defined as follows

$$\tau = \langle \langle C \rangle \rangle := \frac{\langle \mathbf{1} \rangle_C}{\langle \mathbf{1} \rangle_0}.$$

PROPOSITION 2.1. *The τ -function of the contour C is well-defined and is given explicitly by the following expression*

$$\begin{aligned} \log \tau &= -\frac{1}{\pi^2} \int_{\Omega} \int_{\Omega} \log |z - w| d^2 z d^2 w + \frac{2}{\pi^2} A(\Omega) \int_{\Omega} \log |z| d^2 z \\ &= -\frac{1}{\pi^2} \int_{\Omega} \int_{\Omega} \log \left| \frac{1}{z} - \frac{1}{w} \right| d^2 z d^2 w. \end{aligned}$$

Proof. It is another standard computation. Consider the Gaussian integral

$$\langle \mathbf{1} \rangle_C^{(\varepsilon)} := \int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi] e^{-\frac{1}{\varepsilon} S_C^{(\varepsilon)}(\varphi)}.$$

and make the change of variables $\varphi = \Phi^{(\varepsilon)} + \tilde{\varphi}$, where $\Phi^{(\varepsilon)}$ is uniquely determined by the condition that $S_C^{(\varepsilon)}(\Phi^{(\varepsilon)} + \tilde{\varphi})$ does not contain linear terms in $\tilde{\varphi}$ and by the normalization $\Phi^{(\varepsilon)}(\infty) = 0$. Using the Stokes theorem we get

$$-\frac{\partial^2 \Phi^{(\varepsilon)}(z)}{\partial z \partial \bar{z}} = \lambda^{(\varepsilon)}(z),$$

where $\lambda^{(\varepsilon)} := \chi_{\Omega}^{(\varepsilon)} - A(\Omega)\delta_0^{(\varepsilon)}$. The function $\lambda^{(\varepsilon)}$ is smooth, has compact support and $\int_{\mathbb{C}} \lambda^{(\varepsilon)}(z) d^2 z = 0$, so that

$$\Phi^{(\varepsilon)}(z) = -\frac{2}{\pi} \int_{\mathbb{C}} \log |z - w| \lambda^{(\varepsilon)}(w) d^2 w.$$

Since

$$S_C^{(\varepsilon)}(\Phi^{(\varepsilon)} + \tilde{\varphi}) = S_0(\tilde{\varphi}) + \frac{1}{2} \int_{\mathbb{C}} \Phi^{(\varepsilon)}(z) \lambda^{(\varepsilon)}(z) d^2 z$$

we finally obtain

$$\langle \mathbf{1} \rangle_C^{(\varepsilon)} = \langle \mathbf{1} \rangle_0 \exp \left\{ -\frac{1}{\pi^2} \int_{\mathbb{C}} \int_{\mathbb{C}} \log |z - w| \lambda^{(\varepsilon)}(z) \lambda^{(\varepsilon)}(w) d^2 z d^2 w \right\}.$$

Multiplying by the regularization factor and passing to the limit $\varepsilon \rightarrow 0$ we see that $\log \tau_C$ is well-defined and is given by the formula above. \square

COROLLARY 2.2. *The τ -function is $-1/\pi^2$ times regularized energy of the pseudo-measure $d\mu = d^2 z - A(\Omega)\delta_0$ on the domain Ω , where $d^2 z$ is the Lebesgue measure and δ_0 — the delta-measure at 0.*

Proof. Indeed, the energy $I(\nu)$ of a Borel measure ν on Ω is defined by (see, e.g., [38])

$$I(\nu) := \int_{\Omega} \int_{\Omega} \log |z - w| \nu(z) \nu(w).$$

In our case, due to the presence of a delta-measure, we formally have $I(\mu) = -\infty$. However, with the above regularization $I(\mu) = -\pi^2 \tau$. \square

Remark 2.3. It follows from the Proposition 2.1 that the τ -function τ_C coincides with Wiegmann–Zabrodin τ -function τ_{WZ} for the analytic closed Jordan curve C , introduced in [45] and computed in [24]! It is interesting to compare these two approaches. Specifically, in [24, 45] the τ -function τ_{WZ} appears as a dispersionless limit of the Hirota’s τ -function for the integrable two-dimensional Toda hierarchy and also as a large N limit of a partition function of a certain ensemble of random $N \times N$ matrices. In our approach the τ -function τ_C is just a partition function of a quantum field theory of free bosons on \mathbb{P}^1 parameterized by a smooth contour $C \in \mathcal{C}$. This is quite analogous to the definition of the τ -function in Part I as a partition function of a quantum field theory of free bosons on X parameterized by a holomorphic line bundle $L \in \text{Pic}^{g-1}(X) \setminus \mathcal{E}$.

