

Renormalization for irrationally indifferent fixed points of holomorphic functions

Mitsuhiro Shishikura

Kyoto University

Frontiers in Complex Dynamics
Celebrating John Milnor's 80th birthday
BIRS, Banff, Canada
February 23, 2011

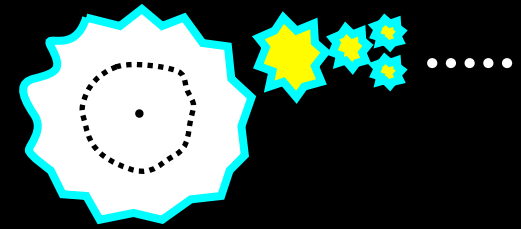
Frontiers in Complex Dynamics

use renormalization!

Frontiers in Complex Dynamics
include *boundaries* of Siegel disks

use renormalization!

Frontiers in Complex Dynamics
include *boundaries* of Siegel disks

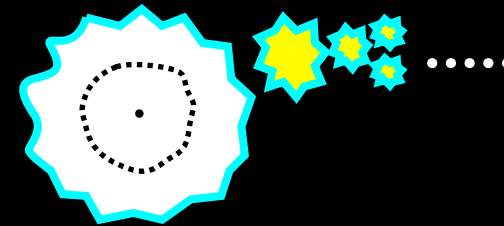


use renormalization!

Frontiers in Complex Dynamics

include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition



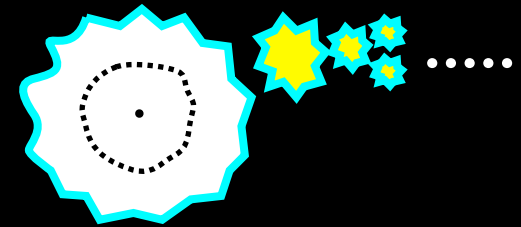
use renormalization!

Frontiers in Complex Dynamics

include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



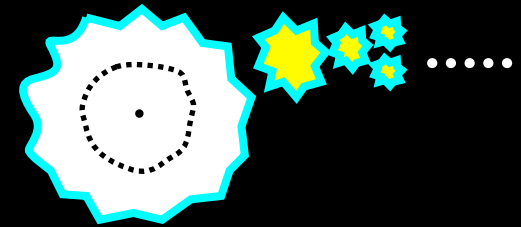
use renormalization!

Frontiers in Complex Dynamics

include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



Typical argument to show that some objects are continuous curves: for example, $J(z^2 + \varepsilon)$ is a Jordan curve

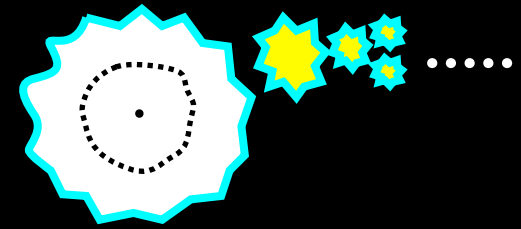
use renormalization!

Frontiers in Complex Dynamics

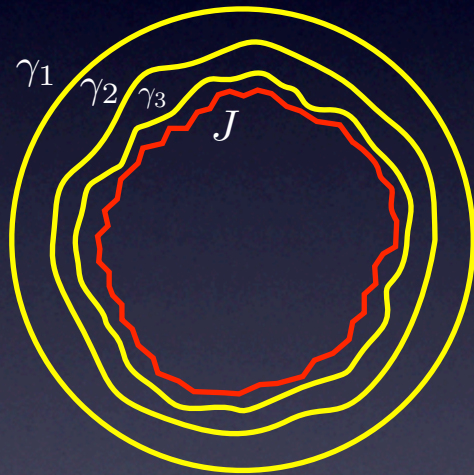
include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



Typical argument to show that some objects are continuous curves: for example, $J(z^2 + \varepsilon)$ is a Jordan curve



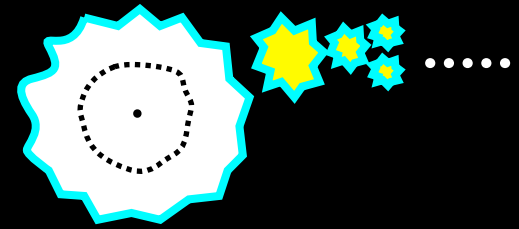
use renormalization!

