Mass, Kähler Manifolds, &

Symplectic Geometry

Claude LeBrun Stony Brook University

Differential Geometry Seminar, University of California at Irvine. April 2, 2019. Mass, Kähler Manifolds, and Symplectic Geometry

arXiv: 1810.11417 [math.DG]

Mass, Kähler Manifolds, and Symplectic Geometry

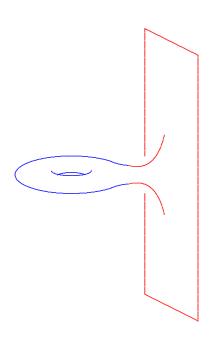
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Builds on previous paper

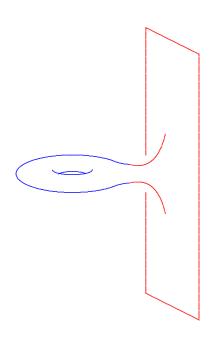
Mass in Kähler Geometry Comm. Math. Phys. 347 (2016) 621–653.

(Joint with Hans-Joachim Hein)

Definition. A complete, non-compact Riemannian n-manifold (M^n, g)

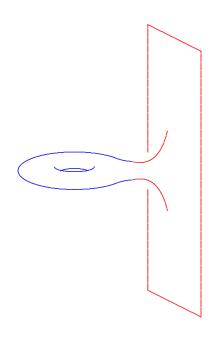


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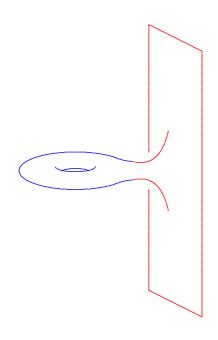
 $n \ge 3$

Definition. A complete, non-compact Riemannian n-manifold (M^n, g) is called asymptotically Euclidean



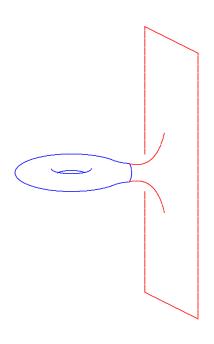
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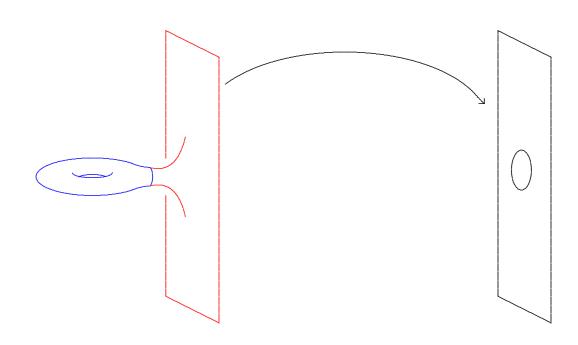
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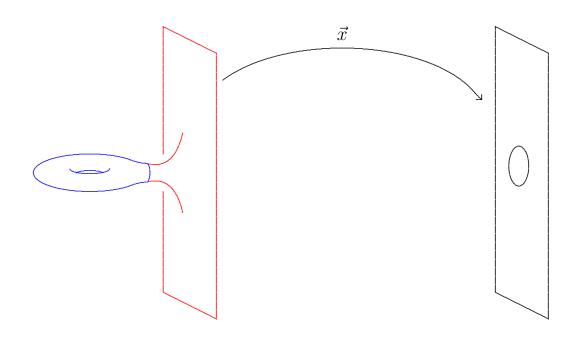


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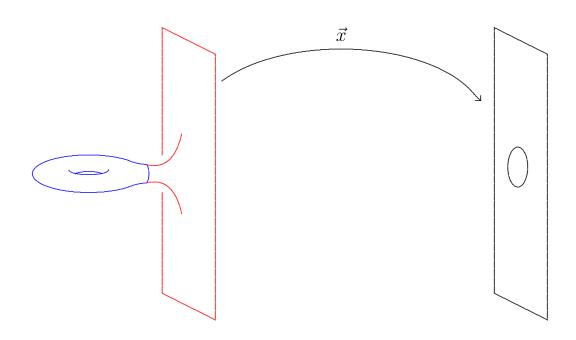
Definition. A complete, non-compact Riemannian n-manifold (M^n, g) is called asymptotically Euclidean (AE) if there is a compact set $K \subset M$



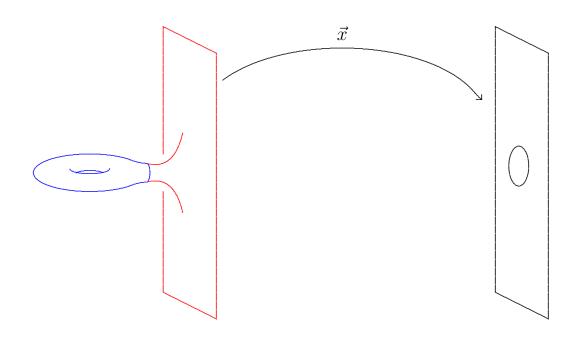




 $g_{jk} = \delta_{jk} + \text{terms that fall-off at infinity}$

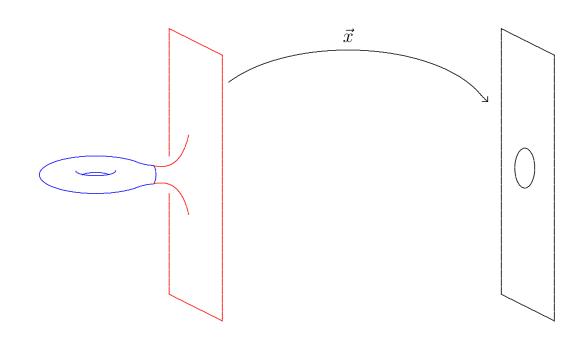


Weakest reasonable assumption:

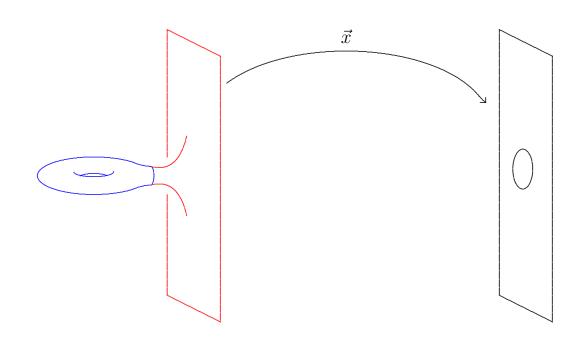


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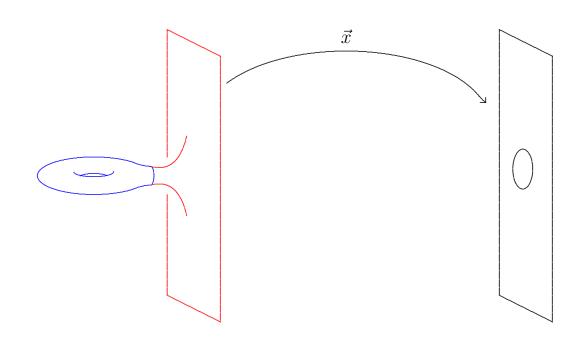
Chruściel-type fall-off:



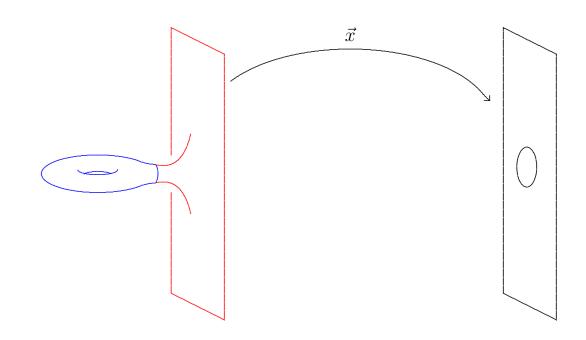
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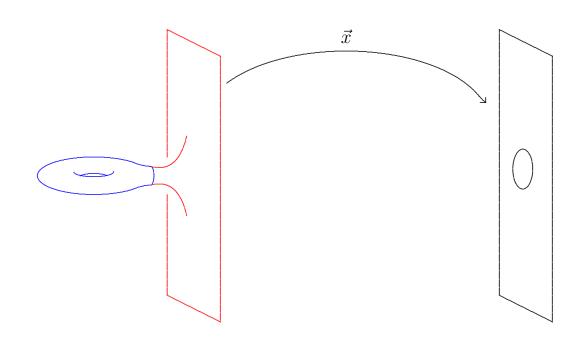