We also consider the quantum theory of free bosons in the exterior domain $\mathbb{P}^1 \setminus \Omega$, defined by the following functional

$$S_{\text{ext}}(\varphi) := \frac{1}{8} \int_{\mathbb{P}^1 \setminus \Omega} d\varphi \wedge *d\varphi = \frac{i}{4} \int_{\mathbb{P}^1 \setminus \Omega} \partial\varphi \wedge \bar{\partial}\varphi,$$

where $\varphi \in C_{\text{DBC}}^\infty(\mathbb{P}^1 \setminus \Omega)$ is smooth (including ∞) and satisfies the Dirichlet boundary condition $\varphi|_C = 0$. The corresponding partition function is defined by the functional integral

$$\langle \mathbf{1} \rangle_{\text{DBC}} := \int_{C_{\text{DBC}}^\infty(\mathbb{P}^1 \setminus \Omega, \mathbb{R})} [\mathcal{D}\varphi] e^{-\frac{1}{\pi} S_{\text{ext}}(\varphi)}$$

and is given explicitly as

$$\langle \mathbf{1} \rangle_{\text{DBC}} = \left(\frac{\text{Area}(\mathbb{P}^1 \setminus \Omega)}{\det_\zeta \Delta_0} \right)^{1/2},$$

where Δ_0 is a $\bar{\partial}$ -Laplacian of the metric ds^2 acting on functions on $\mathbb{P}^1 \setminus \Omega$ satisfying the Dirichlet boundary condition.

3. Current Ward Identities

Introduce holomorphic and anti-holomorphic components of the bosonic field current $d\varphi$ as $d\varphi = j + \bar{j}$, where $j = \partial\varphi$ and $\bar{j} = \bar{\partial}\varphi$. By definition, multi-point correlation functions of current components are given by the following functional integral

$$\begin{aligned} & \langle j(z_1) \dots j(z_m) \bar{j}(w_1) \dots \bar{j}(w_n) \rangle \\ & := \int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [\mathcal{D}\varphi] j(z_1) \dots j(z_m) \bar{j}(w_1) \dots \bar{j}(w_n) e^{-\frac{1}{\pi} S_C(\varphi)}. \end{aligned}$$

Correlation functions for the standard theory of free bosons on \mathbb{P}^1 with action functional S_0 (which formally corresponds to $C = \emptyset$ – the empty set) and for the

theory on $\mathbb{P}^1 \setminus \Omega$ with action functional S_{ext} and Dirichlet boundary condition are defined similarly. We denoted them by $\langle \dots \rangle_0$ and $\langle \dots \rangle_{\text{DBC}}$ respectfully.

We start with the computation of the normalized 1-point correlation function of the holomorphic component J , defined by

$$\langle \langle J(z) \rangle \rangle := \frac{\langle J(z) \rangle}{\langle \mathbf{1} \rangle_C}.$$

Repeating the proof of the Proposition 2.1, which is based on the change of variables $\varphi = \Phi^{(\varepsilon)} + \tilde{\varphi}$, and using the standard fact that $\langle \tilde{j}(z) \rangle_0 = 0$, where $\tilde{j}(z) = \partial \tilde{\varphi}(z)$ (cf. Section 4 in Part 1) we get

$$\int_{C^\infty(\mathbb{P}^1, \mathbb{R})/\mathbb{R}} [D\varphi] J(z) e^{-\frac{1}{\varepsilon} S_C^{(\varepsilon)}(\varphi)} = \langle \mathbf{1} \rangle_C^{(\varepsilon)} \frac{\partial \Phi^{(\varepsilon)}(z)}{\partial z} dz.$$

Passing to the limit $\varepsilon \rightarrow 0$ we obtain

$$\langle \langle J(z) \rangle \rangle = \frac{\partial \Phi(z)}{\partial z} dz,$$

where $\Phi := \lim_{\varepsilon \rightarrow 0} \Phi^{(\varepsilon)}$ and is given explicitly by

$$\Phi(z) = \frac{2A(\Omega)}{\pi} \log |z| - \frac{2}{\pi} \int_{\Omega} \log |z - w| d^2 w.$$