Frontiers in Complex Dynamics

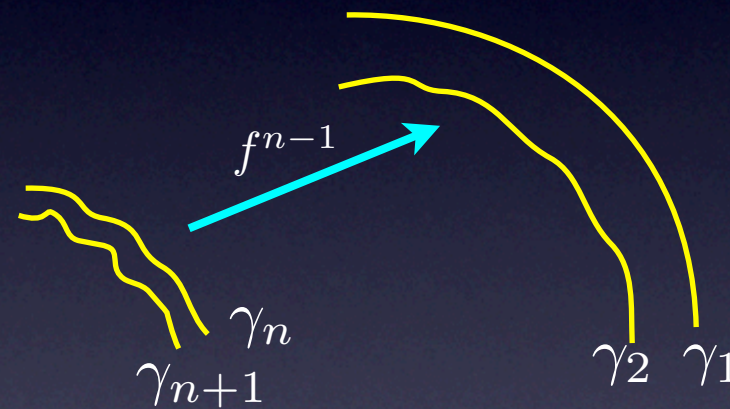
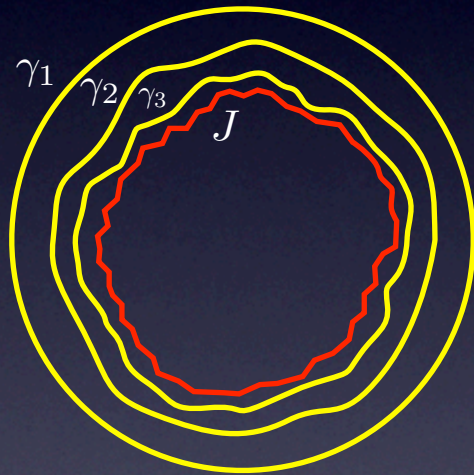
include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



Typical argument to show that some objects are continuous curves: for example, $J(z^2 + \varepsilon)$ is a Jordan curve



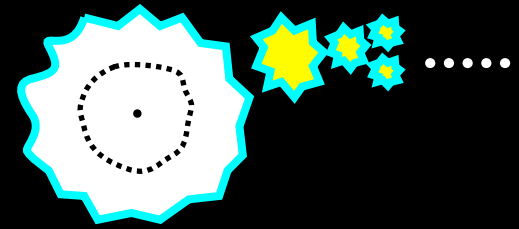
use renormalization!

Frontiers in Complex Dynamics

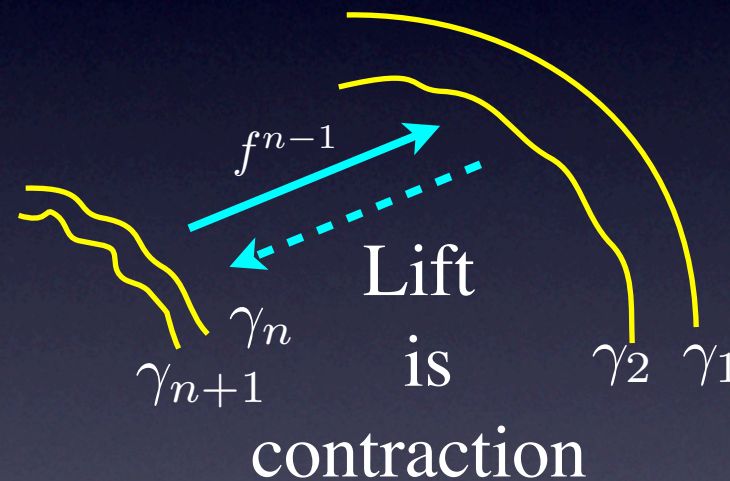
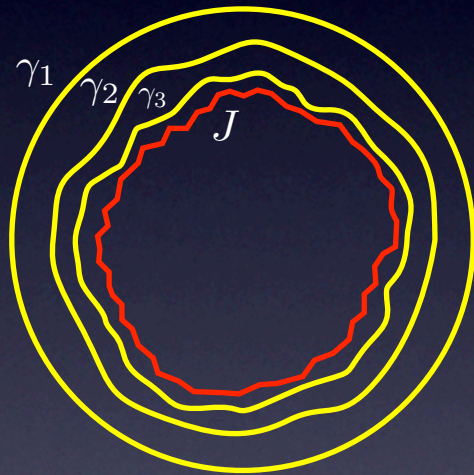
include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



Typical argument to show that some objects are continuous curves: for example, $J(z^2 + \varepsilon)$ is a Jordan curve



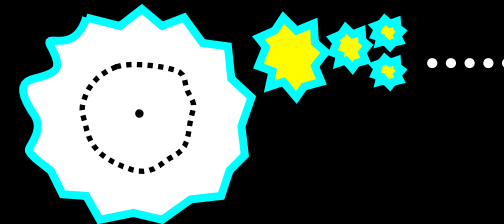
use renormalization!