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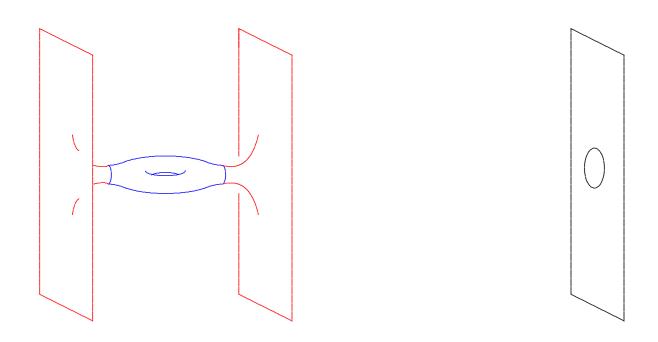


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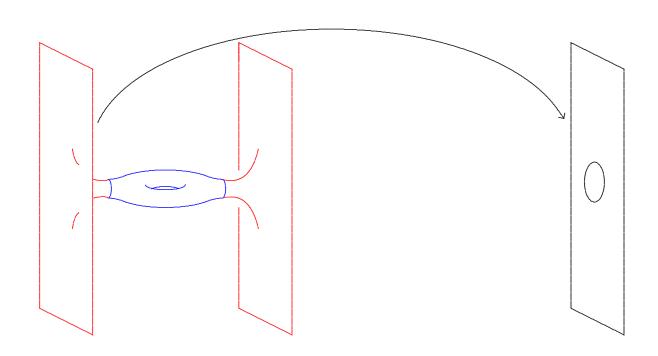
$$g_{jk,\ell} = O(|x|^{-\frac{n}{2} - \varepsilon}), \quad \text{scalar curvature } \in L^1$$



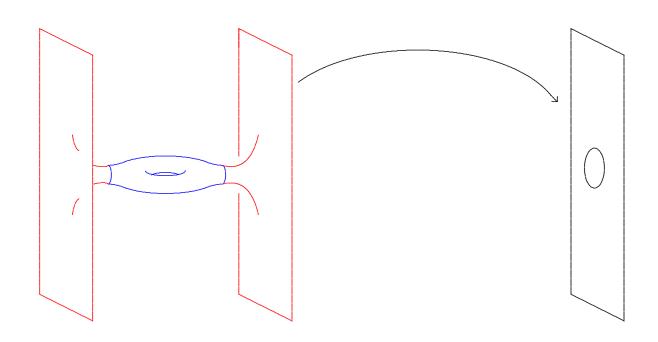
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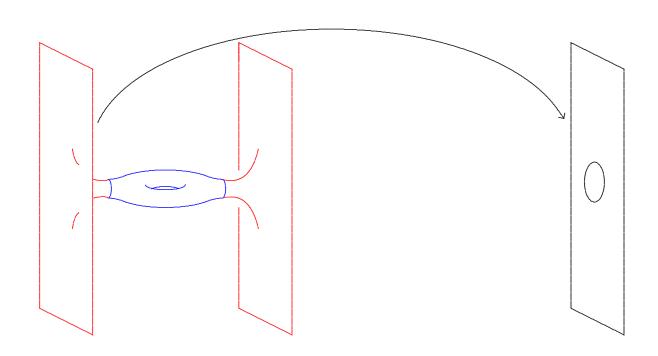
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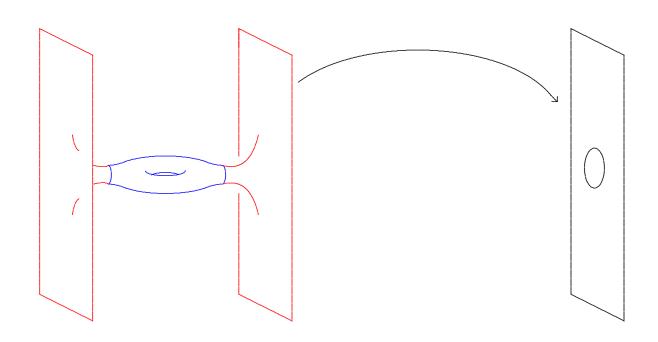
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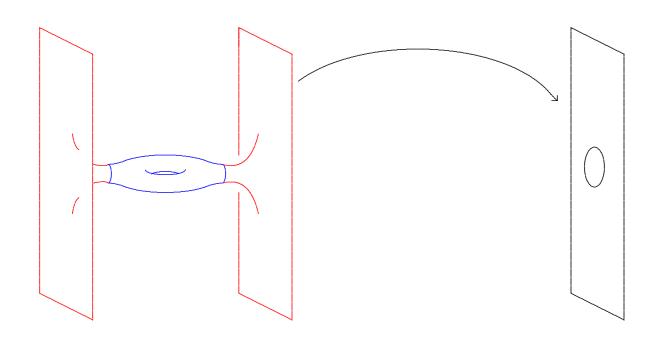
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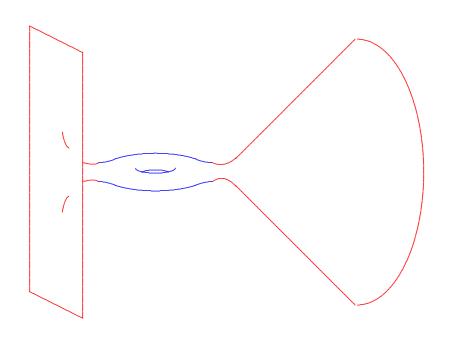


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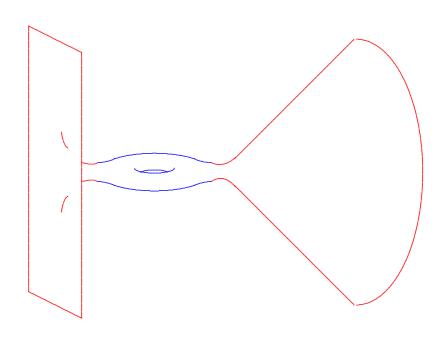


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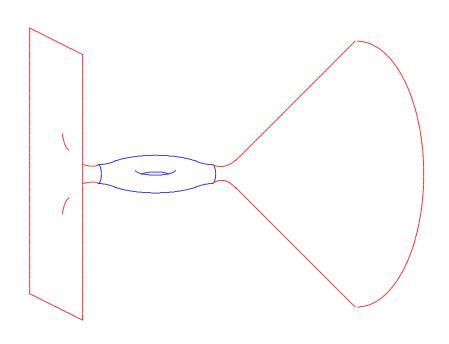
Definition. Complete, non-compact n-manifold (M^n, g) is asymptotically locally Euclidean



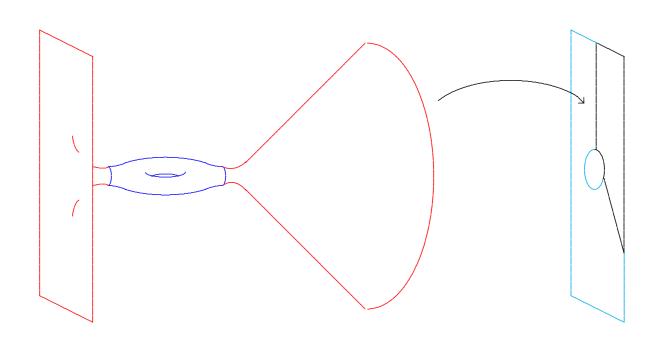
Definition. Complete, non-compact n-manifold (M^n, g) is asymptotically locally Euclidean (ALE)