The function $\Phi(z)$ can be characterized as a continuous solution of the equation

$$-\frac{\partial^2 \Phi(z)}{\partial z \partial \bar{z}} = \begin{cases} \chi_{\Omega}(z) - A(\Omega) \delta_0(z) & \text{if } z \in \Omega, \\ 0 & \text{if } z \in \mathbb{C} \setminus \Omega, \end{cases}$$

normalized by the condition $\Phi(\infty) = 0$, and is a logarithmic potential of the pseudo-measure $d\mu = d^2 z - A(\Omega) \delta_0$ on Ω . It also follows from the integral representation that holomorphic on $\mathbb{P}^1 \setminus \Omega$ function $\partial \Phi / \partial z$ coincides with the function $S_{-}(z) + t_0/z$ (see Section 1.3), and has the following Laurent expansion at $z = \infty$

$$\frac{\partial \Phi(z)}{\partial z} = - \sum_{n=1}^{\infty} v_n z^{-n-1}.$$

On the other hand, consider the partial derivative

$$\frac{\partial \log \tau}{\partial t_n} = \frac{1}{\langle \mathbf{1} \rangle_C} \frac{\partial \langle \mathbf{1} \rangle_C}{\partial t_n}, \quad n \in \mathbb{N}.$$

We can evaluate it by differentiating under the functional integral sign, which can be easily justified using the rigorous definition of the partition function $\langle \mathbf{1} \rangle_C$. Since only topological term in the action functional depends on the domain Ω , the computation is based on another calculus formula

$$\frac{\partial}{\partial t_n} \int_{\Omega} \varphi d^2 z = \frac{1}{2i} \int_C \varphi \dot{\omega}_C^{(n)}.$$

Arguing as in the proof of Proposition 2.1 we get

$$\frac{\partial \log \tau}{\partial t_n} = \frac{1}{2\pi i} \left\langle \int_C \varphi \dot{\omega}_C^{(n)} \right\rangle / \langle \mathbf{1} \rangle_C = \frac{1}{2\pi i} \int_C \Phi \dot{\omega}_C^{(n)}.$$

Using Krichever's lemma and integration by parts, we obtain

$$\frac{\partial \log \tau}{\partial t_n} = \frac{1}{2\pi i} \int_C \Phi d(F_n \circ G) = -\frac{1}{2\pi i} \int_C F_n \circ G \left(\frac{\partial \Phi}{\partial z} dz + \frac{\partial \Phi}{\partial \bar{z}} d\bar{z} \right),$$

where $\partial \Phi / \partial z$ and $\partial \Phi / \partial \bar{z}$ are boundary values on C of holomorphic and anti-holomorphic functions on the exterior domain $\mathbb{P}^1 \setminus \Omega$. The latter can be easily justified by considering smooth functions $\Phi^{(\varepsilon)}$ first and then passing to the limit $\varepsilon \rightarrow 0$. Using the characteristic property of Faber polynomials and Cauchy theorem, we have

$$-\frac{1}{2\pi i} \int_C (F_n \circ G) \frac{\partial \Phi}{\partial z} dz = -\frac{1}{2\pi i} \int_C \frac{\partial \Phi}{\partial z} z^n dz = v_n.$$

We claim that the second integral in the formula for $\partial \log \tau / \partial t_n$ is 0. Using that $|G| = 1$ on C , we have

$$\overline{\int_C (F_n \circ G) \frac{\partial \Phi}{\partial \bar{z}} d\bar{z}} = \int_C (\bar{F}_n \circ 1/G) \frac{\partial \Phi}{\partial z} dz,$$

where $\bar{F}_n(z) := \overline{F_n(\bar{z})}$. Since $\partial \Phi / \partial z = O(z^{-2})$ as $z \rightarrow \infty$, the integral indeed vanishes.

Thus, we proved the Ward identity for the 1-point correlation function

$$\frac{\partial \log \tau}{\partial t_n} = -\frac{1}{2\pi i} \int_C z^n \langle \langle J(z) \rangle \rangle, \quad n \in \mathbb{N}.$$

The case $n = 0$ can be considered similarly with the only difference that since $A(\Omega) = \pi t_0$ one needs to differentiate the regularization factor as well. We have the following computation

$$\begin{aligned} \frac{\partial \log \tau}{\partial t_0} &= \lim_{\varepsilon \rightarrow 0} \left(\frac{1}{2\pi i} \int_C \Phi^{(\varepsilon)} \frac{dG}{G} - \int_C \Phi^{(\varepsilon)}(z) \delta^{(\varepsilon)}(z) d^2 z + \right. \\ &\quad \left. + \frac{2A(\Omega)}{\pi} \int_C \int_C \log |z - w| \delta_0^{(\varepsilon)}(z) \delta_0^{(\varepsilon)}(w) d^2 z d^2 w \right) \\ &= \frac{1}{2\pi i} \int_C \Phi \frac{dG}{G} + \lim_{\varepsilon \rightarrow 0} \int_C \left(\frac{2A(\Omega)}{\pi} \log |z| - \Phi^{(\varepsilon)}(z) \right) \delta_0^{(\varepsilon)}(z) d^2 z \\ &= \lim_{z \rightarrow 0} \left(\frac{2A(\Omega)}{\pi} \log |z| - \Phi(z) \right) \\ &= v_0, \end{aligned}$$

with an easy justification of all steps. Here we have also used the equation $\int_C \Phi(dG/G) = 0$, which can be easily proved as follows. Let Ψ be a holomorphic

function on $\mathbb{P}^1 \setminus \Omega$ such that

$$\Psi'(z) = \frac{\partial \Phi(z)}{\partial z} \quad \text{and} \quad \Psi(\infty) = 0.$$