Frontiers in Complex Dynamics

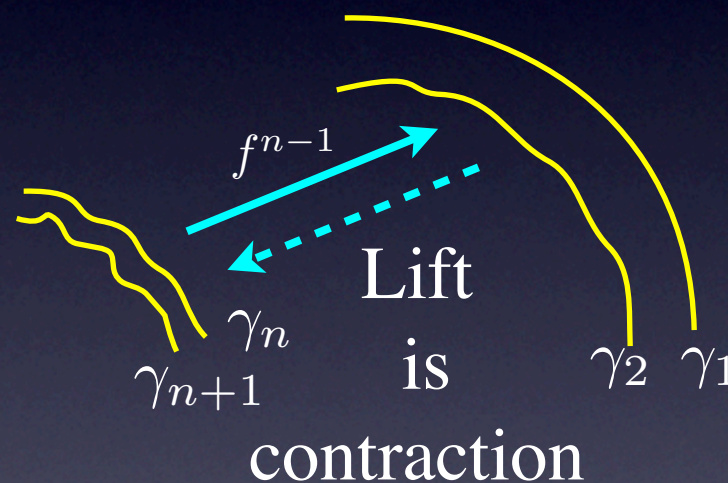
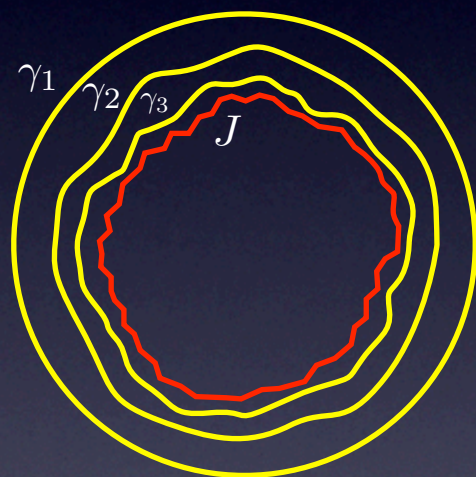
include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



Typical argument to show that some objects are continuous curves: for example, $J(z^2 + \varepsilon)$ is a Jordan curve



Conclusions from Informal Discussion:

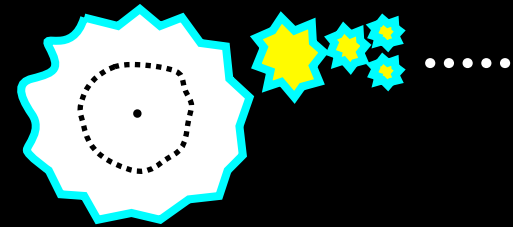
use renormalization!

Frontiers in Complex Dynamics

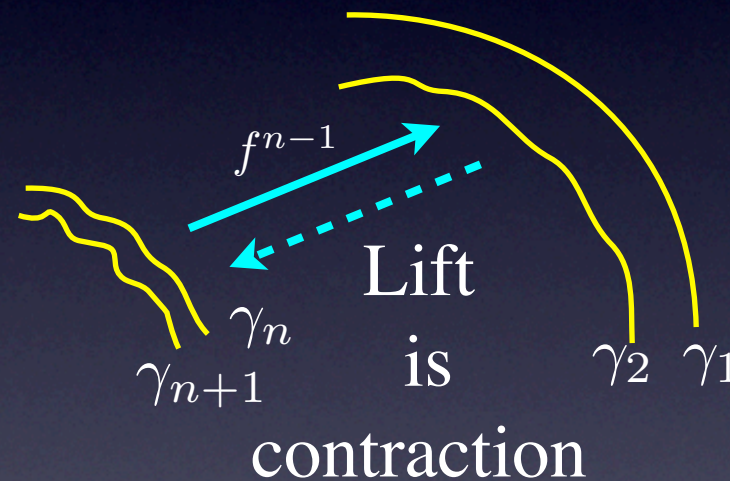
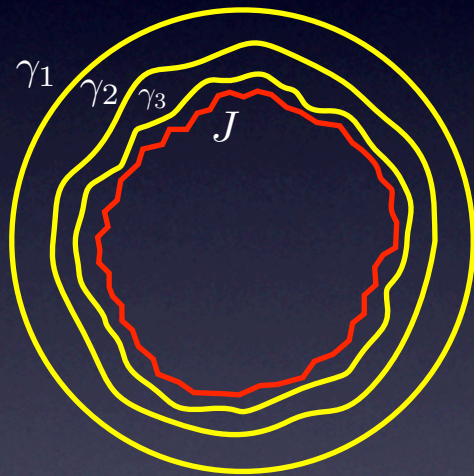
include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



Typical argument to show that some objects are continuous curves: for example, $J(z^2 + \varepsilon)$ is a Jordan curve



Conclusions from Informal Discussion:

Lifts are important.