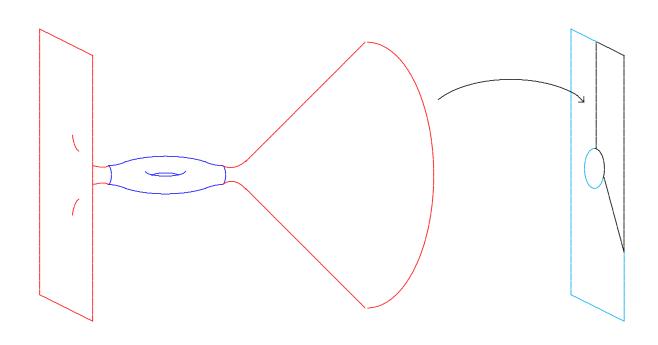
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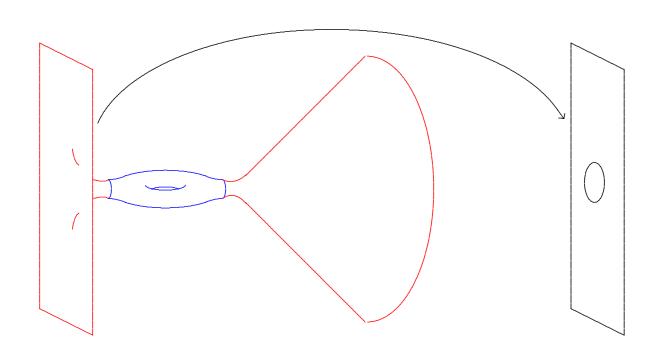
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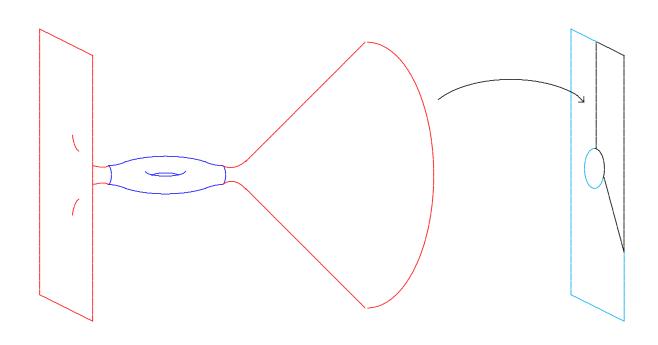
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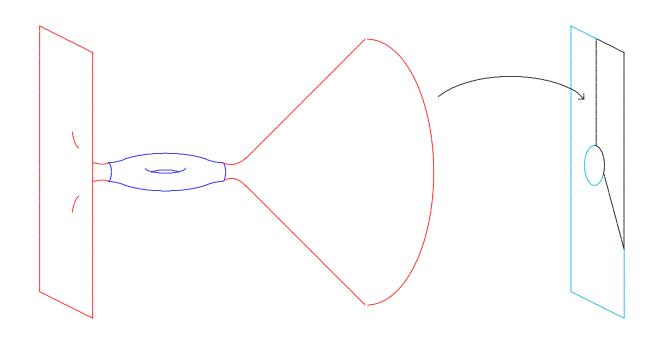
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$$g_{jk} = \delta_{jk} + O(|x|^{1 - \frac{n}{2} - \varepsilon})$$
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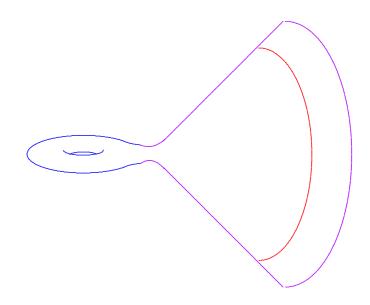
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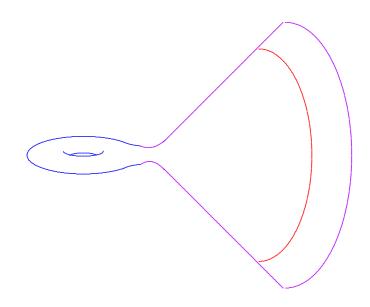
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Bartnik/Chruściel (1986): With weak fall-off conditions,

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Bartnik/Chruściel (1986): With weak fall-off conditions, the mass is well-defined & coordinate independent.

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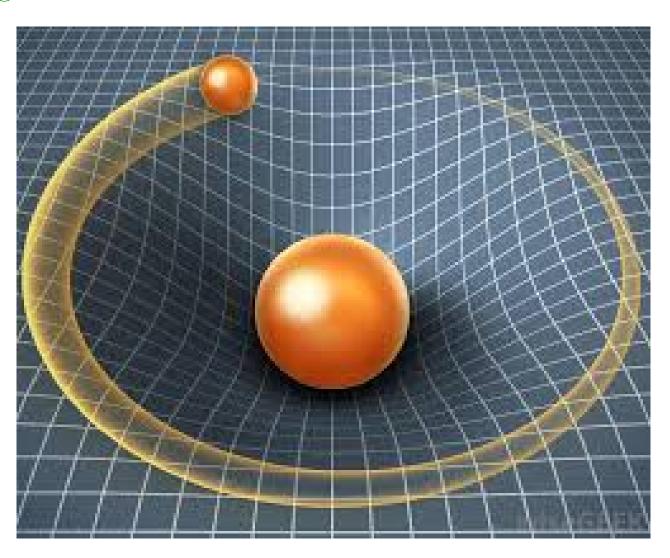
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When n = 3, ADM mass in general relativity.

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Reads off "apparent mass" from strength of the gravitational field far from an isolated source.

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$$g = -\left(1 - \frac{2m}{\varrho^{n-2}}\right)dt^2 + \left(1 - \frac{2m}{\varrho^{n-2}}\right)^{-1}d\varrho^2 + \varrho^2 h_{S^{n-1}}$$

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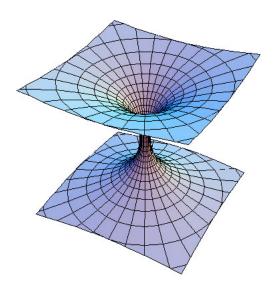
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Reads off "apparent mass" from strength of the gravitational field far from an isolated source.

In any dimension, reproduces "mass" of t=0 hypersurface in (n+1)-dimensional Schwarzschild

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Two such regions fit together to form the wormhole metric.



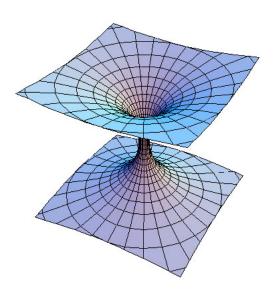
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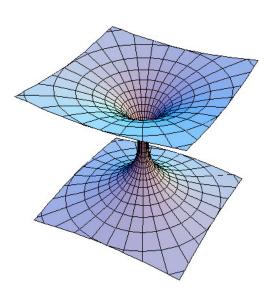
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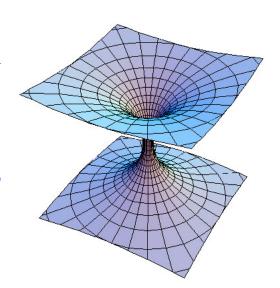
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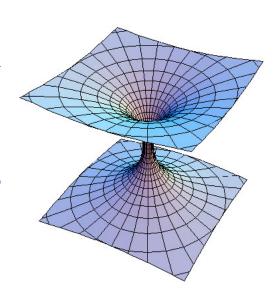
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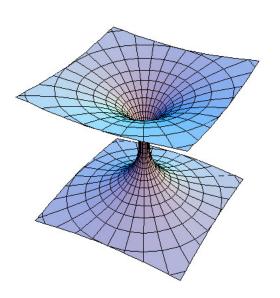
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$$g = \left(1 + \frac{m/2}{r^{n-2}}\right)^{4/(n-2)} \left[\sum (dx^j)^2\right]$$

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But notice that this crude model for the mass in particular assumes faster metric fall-off!

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Burns metric on \mathbb{C}^2

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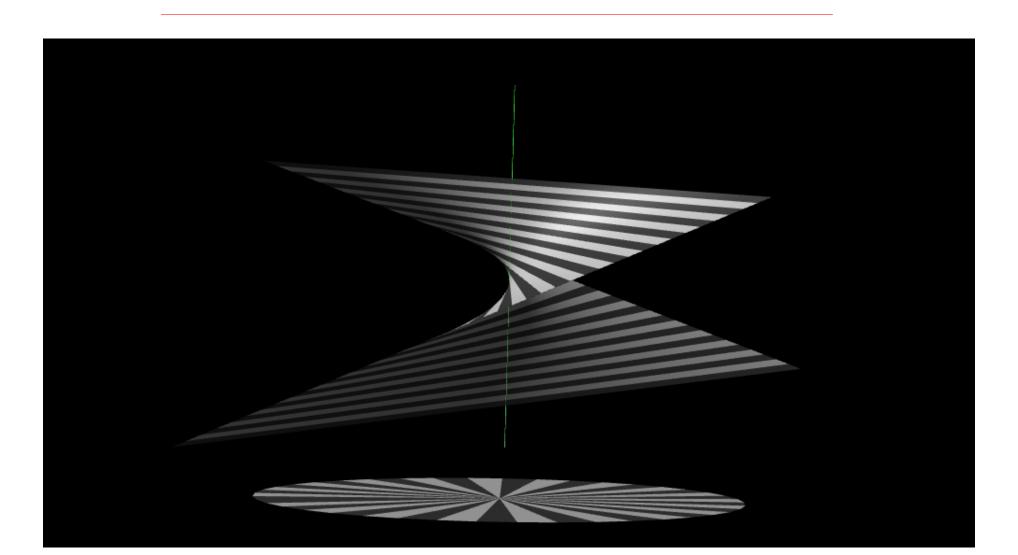
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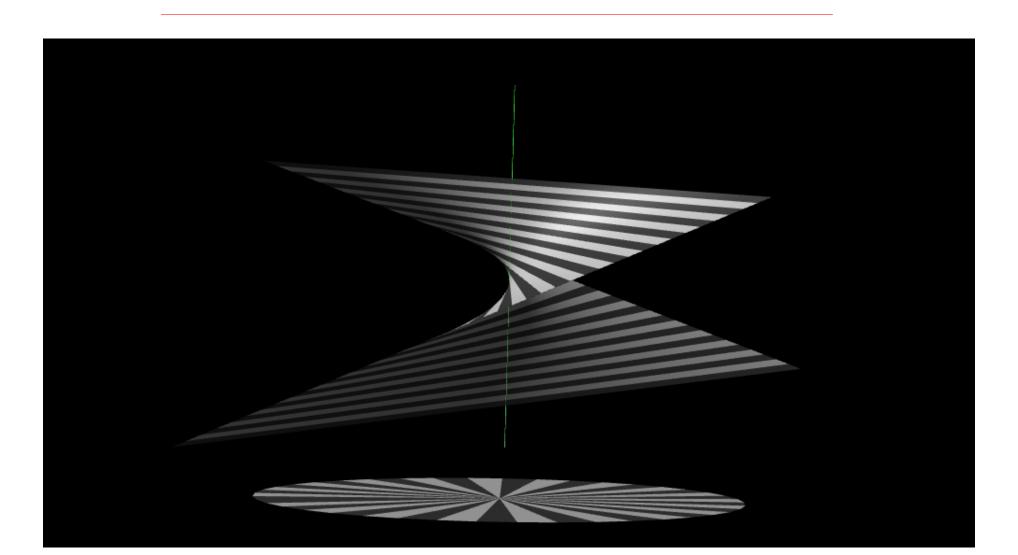
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Restrict Euclidean × round metric to
$$\begin{vmatrix} z_1 & z_2 \\ \zeta_1 & \zeta_2 \end{vmatrix} = 0.$$