Such function exists since $\partial \Phi(z)/\partial z = O(z^{-2})$ as $z \rightarrow \infty$. We have $\Phi = \Psi + \bar{\Psi}$ on $\mathbb{P}^1 \setminus \Omega$ and by Cauchy theorem $\int_C \Psi(dG/G) = 0$. Using the same argument as for the case $n > 0$ above shows that the integral with $\bar{\Psi}$ also vanishes.

We summarize these results as the following statement (cf. [24]).

THEOREM 3.1. *The normalized 1-point current correlation functions of free bosons on \mathbb{P}^1 parameterized by $C \in \tilde{\mathcal{C}}_a$ for every $a > 0$ satisfy the Ward identities, given by the following Laurent expansions at $z = \infty$*

$$\langle\langle J(z) \rangle\rangle = - \sum_{n=1}^{\infty} z^{-n-1} \frac{\partial \log \tau}{\partial t_n} dz = \mathbf{d}' \log \tau,$$

and

$$\langle\langle \bar{J}(z) \rangle\rangle = - \sum_{n=1}^{\infty} \bar{z}^{-n-1} \frac{\partial \log \tau}{\partial \bar{t}_n} d\bar{z} = \mathbf{d}'' \log \tau.$$

Remark 3.2. Since $\langle J(z) \rangle_{\text{DBC}} = \langle \bar{J}(z) \rangle_{\text{DBC}} = 0$ the theorem can be also trivially interpreted as computing in terms of the τ -function the difference between 1-point correlation functions of free bosons on \mathbb{P}^1 parameterized by C and of free bosons on $\mathbb{P}^1 \setminus \Omega$ with the Dirichlet boundary condition. This will be relevant for the 2-point correlation functions.

COROLLARY 3.3 ([24, 26, 45]). *The function $\log \tau \in C^\infty(\mathcal{C}, \mathbb{R})$ is a generating function for the harmonic moments of interior:*

$$v_0 = \frac{\partial \log \tau}{\partial t_0} \quad \text{and} \quad v_n = \frac{\partial \log \tau}{\partial t_n}, \quad n \in \mathbb{N}.$$

Remark 3.4. It is instructive to compare this corollary with our results with P. Zograf [48, 49] on the solution of the Poincaré problem of accessory parameters (see also [40] for an overview). Namely, let $X := \mathbb{P}^1 \setminus \{z_1, \dots, z_n\}$ be an n -punctured Riemann sphere, $n > 3$, normalized by $z_{n-2} = 0, z_{n-1} = 1$ and $z_n = \infty$, and let $J: \mathbb{H} \rightarrow X$ be the uniformization map — a complex-analytic covering of X by the upper half-plane \mathbb{H} . The Schwarzian derivative $S(J^{-1})$ of the inverse map J^{-1} is the following rational function on \mathbb{P}^1

$$S(J^{-1})(z) = \sum_{i=1}^{n-3} \left(\frac{1/2}{(z - z_i)^2} + \frac{c_i}{z - z_i} \right) + \frac{2 - n}{2z(z - 1)}.$$

The coefficients c_i , $i = 1, \dots, n-3$, are smooth functions on the space of punctures

$$\mathcal{Z}_n := \{(z_1, \dots, z_{n-3}) \in \mathbb{C}^{n-3} \mid z_i \neq z_j \text{ for } i \neq j \text{ and } z_i \neq 0, 1\}$$

and are called accessory parameters of the Fuchsian uniformization of the n -punctured Riemann sphere. We proved in [48, 49] that, in accordance with a conjecture of A. Polyakov [33], there exists a real-valued smooth function S on \mathcal{Z}_n such that

$$c_i = -\frac{1}{2\pi} \frac{\partial S}{\partial z_i}, \quad i = 1, \dots, n-3.$$

The function S is the critical value of the action functional for the *Liouville theory* — *the two-dimensional quantum gravity*, and the formulas for accessory parameters follow from the *semi-classical Ward identity for the 1-point correlation function with holomorphic component of the stress-energy tensor* (see [41] for the details).