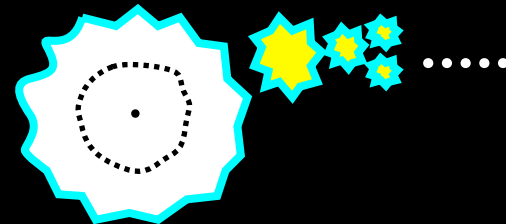
use renormalization!

Frontiers in Complex Dynamics

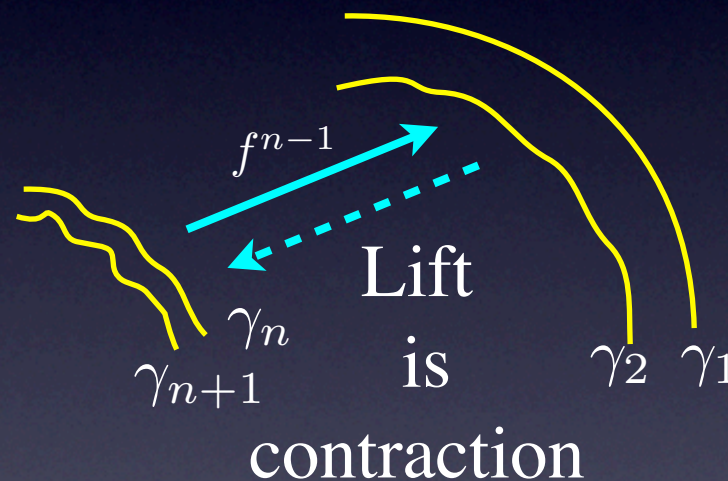
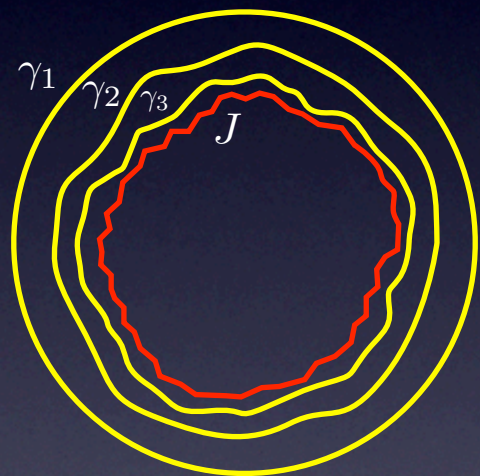
include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



Typical argument to show that some objects are continuous curves: for example, $J(z^2 + \varepsilon)$ is a Jordan curve



Conclusions from Informal Discussion:

Lifts are important.

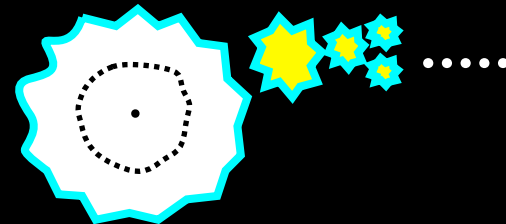
When lifts are not available, need to look for alternatives.
use renormalization!

Frontiers in Complex Dynamics

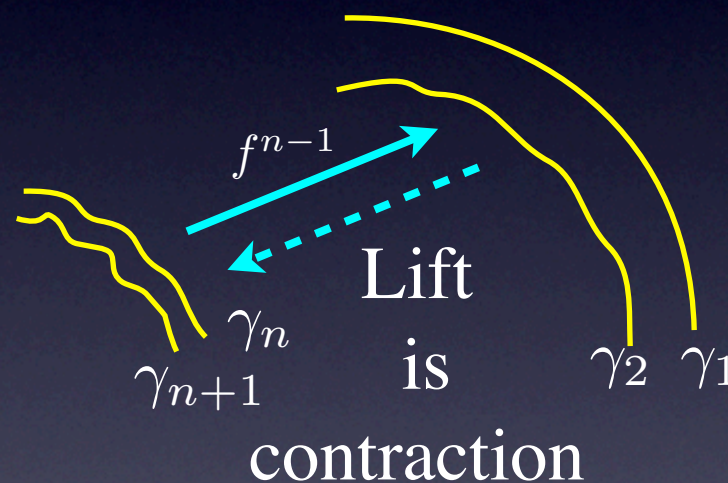
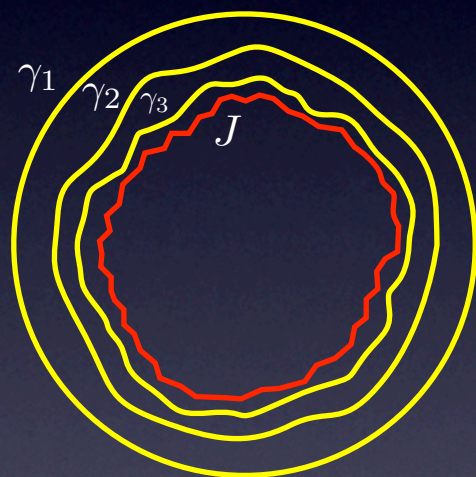
include *boundaries* of Siegel disks