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$$\begin{vmatrix} z_1 & z_2 \\ \zeta_1 & \zeta_2 \end{vmatrix} = 0.$$

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Reads off "apparent mass" from strength of the gravitational field far from an isolated source.

In any dimension, reproduces "mass" of t=0 hypersurface in (n+1)-dimensional Schwarzschild

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Any AE manifold

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unless $\varepsilon > \frac{1}{2}$, when Chruściel fall-off sufficed.

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This entails new proofs in real dimension four that are based on results in symplectic geometry.

This is rather unsatisfactory, since real dimension four is the setting for many of the most interesting applications!

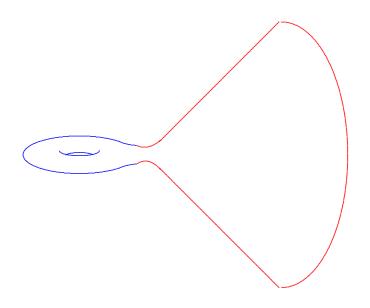
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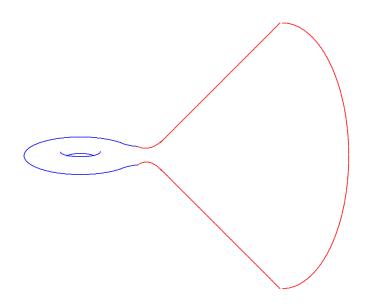
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In particular, Chruściel fall-off suffices to imply all the following results:

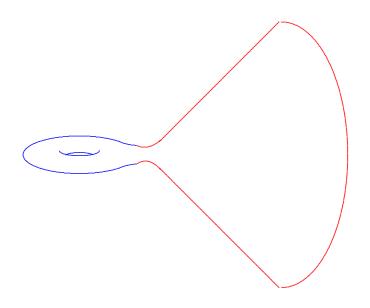
Lemma. Any ALE Kähler manifold has only one end.

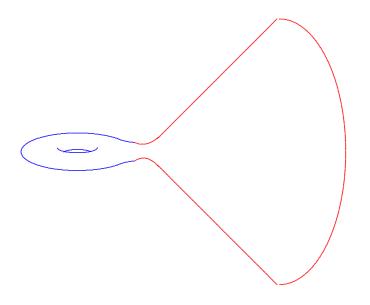




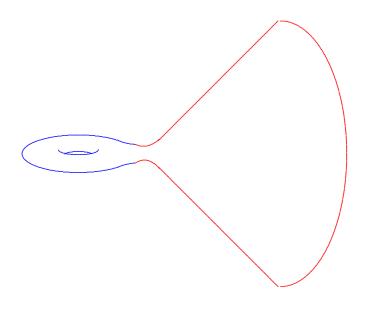
$$n = 2m \ge 4$$

Lemma. Any ALE Kähler manifold has only one end.



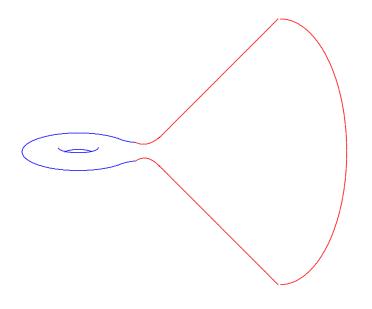


Upshot:



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Mass of an ALE Kähler manifold is unambiguous.



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Mass of an ALE Kähler manifold is unambiguous.

Does not depend on the choice of an end!

Theorem C.

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$$m(M,g) = +$$

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$$m(M,g) = + \frac{(m-1)!}{4(2m-1)\pi^m} \int_M s_g d\mu_g$$

$$m(M,g) = -\frac{\langle \mathbf{A}(c_1), [\omega]^{m-1} \rangle}{4(2m-1)\pi^m} \int_M s_g d\mu_g$$

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- $\clubsuit: H^2(M) \xrightarrow{\cong} H^2_c(M)$ inverse of natural map.

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New proof shows this follows from Chruściel fall-off.

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In particular, this shows mass is coordinate-invariant, without ever invoking Bartnik-Chruściel!

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Theorems A & B are corollaries concerning scalarflat Kähler metrics.

Another key consequence...

Theorem E (Penrose Inequality).

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Theorem E (Penrose Inequality). Let (M^{2m}, g, J) be an AE Kähler manifold with scalar curvature $s \geq 0$. Then (M, J) carries a canonical divisor D that is expressed as a sum $\sum_{j} \mathbf{n}_{j} D_{j}$ of compact complex hypersurfaces with positive integer coefficients, with the property that $\bigcup_{j} D_{j} \neq \emptyset$

$$m(M,g) \ge \sum Vol(D_j)$$

$$m(M,g) \ge \frac{(m-1)!}{(2m-1)\pi^{m-1}} \sum_{j=1}^{n} \operatorname{Vol}(D_j)$$

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This has an interesting corollary...

Theorem D (Positive Mass Theorem).

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Equality \Longrightarrow Ricci-flat because Ricci-form L^2 harmonic and $c_1=0$.

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Now use Bishop-Gromov inequality.

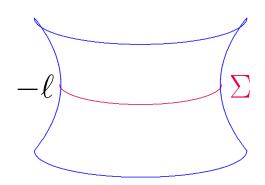
Some applications . . .

Theorem C. Any ALE Kähler manifold (M, g, J) of complex dimension m has mass given by

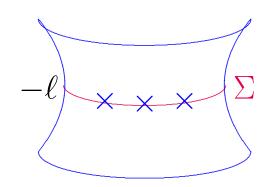
$$\mathbf{m}(\mathbf{M},g) = -\frac{\langle \mathbf{A}(\mathbf{c}_1), [\boldsymbol{\omega}]^{m-1} \rangle}{(2m-1)\pi^{m-1}} + \frac{(m-1)!}{4(2m-1)\pi^m} \int_{\mathbf{M}} \mathbf{s}_g d\mu_g$$

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$$-\ell-k$$
 $E_1 E_2 E_k$
 $\tilde{\Sigma}$

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L, Calderbank-Singer, Lock-Viaclovsky...

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That is, m(M, g, J) is completely determined by

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Theorem B. Let (M^4, g, J) be an ALE scalar-flat Kähler surface, and suppose that (M, J) is the minimal resolution of a surface singularity.

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Theorem B. Let (M^4, g, J) be an ALE scalarflat Kähler surface, and suppose that (M, J) is the minimal resolution of a surface singularity. Then $m(M, g) \leq 0$, **Theorem A.** The mass of an ALE scalar-flat Kähler manifold is a topological invariant.

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In particular, this shows mass is coordinate-invariant, without ever invoking Bartnik-Chruściel!

Theorems A & B are corollaries concerning scalarflat Kähler metrics. Theorem C. Any ALE Kähler manifold (M, g, J) of complex dimension m has mass given by

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Theorems A & B are corollaries concerning scalarflat Kähler metrics. But for such metrics, faster fall-off is guaranteed, so new proof is not actually needed!