This comparison shows similarity between quantum theory of free bosons on \mathbb{P}^1 in the presence of the contour C and quantum Liouville theory. Namely, 1-point Ward identities for both theories imply that the logarithm of the τ -function and the critical value of the Liouville action (the logarithm of the semi-classical approximation to the partition function) are generating functions for the harmonic moments of interior and accessory parameters correspondingly.

Remark 3.5. In the case when the punctures z_1, \dots, z_{n-3} are real, the Riemann surface X from the previous remark possesses an anti-holomorphic involution $z \mapsto \bar{z}$. It is a classical result (see, e.g., [19]) that for this case the map J^{-1} is given by the conformal mapping of the upper half-plane \mathbb{H} (or the unit disk D) onto a circular n -sided polygon with zero angles, inscribed into S^1 . Under this map the marked points $z_1, \dots, z_{n-3}, 0, 1, \infty$ on the boundary $\mathbb{R} \cup \{\infty\}$ of \mathbb{H} (or corresponding n marked points on S^1) are mapped onto the vertices of the polygon and this map is unique if the last three vertices are normalized as $-1, -i, 1$. The harmonic moments of the boundary of the polygon depend on $n-3$ real parameters z_1, \dots, z_{n-3} . It would be instructive to express accessory parameters through harmonic moments and compare results [48, 49] with the Corollary 3.3 directly. In order to get the analog of the Corollary 3.3 for this case, one needs a generalization to the case of piece-wise smooth contours like boundaries of circular polygons. We do not address this interesting question here.

Remark 3.6. As it was pointed out in the Remark 1.1 there is an inclusion of $\text{Diff}_+(S^1)/\text{Möb}(S^1)$ into the Bers universal Teichmüller space $T(1)$, which contains all Teichmüller spaces $T_{g,n}$ of Riemann surfaces of type (g, n) as complex submanifolds. Similar to the smooth case, there is a correspondence $\gamma \mapsto C$, where γ is a quasi-conformal homeomorphism of S^1 and C is a quasi-circle — an image of S^1 under quasi-conformal homeomorphism of the complex plane \mathbb{C} which is conformal outside S^1 . Extension of the above formulation from smooth contours to

quasi-circles would naturally allow to consider Riemann surfaces of type (g, n) by the same method. We do not address this important question here.

For completeness, let us show how to determine conformal map G ‘explicitly’ [24, 45] from the Ward identity proved above. As it follows from the proof of the Krichever’s lemma and equation

$$S_-(z) = \frac{\partial\Phi(z)}{\partial z} - \frac{t_0}{z},$$

$$\frac{G'(z)}{G(z)} = -\frac{\partial S_-(z)}{\partial t_0} = -\frac{\partial^2\Phi(z)}{\partial t_0\partial z} + \frac{1}{z}.$$

Therefore

$$\frac{G'(z)}{G(z)} = \frac{1}{z} + \sum_{n=1}^{\infty} z^{-n-1} \frac{\partial^2 \log \tau}{\partial t_0 \partial t_n}$$

and integrating

$$\log G(z) = \log b_{-1} + \log z - \sum_{n=1}^{\infty} \frac{z^{-n}}{n} \frac{\partial^2 \log \tau}{\partial t_0 \partial t_n}$$

It is also possible [45] to express the Robin’s constant $\log b_{-1} = V$ (see Section 1.2) through $\log \tau$. Namely, consider

$$\frac{\partial^2 \log \tau}{\partial t_0^2} = \frac{\partial v_0}{\partial t_0}.$$

Using the definition of v_0 and the calculus formula, we get

$$\frac{\partial v_0}{\partial t_0} = \frac{1}{2\pi i} \int_C \log |z|^2 \frac{dG}{G} = -2 \log b_{-1},$$

as can be readily shown by integration by parts. Therefore, one gets the result in [24, 45].

COROLLARY 3.7. *The conformal map G is given by the ‘explicit formula’*

$$\log G(z) = \log z - \frac{1}{2} \frac{\partial^2 \log \tau}{\partial t_0^2} - \sum_{n=1}^{\infty} \frac{z^{-n}}{n} \frac{\partial^2 \log \tau}{\partial t_0 \partial t_n}.$$

Remark 3.8. In complex analysis there is the following relation between moments M_n of the equilibrium distribution for the domain Ω and the conformal mapping G of the exterior domain $\mathbb{P}^1 \setminus \Omega$ (see, e.g., [18])

$$\frac{G'(z)}{G(z)} = \sum_{n=0}^{\infty} M_n z^{-n-1}.$$

As it follows for the above formula for $G'(z)/G(z)$,

$$M_n = \frac{\partial^2 \log \tau}{\partial t_n \partial t_0}, \quad n \in \mathbb{N},$$

so that function $v_0 = \partial \log \tau / \partial t_0$ on \mathcal{C} is a generating function for the moments M_n .