want to show that they are Jordan curves under some condition

Other approaches: Herman, Petersen-Zackeri, Buff-Chéritat



Typical argument to show that some objects are continuous curves: for example, $J(z^2 + \varepsilon)$ is a Jordan curve



Conclusions from Informal Discussion:

Lifts are important.

When lifts are not available, need to look for alternatives.

Gondola, T-bar, ... *use renormalization!*

Two extremes in dynamics

Two extremes in dynamics

Tame (zero entropy, minimal)

Two extremes in dynamics

Tame (zero entropy, minimal)

Chaotic (positive entropy)

Two extremes in dynamics

Tame (zero entropy, minimal)

Chaotic (positive entropy)

Fragile (easy to destroy)

Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Chaotic (positive entropy)

Robust (stable under perturbation)

Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Rigid (conjugacy is smooth)

Chaotic (positive entropy)

Robust (stable under perturbation)

Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Rigid (conjugacy is smooth)

Chaotic (positive entropy)

Robust (stable under perturbation)

Non-rigid (conj. not smooth)

Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Rigid (conjugacy is smooth)

Chaotic (positive entropy)

Robust (stable under perturbation)

Non-rigid (conj. not smooth)

S^1

Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Rigid (conjugacy is smooth)

Chaotic (positive entropy)

Robust (stable under perturbation)

Non-rigid (conj. not smooth)

irrat. rotation

α irrational



S^1

$$t \mapsto t + \alpha \pmod{\mathbb{Z}}$$

Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Rigid (conjugacy is smooth)

Chaotic (positive entropy)

Robust (stable under perturbation)

Non-rigid (conj. not smooth)

irrat. rotation
 α irrational



$$t \mapsto t + \alpha \pmod{\mathbb{Z}}$$

\mathbb{S}^1

doubling map



$$t \mapsto 2t \pmod{\mathbb{Z}}$$

Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Rigid (conjugacy is smooth)

Chaotic (positive entropy)

Robust (stable under perturbation)

Non-rigid (conj. not smooth)

irrat. rotation
 α irrational



$$t \mapsto t + \alpha \pmod{\mathbb{Z}}$$

\mathbb{S}^1

doubling map



$$t \mapsto 2t \pmod{\mathbb{Z}}$$

$T_1 \Sigma_g$

(Σ_g surface $g \geq 2$)

Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Rigid (conjugacy is smooth)

Chaotic (positive entropy)

Robust (stable under perturbation)

Non-rigid (conj. not smooth)

irrat. rotation
 α irrational



$$t \mapsto t + \alpha \pmod{\mathbb{Z}}$$

\mathbb{S}^1

doubling map

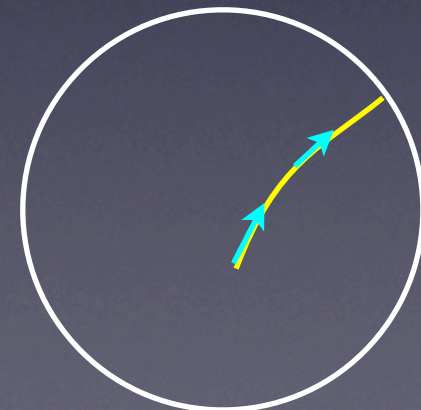


$$t \mapsto 2t \pmod{\mathbb{Z}}$$

$T_1\Sigma_g$

(Σ_g surface $g \geq 2$)

geodesic flow



Two extremes in dynamics

Tame (zero entropy, minimal)

Fragile (easy to destroy)

Rigid (conjugacy is smooth)