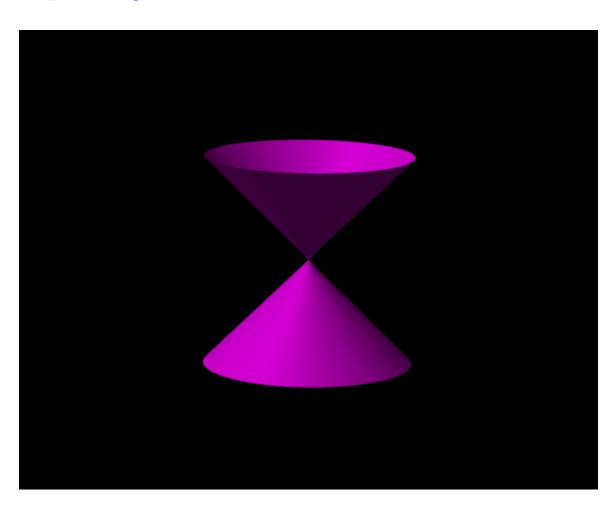
How does one prove main results?

That is, there exist asymptotic complex coordinates.

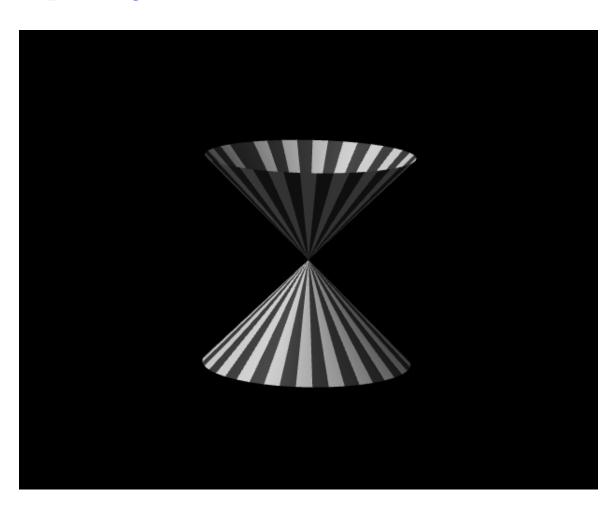
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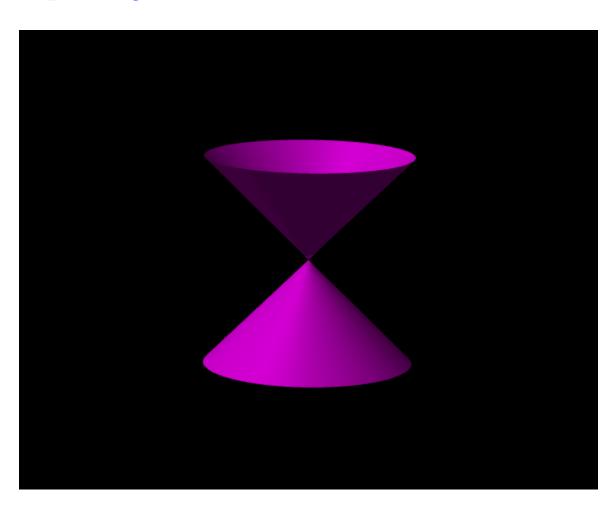
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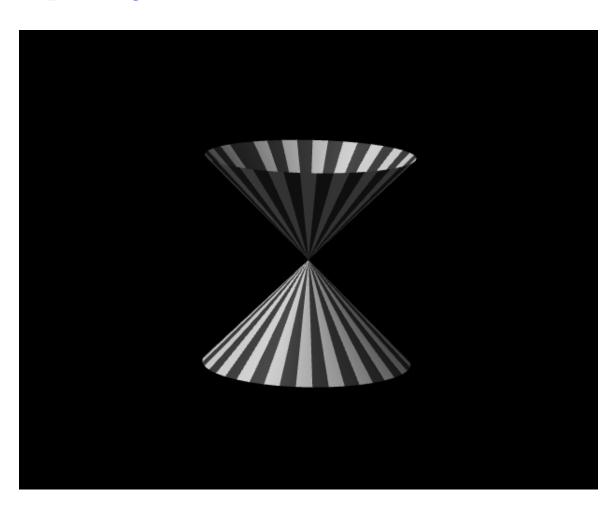
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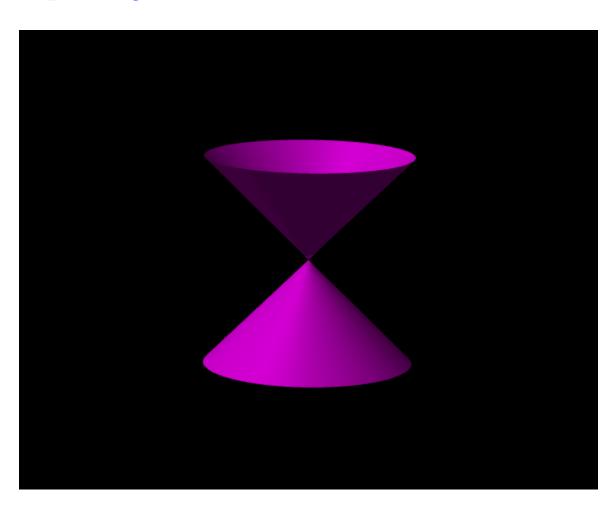
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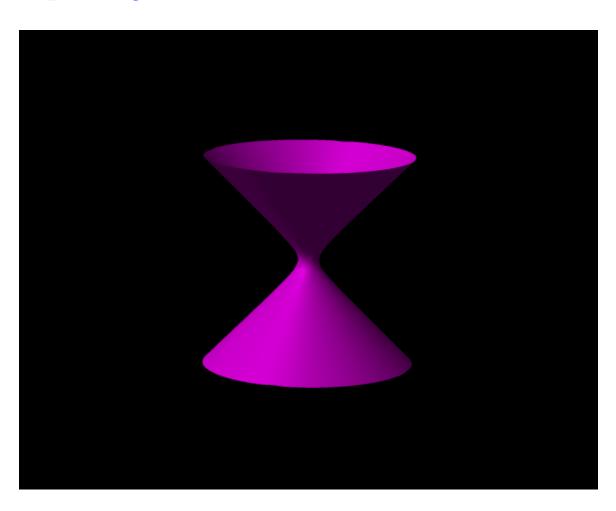
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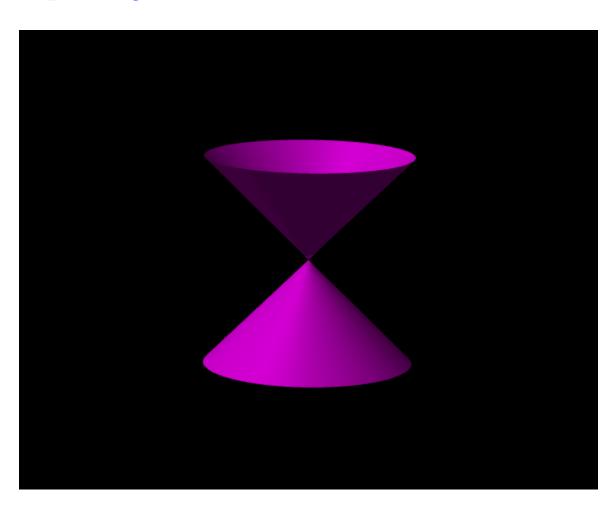
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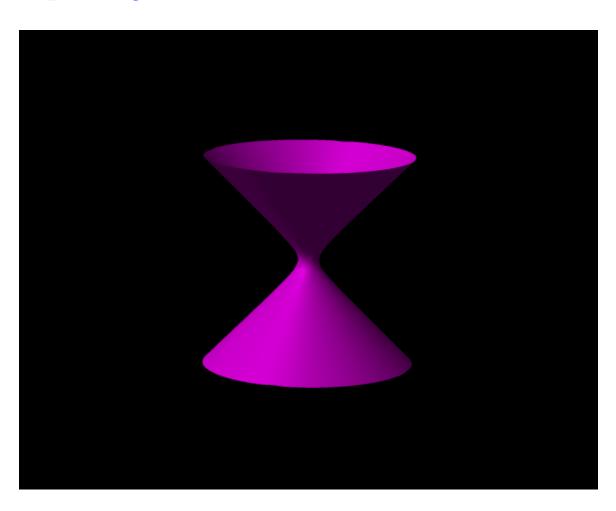
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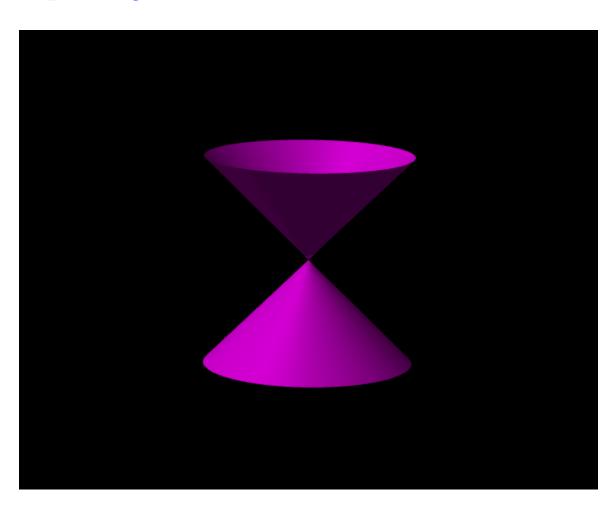
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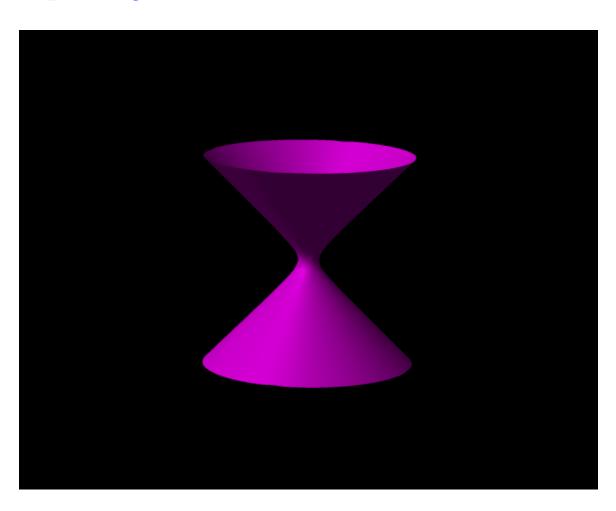
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Fortunately, however, the symplectic structure is always is standard at infinity!