Next, consider the normalized reduced 2-point current correlation function

$$\langle \langle j(z)j(w) \rangle \rangle := \frac{\langle j(z)j(w) \rangle}{\langle \mathbf{1} \rangle_{\mathcal{C}}} - \langle \langle j(z) \rangle \rangle \langle \langle j(w) \rangle \rangle.$$

Using the same arguments as in the proof of Proposition 2.1: the change of variables $\varphi = \Phi^{(\varepsilon)} + \tilde{\varphi}$ and passage to the limit $\varepsilon \rightarrow 0$, as well as the standard fact that $\langle \tilde{j}(z) \rangle_0 = 0$ we obtain, as in Section 4.1 of Part 1,

$$\langle j(z)j(w) \rangle = \langle \tilde{j}(z)\tilde{j}(w) \rangle_0 + \langle \mathbf{1} \rangle_{\mathcal{C}} \frac{\partial \Phi(z)}{\partial z} \frac{\partial \Phi(w)}{\partial w} dz \otimes dw.$$

As the result

$$\begin{aligned} \langle \langle j(z)j(w) \rangle \rangle &= \langle \langle \tilde{j}(z)\tilde{j}(w) \rangle \rangle_0 = \pi \frac{\partial^2 G(z, w)}{\partial z \partial w} dz \otimes dw \\ &= -\frac{dz \otimes dw}{(z - w)^2}, \end{aligned}$$

where we used a simple expression for the Schiffer kernel S on \mathbb{P}^1 from Section 1.3.

Our goal is to compare this correlation function with the corresponding reduced normalized 2-point correlation function for free bosons on $\mathbb{P}^1 \setminus \Omega$ with the Dirichlet boundary condition. Since $\langle j(z) \rangle_{\text{DBC}} = 0$ we have

$$\langle \langle j(z)j(w) \rangle \rangle_{\text{DBC}} := \frac{\langle j(z)j(w) \rangle_{\text{DBC}}}{\langle \mathbf{1} \rangle_{\text{DBC}}}.$$

Arguing as in the Section 4.1 of Part 1 and using results of Section 1.3, we get

$$\langle \langle j(z)j(w) \rangle \rangle_{\text{DBC}} = \pi \frac{\partial^2 G_{\text{DBC}}(z, w)}{\partial z \partial w} dz \otimes dw = -\frac{G'(z)G'(w)}{(G(z) - G(w))^2} dz \otimes dw.$$

On the other hand, consider $\partial^2 \log \tau / \partial t_m \partial t_n$, $m, n > 0$. Using the Ward identity for the 1-point correlation function, holomorphy of $\partial \Phi / \partial z$ on $\mathbb{P}^1 \setminus \Omega$ and the calculus formula, we get

$$\frac{\partial^2 \log \tau}{\partial t_m \partial t_n} = -\frac{1}{2\pi i} \frac{\partial}{\partial t_m} \int_{\mathcal{C}} \frac{\partial \Phi(z)}{\partial z} z^n dz = -\frac{1}{2\pi i} \int_{\mathcal{C}} \frac{\partial^2 \Phi(z)}{\partial t_m \partial z} z^n dz.$$

Using $\partial \Phi(z) / \partial z = S_-(z) + t_0/z$ and the formula

$$\dot{\omega}_{\mathcal{C}}^{(m)} = dz^m - \frac{\partial S_-}{\partial t_m} dz, \quad z \in \mathbb{P}^1 \setminus \Omega$$

(see the proof of part (i) of the Krichever’s lemma), we have

$$\frac{\partial^2 \log \tau}{\partial t_m \partial t_n} = \frac{1}{2\pi i} \int_C w^n (\dot{\omega}_C^{(m)} - dw^m) = \frac{1}{2\pi i} \int_C w^n d(F_m \circ G - w^m),$$

where in the last equation we have used Krichever’s lemma again. Next, it follows from the definition of Faber polynomials (see Section 1.3) that

$$\frac{G'(z)G'(w)}{(G(z) - G(w))^2} = \sum_{m=0}^{\infty} \frac{dF_m(G(w))}{dw} z^{-m-1},$$

where $|G(z)| > |G(w)|$. From here we get

$$\frac{dF_m(G(w))}{dw} = \frac{1}{2\pi i} \int_{C_w} \frac{G'(z)G'(w)}{(G(z) - G(w))^2} z^m dz,$$

where contour C_w is such that its image under the map G contains the circle of radius $|G(w)|$ inside. Similarly, from the expansion