Chaotic (positive entropy)

Robust (stable under perturbation)

Non-rigid (conj. not smooth)

irrat. rotation
 α irrational



$$t \mapsto t + \alpha \pmod{\mathbb{Z}}$$

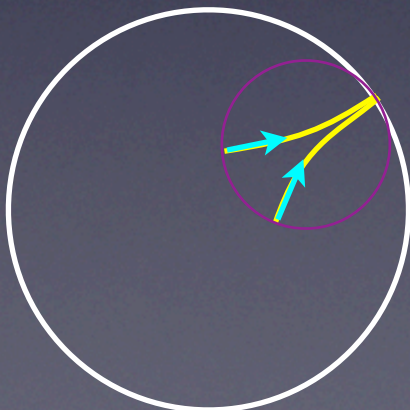
S^1

doubling map



$$t \mapsto 2t \pmod{\mathbb{Z}}$$

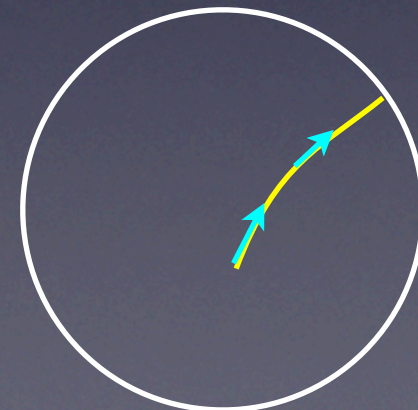
horocyclic flow



$T_1\Sigma_g$

(Σ_g surface $g \geq 2$)

geodesic flow



Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

Chaotic

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

Chaotic

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

Chaotic

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
(stable)

Chaotic

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

geodesic flow g^t on $T_1\Sigma_g$

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
(stable)

$$h^{\lambda t} \circ g^t = g^t \circ h^t$$

Chaotic

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

geodesic flow g^t on $T_1\Sigma_g$

Two extreme dynamics are often related
("intertwining relation" or "renormalization pair")

Tame

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
(stable)

$$h^{\lambda t} \circ g^t = g^t \circ h^t$$

Chaotic

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

geodesic flow g^t on $T_1\Sigma_g$

shift σ on $\Sigma = \{0, 1\}^{\mathbb{N}}$
 $x_0x_1x_2 \dots \mapsto x_1x_2x_3 \dots$

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

Chaotic

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
 (stable)

geodesic flow g^t on $T_1\Sigma_g$

$$h^{\lambda t} \circ g^t = g^t \circ h^t$$

adding machine τ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

shift σ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

$$\begin{array}{c} 1110010 \dots \\ \cdot \quad \cdot \quad \cdot \end{array} \mapsto \begin{array}{c} 0001010 \dots \\ \cdot \quad \cdot \quad \cdot \end{array} \dots$$

$$x_0x_1x_2 \dots \mapsto x_1x_2x_3 \dots$$

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

Chaotic

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
 (stable)

geodesic flow g^t on $T_1\Sigma_g$

$$h^{\lambda t} \circ g^t = g^t \circ h^t$$

adding machine τ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

shift σ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

$$\begin{array}{c} 1110010 \dots \\ \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \\ \dots \end{array} \mapsto \begin{array}{c} 0001010 \dots \\ \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \quad \cdot \\ \dots \end{array}$$

$$x_0x_1x_2 \dots \mapsto x_1x_2x_3 \dots$$

$$\sigma \circ \tau^2 = \tau \circ \sigma$$

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

Chaotic

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
 (stable)

geodesic flow g^t on $T_1\Sigma_g$

$$h^{\lambda t} \circ g^t = g^t \circ h^t$$

adding machine τ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

shift σ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

$$\begin{array}{c} 1110010 \dots \\ \cdot \cdot \cdot \end{array} \mapsto \begin{array}{c} 0001010 \dots \\ \cdot \end{array} \dots$$

$$x_0x_1x_2 \dots \mapsto x_1x_2x_3 \dots$$

$$\sigma \circ \tau^2 = \tau \circ \sigma$$

Anosov diffeo F on T^2

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

Chaotic

$$R_\alpha : t \mapsto t + \alpha \text{ on } \mathbb{R}/\mathbb{Z}$$

$$F : t \mapsto 2t \text{ on } \mathbb{R}/\mathbb{Z}$$

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
 (stable)

geodesic flow g^t on $T_1\Sigma_g$

$$h^{\lambda t} \circ g^t = g^t \circ h^t$$

adding machine τ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

shift σ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

$$\begin{array}{c} 1110010 \dots \\ \cdot \quad \cdot \quad \cdot \end{array} \mapsto \begin{array}{c} 0001010 \dots \\ \cdot \quad \cdot \quad \cdot \end{array}$$

$$x_0x_1x_2 \dots \mapsto x_1x_2x_3 \dots$$

$$\sigma \circ \tau^2 = \tau \circ \sigma$$

irrat. flow ϕ^t on T^2
 (along expanding direction of F)