Proposition.

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Proposition. Let (M^4, g, J) be an ALE Kähler surface, let $M_{\infty,i}$ be an end of M, let $\widetilde{M}_{\infty,i}$ be the universal cover of $M_{\infty,i}$, and let

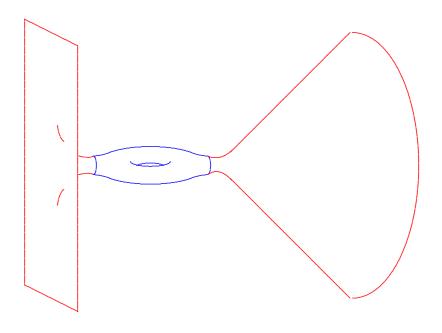
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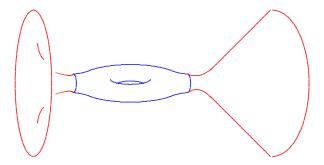
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Quantitative version of Moser stability argument...

Symplectically compactify any ALE Kähler surface:

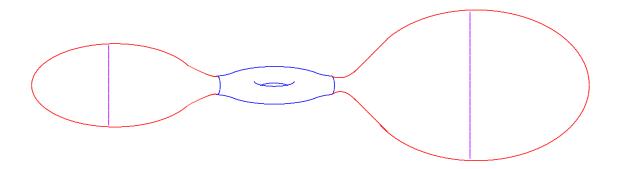


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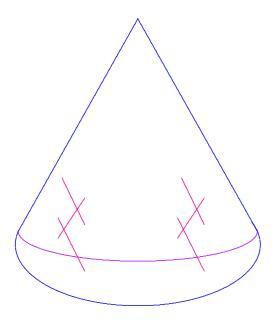


- Truncate at large radius;
- Glue in standard plugs.

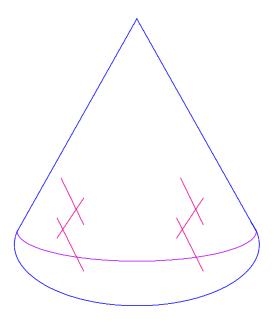
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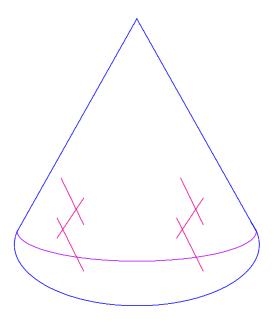
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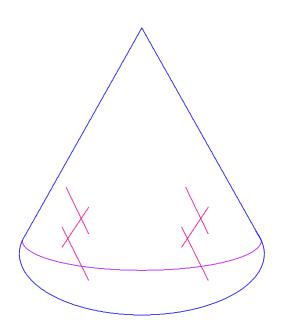
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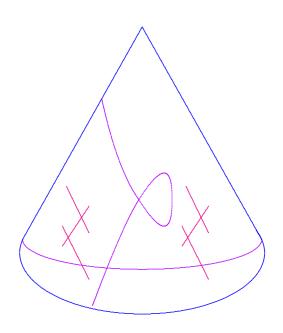
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for some, and hence any, ω -compatible almost-complex structure J.

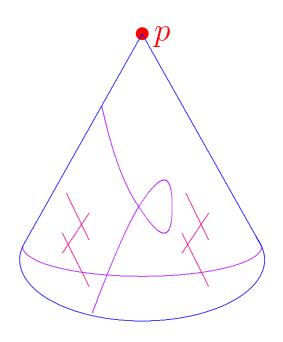
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Proof uses *J*-holomorphic curves.

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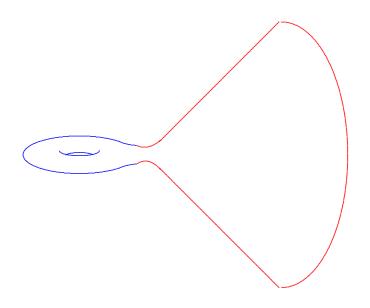
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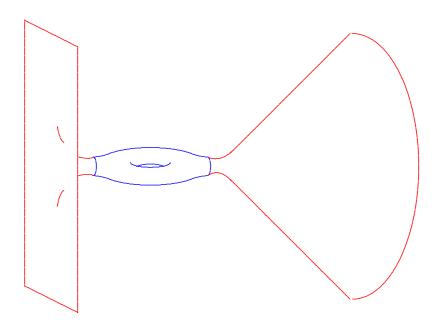
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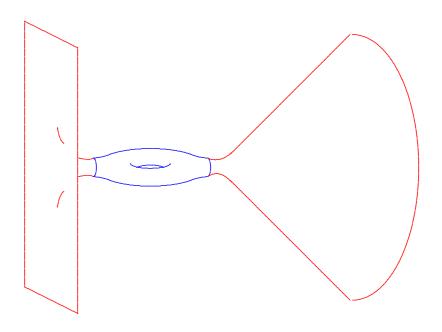
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Lemma. Any ALE Kähler manifold has only one end.

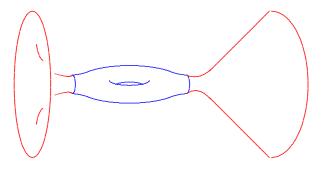


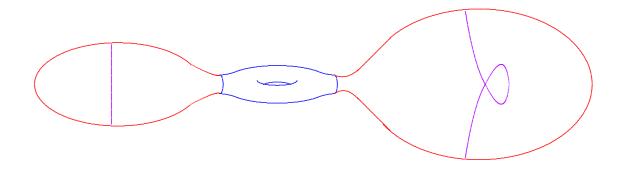


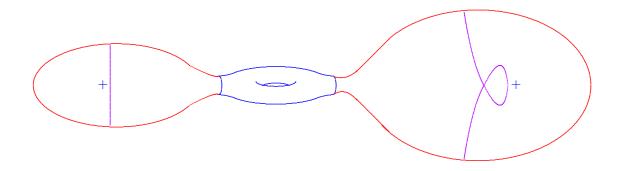
What if M^4 has more than one end?



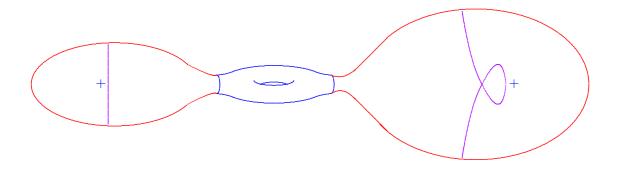
Compactify M as symplectic 4-manifold \widehat{M} .