$$\frac{1}{(z - w)^2} = \sum_{m=0}^{\infty} m w^{m-1} z^{-m-1},$$

where $|z| > |w|$, we get

$$m w^{m-1} = \frac{1}{2\pi i} \int_{|z|=R} \frac{1}{(z - w)^2} z^m dz,$$

where $|R| > |w|$. Thus we obtain

$$\frac{dF_m(G(w))}{dw} - m w^{m-1} = \frac{1}{2\pi i} \int_C \left(\frac{G'(z)G'(w)}{(G(z) - G(w))^2} - \frac{1}{(z - w)^2} \right) z^m dz,$$

where we moved the contour of integration from the neighborhood of ∞ to C because the integrand is regular for all $z, w \in \mathbb{C} \setminus \Omega$. Thus, we finally get

$$\begin{aligned} \frac{\partial^2 \log \tau}{\partial t_m \partial t_n} &= \frac{1}{(2\pi i)^2} \int_C \int_C \left(\frac{G'(z)G'(w)}{(G(z) - G(w))^2} - \frac{1}{(z - w)^2} \right) z^m w^n dz dw \\ &= \frac{1}{(2\pi i)^2} \int_C \int_C z^m w^n (\langle J(z)J(w) \rangle - \langle J(z)J(w) \rangle_{\text{DBC}}) \end{aligned}$$

— the Ward identity for the difference between normalized reduced 2-point correlation functions of holomorphic current components for free bosons on \mathbb{P}^1 parameterized by C and free bosons on $\mathbb{P}^1 \setminus \Omega$ with the Dirichlet boundary condition.

Next, consider normalized reduced 2-point correlation function of holomorphic and anti-holomorphic current components

$$\langle J(z)\bar{J}(w) \rangle := \frac{\langle J(z)\bar{J}(w) \rangle}{\langle \mathbf{1} \rangle_C} - \langle J(z) \rangle \langle \bar{J}(w) \rangle.$$

As before, we get

$$\langle\langle J(z)\bar{j}(w)\rangle\rangle = \langle J(z)\bar{j}(w)\rangle_0 = \pi \frac{\partial^2 G(z, w)}{\partial z \partial \bar{w}} = 0,$$

since the Bergman reproducing kernel on \mathbb{P}^1 is 0 (see Section 1.2). Similarly, arguing as in Section 4.1 of Part 1 and using results in Section 1.2 on the Bergman kernel on $\mathbb{P}^1 \setminus \Omega$, we get

$$\langle\langle J(z)\bar{j}(w)\rangle\rangle_{\text{DBC}} = \pi \frac{\partial^2 G_{\text{DBC}}(z, w)}{\partial z \partial \bar{w}} dz \otimes d\bar{w} = -\frac{G'(z)\overline{G'(w)}}{(1 - G(z)\overline{G(w)})^2} dz \otimes d\bar{w}.$$

Computation of $\partial^2 \log \tau / \partial t_m \partial \bar{t}_n$, $m, n > 0$ is also similar to the one done before. Namely, since the vector field $\partial / \partial \bar{t}_n$ corresponds to the meromorphic $(1, 0)$ -form $\dot{\omega}_C^{(\bar{n})}$ on \mathbb{P}^1_C which coincides with $d(F_n \circ 1/\bar{G})$ on the domain $\mathbb{P}^1 \setminus \Omega$ (see the proof of Krichever's lemma in Section 1.3), we get

$$\begin{aligned} \frac{\partial^2 \log \tau}{\partial t_m \partial \bar{t}_n} &= -\frac{1}{2\pi i} \frac{\partial}{\partial \bar{t}_n} \int_C z^m \frac{\partial \Phi(z)}{\partial z} dz = -\frac{1}{2\pi i} \int_C z^m \frac{\partial S_-(z)}{\partial \bar{t}_n} dz \\ &= \frac{1}{2\pi i} \int_C z^m \dot{\omega}_C^{(\bar{n})} = -\frac{1}{2\pi i} \int_C z^m d(F_n \circ 1/\bar{G}). \end{aligned}$$