Anosov diffeo F on T^2

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

Chaotic

$R_\alpha : t \mapsto t + \alpha$ on \mathbb{R}/\mathbb{Z}

$F : t \mapsto 2t$ on \mathbb{R}/\mathbb{Z}

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
 (stable)

geodesic flow g^t on $T_1\Sigma_g$

$$h^{\lambda t} \circ g^t = g^t \circ h^t$$

adding machine τ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

shift σ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

$\begin{array}{c} 1110010 \dots \\ \cdot \cdot \cdot \end{array} \mapsto \begin{array}{c} 0001010 \dots \\ \cdot \end{array} \dots$

$x_0x_1x_2 \dots \mapsto x_1x_2x_3 \dots$

$$\sigma \circ \tau^2 = \tau \circ \sigma$$

irrat. flow ϕ^t on T^2

Anosov diffeo F on T^2

(along expanding direction of F)

$$\phi^{\lambda t} \circ F = F \circ \phi^t$$

Two extreme dynamics are often related
 (“intertwining relation” or “renormalization pair”)

Tame

Chaotic

$R_\alpha : t \mapsto t + \alpha$ on \mathbb{R}/\mathbb{Z}

$F : t \mapsto 2t$ on \mathbb{R}/\mathbb{Z}

$$R_\alpha^2 \circ F = F \circ R_\alpha$$

horocyclic flow h^t on $T_1\Sigma_g$
 (stable)

geodesic flow g^t on $T_1\Sigma_g$

$$h^{\lambda t} \circ g^t = g^t \circ h^t$$

adding machine τ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

shift σ on $\Sigma = \{0, 1\}^{\mathbb{N}}$

$1110010 \dots \mapsto 0001010 \dots$

$x_0x_1x_2 \dots \mapsto x_1x_2x_3 \dots$

$$\sigma \circ \tau^2 = \tau \circ \sigma$$

irrat. flow ϕ^t on T^2

Anosov diffeo F on T^2

(along expanding direction of F)

$$\phi^{\lambda t} \circ F = F \circ \phi^t$$

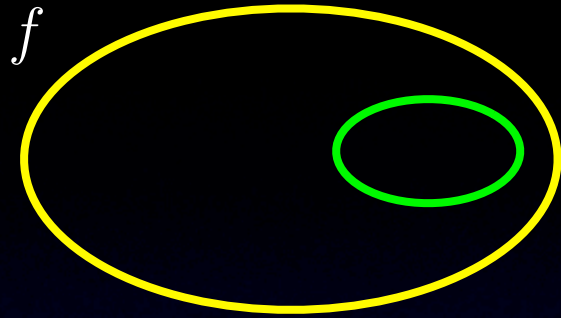
Conjugation by chaotic one is like a time change for tame one