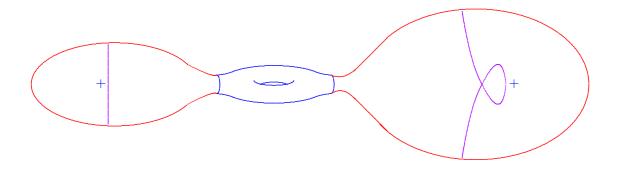


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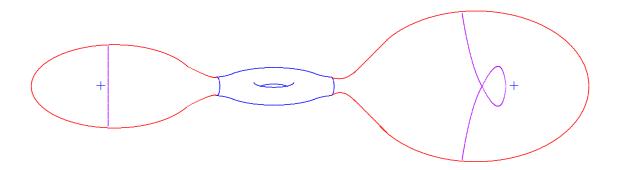
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McDuff $\Longrightarrow \widehat{M} \approx \text{rational complex surface.}$



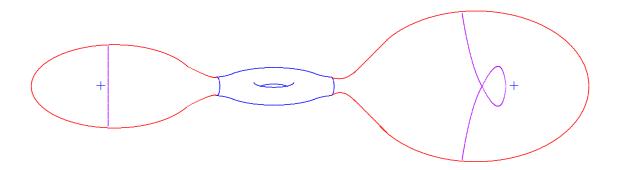
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 $McDuff \Longrightarrow intersection form (+-\cdots-).$



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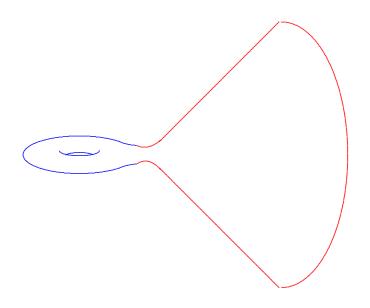


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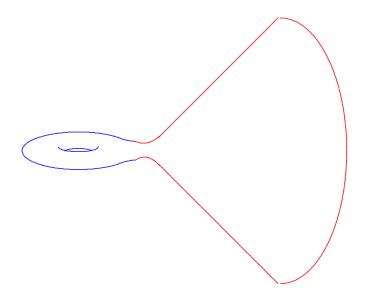
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Since each end contributes positive direction...

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[In higher dimensions, one similarly shows that (M, J) can be compactified as Kähler orbifold. The Hodge theorem on intersection form instead tells one that form on $H^{1,1}(\widehat{M}, \mathbb{R})$ is of type $(+-\cdots-)$.]

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The following result provides the key...

Proposition. Let g be a C^2 Kähler metric on $(\mathbb{R}^{2m} - \mathbf{D}^{2m})/\Gamma$, $m \geq 2$, where $\Gamma \subset \mathbf{SO}(2m)$ is some finite group that acts without fixed-points on S^{2m-1} .

$$g_{jk} - \delta_{jk} = C^1_{-(m-1+\varepsilon)}, \quad s \in L^1.$$

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Then there is a continuously differentiable 1-form θ on $(\mathbb{R}^{2m} - \mathbf{D}^{2m})/\Gamma$ such that

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Proposition. Let g be a C^2 Kähler metric on $(\mathbb{R}^{2m} - \mathbf{D}^{2m})/\Gamma$, $m \geq 2$, where $\Gamma \subset \mathbf{SO}(2m)$ is some finite group that acts without fixed-points on S^{2m-1} . In the given asymptotic coordinate system, suppose that g satisfies the Chuściel fall-off hypothesis

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where ρ is the Ricci form of g, and such that

$$\int_{S_{\varrho}/\Gamma} \left[g_{kj,k} - g_{kk,j} \right] \mathbf{n}^j d\mathfrak{a}_E =$$

$$\frac{2}{(m-1)!} \int_{S_{\varrho}/\Gamma} \theta \wedge \omega^{m-1} + O(\varrho^{-2\varepsilon}).$$

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Then direct computation shows that

$$[g_{kj,k} - g_{kk,j}] \mathbf{n}^j d\mathfrak{a}_E = \frac{2}{(m-1)!} \theta \wedge \omega^{m-1} + d\Omega + O(\varrho^{-2m-1-2\varepsilon})$$

First rotate coordinates so that $J \to J_0$ at ∞ , where J_0 standard complex structure on \mathbb{C}^m .

Now consider smooth trivialization of $\Lambda^{m,0}$ given by orthogonal projection of $dz^1 \wedge \cdots \wedge dz^m$.

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Integrating on S_{ρ}/Γ therefore yields:

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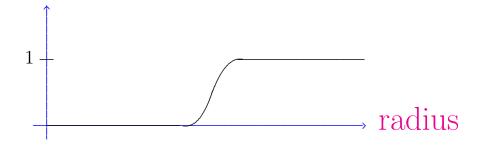
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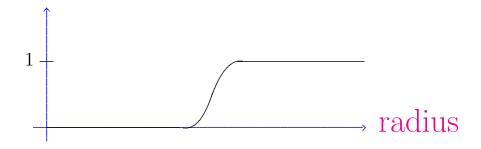
$$\int_{S_{\varrho}/\Gamma} \left[g_{kj,k} - g_{kk,j} \right] \mathbf{n}^{j} d\mathfrak{a}_{E} = \frac{2}{(m-1)!} \int_{S_{\varrho}/\Gamma} \theta \wedge \omega^{m-1} + O(\varrho^{-2\varepsilon})$$

We'll now deduce the mass formula...

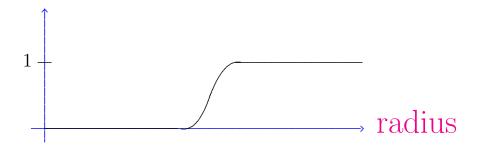
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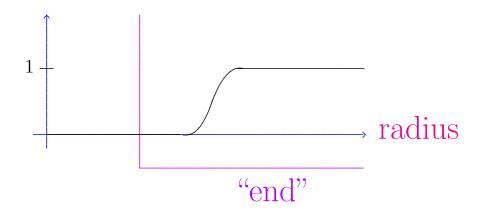
Let $f: M \to \mathbb{R}$ be smooth cut-off function: $\equiv 0$ away from end,



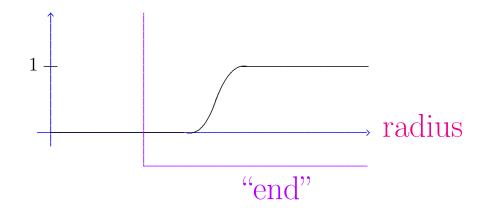
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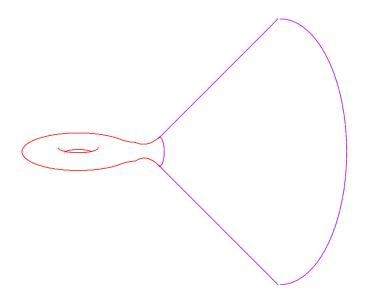


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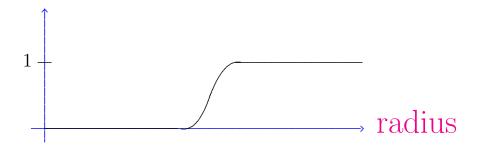


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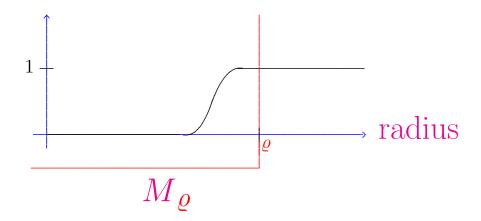




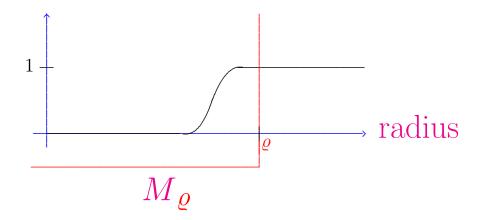
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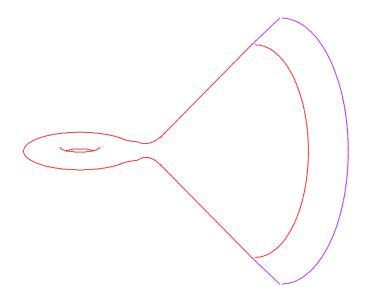


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$$\psi := \rho - d(f\theta)$$

Set

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Compactly supported, because $d\theta = \rho$ near infinity.

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where M_{ϱ} defined by radius $\leq \varrho$.

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$$\langle 2\pi \clubsuit(\mathbf{c}_1), [\omega]^{m-1} \rangle = \int_{M_{\varrho}} \rho \wedge \omega^{m-1} - \int_{M_{\varrho}} d(f\theta) \wedge \omega^{m-1}$$

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$$\langle 2\pi \clubsuit(c_1), [\omega]^{m-1} \rangle = \frac{(m-1)!}{2} \int_{M_{\varrho}} s \, d\mu - \int_{M_{\varrho}} d(f\theta \wedge \omega^{m-1})$$

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by Stokes' theorem.

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because there is only one end!