Using the representation

$$-\frac{G'(z)\overline{G'(w)}}{(1 - G(z)\overline{G(w)})^2} = \sum_{n=0}^{\infty} \frac{dF_n(1/\bar{G}(z))}{dz} \bar{w}^{-n-1},$$

which follows from the definition of Faber polynomials, we get by the Cauchy integral formula

$$\frac{dF_n(1/\bar{G}(z))}{dz} = \frac{1}{2\pi i} \int_{C_+} \frac{G'(z)\overline{G'(w)}}{(1 - G(z)\overline{G(w)})^2} \bar{w}^n d\bar{w},$$

where C_+ is an arbitrary contour around 0 that contains the contour C inside. As a result we finally obtain

$$\begin{aligned} \frac{\partial^2 \log \tau}{\partial t_m \partial \bar{t}_n} &= -\frac{1}{(2\pi i)^2} \int_{C_+} \int_{C_+} \frac{G'(z)\overline{G'(w)}}{(1 - G(z)\overline{G(w)})^2} z^m \bar{w}^n dz d\bar{w} \\ &= -\frac{1}{(2\pi i)^2} \int_{C_+} \int_{C_+} z^m \bar{w}^n (\langle\langle J(z)\bar{j}(w)\rangle\rangle) - \langle\langle J(z)\bar{j}(w)\rangle\rangle_{\text{DBC}} \end{aligned}$$

— the Ward identity for the 2-point correlation function with holomorphic and anti-holomorphic current components that computes the difference between correlation functions of free bosons on \mathbb{P}^1 parameterized by C and of free bosons on $\mathbb{P}^1 \setminus \Omega$ with the Dirichlet boundary condition. We summarize these results as the following statement (cf. [24]).

THEOREM 3.9. *Normalized reduced 2-point current correlation functions for free bosons on \mathbb{P}^1 parameterized by $C \in \tilde{\mathcal{C}}_a$ for every $a > 0$ satisfy the Ward identities, given by the following Laurent series expansions at $z = w = \infty$*

$$\begin{aligned} & \langle\langle j(z)j(w) \rangle\rangle - \langle\langle j(z)j(w) \rangle\rangle_{\text{DBC}} \\ &= \left(\frac{G'(z)G'(w)}{(G(z) - G(w))^2} - \frac{1}{(z - w)^2} \right) dz \otimes dw \\ &= \sum_{m,n=1}^{\infty} z^{-m-1} w^{-n-1} \frac{\partial^2 \log \tau}{\partial t_m \partial t_n} dz \otimes dw \end{aligned}$$

and

$$\begin{aligned} & \langle\langle j(z)\bar{j}(w) \rangle\rangle - \langle\langle j(z)\bar{j}(w) \rangle\rangle_{\text{DBC}} \\ &= \frac{G'(z)\overline{G'(w)}}{(1 - G(z)\overline{G(w)})^2} dz \otimes d\bar{w} \\ &= \sum_{m,n=1}^{\infty} z^{-m-1} \bar{w}^{-n-1} \frac{\partial^2 \log \tau}{\partial t_m \partial \bar{t}_n} dz \otimes d\bar{w} \\ &= \mathbf{d}' \mathbf{d}'' \log \tau. \end{aligned}$$

All higher reduced multi-point current correlation functions vanish.

COROLLARY 3.10. *For every $a > 0$ the Hermitian metric H on $\tilde{\mathcal{C}}_a$ is Kähler with the Kähler potential $\log \tau$.*

Proof. Immediately follows from the definition of the metric H in Section 1.3. \square

Remark 3.11. This corollary should be compared with the result [48–50] that the critical value S of the Liouville action functional is a potential for the Weil–Petersson Kähler metric on the space of punctures \mathcal{Z}_n (and, therefore, on the Teichmüller space $T_{0,n}$)

$$g_{WP}^{ij} = -\frac{\partial^2 S}{\partial z_i \partial \bar{z}_j}, \quad i, j = 1, \dots, n - 3.$$

As was noted in the Remark 1.7, the metric H on $\tilde{\mathcal{C}}_a$ is a simplified analog of the Weil–Petersson metric on Teichmüller spaces $T_{g,n}$. Theorem 3.9 expresses it in terms of the 2-point current correlation function with holomorphic and anti-holomorphic components. Similarly, the Weil–Petersson metric on $T_{g,n}$ can be characterized as a semi-classical limit of a 2-point correlation function with holomorphic and anti-holomorphic components of the stress-energy tensor [41].

As before, it is possible to integrate explicitly the formulas in the Theorem 3.9, using the formula for $\log b_1$. We obtain the following result [24, 45].

COROLLARY 3.12.

$$\log \frac{G(z) - G(w)}{z - w} = -\frac{1}{2} \frac{\partial^2 \log \tau}{\partial t_0^2} + \sum_{m,n=1}^{\infty} \frac{z^{-m} w^{-n}}{mn} \frac{\partial^2 \log \tau}{\partial t_m \partial t_n}$$

and

$$\log \left(\frac{G(z) \overline{G(w)}}{G(z) \overline{G(w)} - 1} \right) = \sum_{m,n=1}^{\infty} \frac{z^{-m} \bar{w}^{-n}}{mn} \frac{\partial^2 \log \tau}{\partial t_m \partial \bar{t}_n}.$$

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