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$$\int_{S_{\varrho}/\Gamma}\theta\wedge\omega^{m-1}=-\langle 2\pi \clubsuit(c_1),[\omega]^{m-1}\rangle+\frac{(m-1)!}{2}\int_{M_{\varrho}}s\,d\mu$$

Let $f: M \to \mathbb{R}$ be smooth cut-off function:

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$$[\psi] = \clubsuit([\rho]) = 2\pi \clubsuit(c_1) \in H_c^2(M)$$

$$\frac{2}{(m-1)!} \int_{S_{\varrho}/\Gamma} \theta \wedge \omega^{m-1} = -\frac{4\pi \langle \clubsuit(c_1), [\omega]^{m-1} \rangle}{(m-1)!} + \int_{M_{\varrho}} s \, d\mu$$

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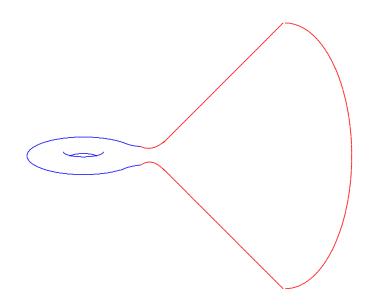
$$-\frac{4\pi\langle \mathbf{A}(\mathbf{c}_1), [\boldsymbol{\omega}]^{m-1}\rangle}{(m-1)!} + \int_{M_{\varrho}} s \, d\mu + O(\varrho^{-2\varepsilon})$$

Limit as $\varrho \to \infty$ now yields the mass formula.

$$\frac{4\pi^m(2m-1)}{(m-1)!} \mathbf{m}(M,g) = -\frac{4\pi}{(m-1)!} \langle \mathbf{A}(\mathbf{c}_1), [\omega]^{m-1} \rangle + \int_M \mathbf{s}_g d\mu_g$$

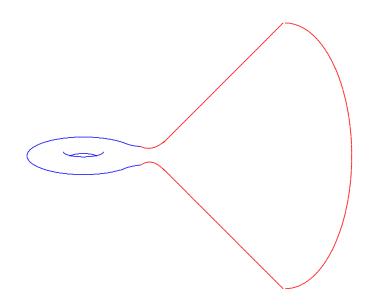
$$m(M,g) = -\frac{\langle \mathbf{A}(c_1), [\omega]^{m-1} \rangle}{(2m-1)\pi^{m-1}} + \frac{(m-1)!}{4(2m-1)\pi^m} \int_M s_g d\mu_g$$

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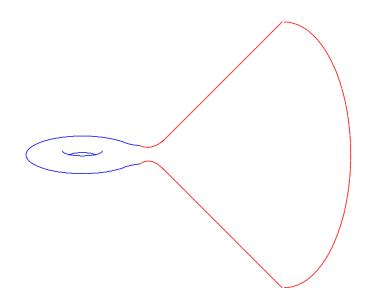
But the Penrose-type inequality is more subtle.

Theorem E (Penrose Inequality). Let (M^{2m}, g, J) be an AE Kähler manifold with scalar curvature $s \geq 0$. Then (M, J) carries a canonical divisor D that is expressed as a sum $\sum_{j} \mathbf{n}_{j} D_{j}$ of compact complex hypersurfaces with positive integer coefficients, with the property that $\bigcup_{j} D_{j} \neq \emptyset$ whenever $M \not\approx \mathbb{R}^{2m}$. In terms of this divisor, we then have

$$m(M,g) \ge \frac{(m-1)!}{(2m-1)\pi^{m-1}} \sum_{j=1}^{n} \operatorname{Vol}(D_j)$$

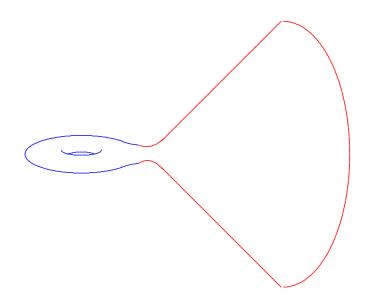
 $with = \iff (M, g, J) \text{ is scalar-flat K\"{a}hler.}$

$$\mathbf{m}(\mathbf{M},g) = \frac{\langle \mathbf{A}(-\mathbf{c}_1), [\boldsymbol{\omega}]^{m-1} \rangle}{(2m-1)\pi^{m-1}} + \frac{(m-1)!}{4(2m-1)\pi^m} \int_{\mathbf{M}} \mathbf{s}_g d\mu_g$$



m = 2:

$$m(M,g) = \frac{\langle \clubsuit(-c_1), [\omega] \rangle}{3\pi} + \frac{1}{12\pi} \int_M s_g d\mu_g$$



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Choose almost-complex structure to coincide with integrable J except in roughly conical asymptotic region and standard neighborhood of line at infinity.

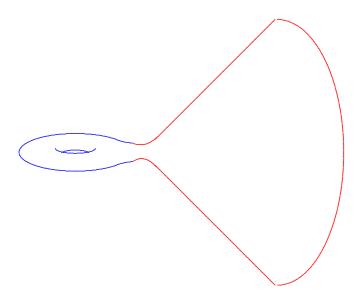
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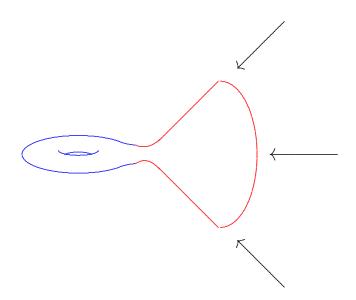
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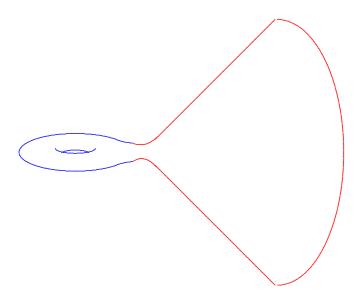
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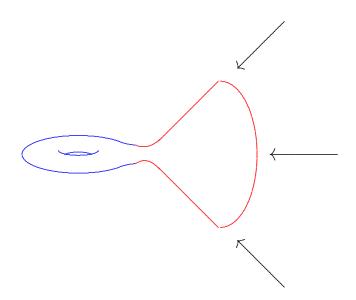
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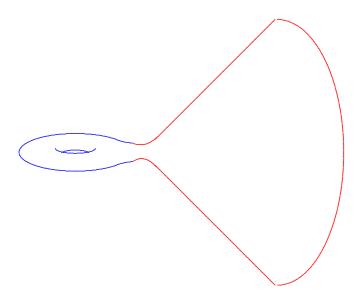
Technical challenge: Loss of control of derivatives!

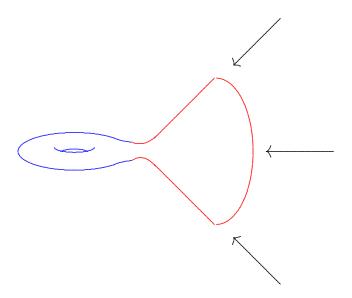












Robust under distortion of metric in outer region.

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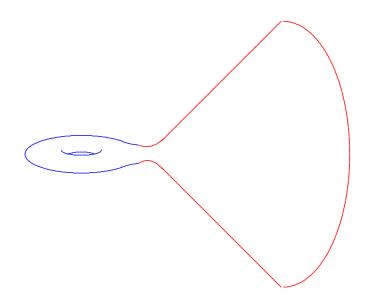
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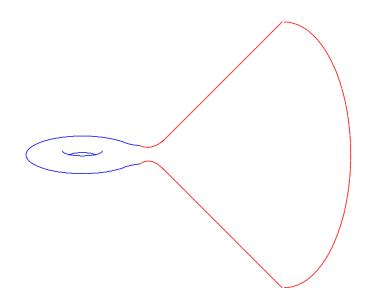
In (M, J), this gives desired Poincaré dual of $-c_1$.

m = 2:

$$m(M,g) = \frac{\langle \clubsuit(-c_1), [\omega] \rangle}{3\pi} + \frac{1}{12\pi} \int_M s_g d\mu_g$$

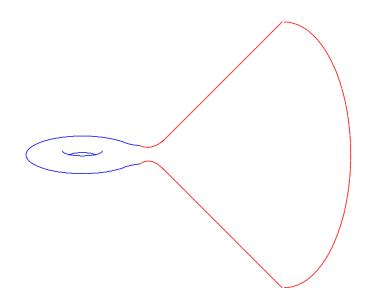


$$\mathbf{m}(\mathbf{M},g) = \frac{\langle \mathbf{A}(-\mathbf{c}_1), [\boldsymbol{\omega}]^{m-1} \rangle}{(2m-1)\pi^{m-1}} + \frac{(m-1)!}{4(2m-1)\pi^m} \int_{\mathbf{M}} \mathbf{s}_g d\mu_g$$



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