Return map and renormalization

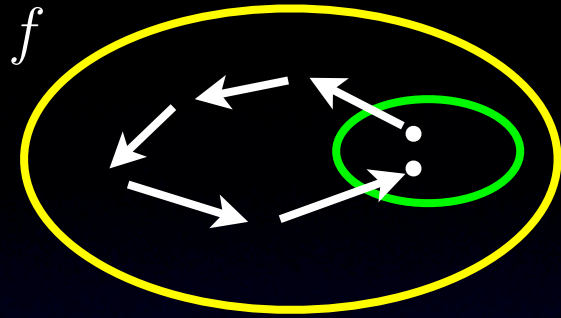
Return map and renormalization



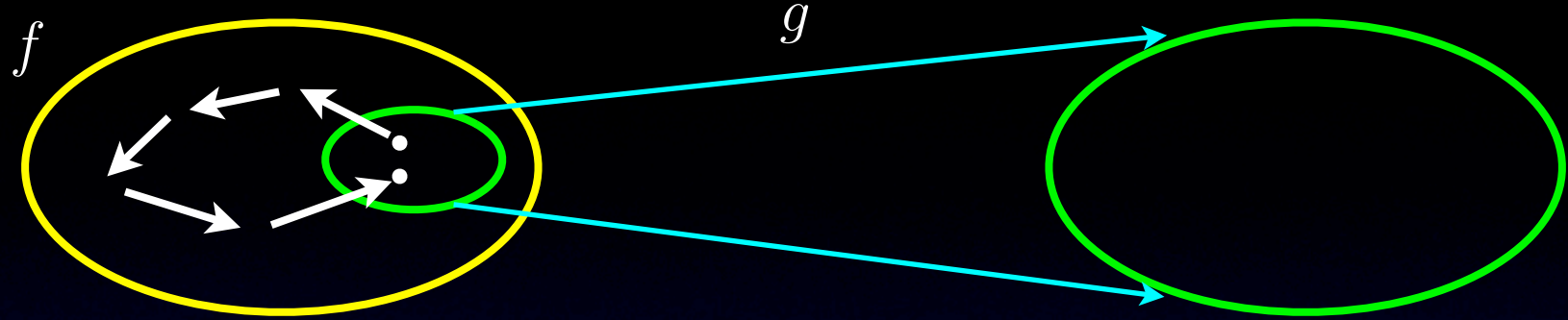
Return map and renormalization



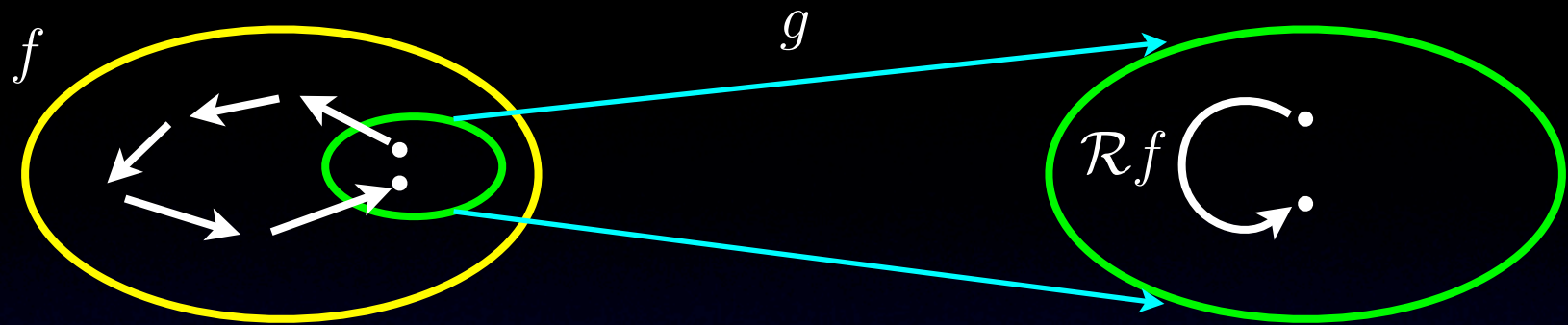
Return map and renormalization



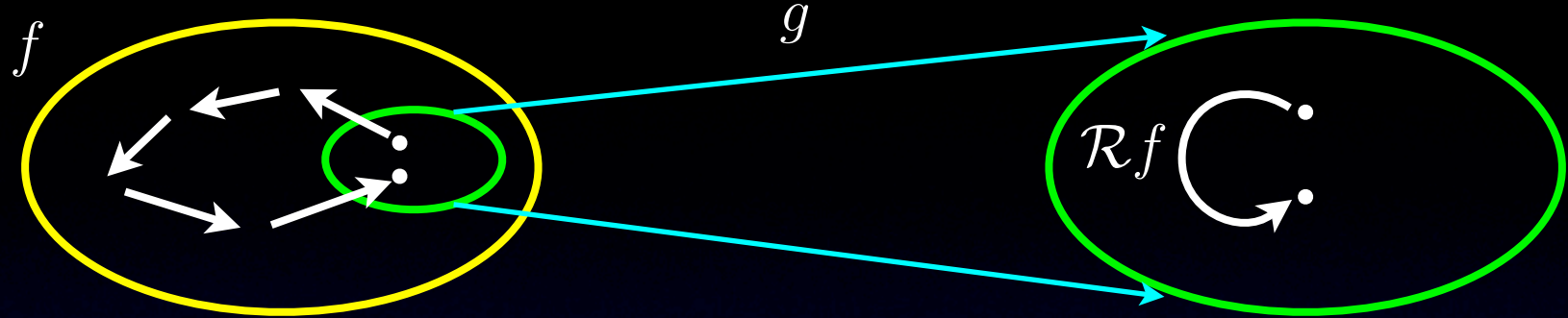
Return map and renormalization



Return map and renormalization

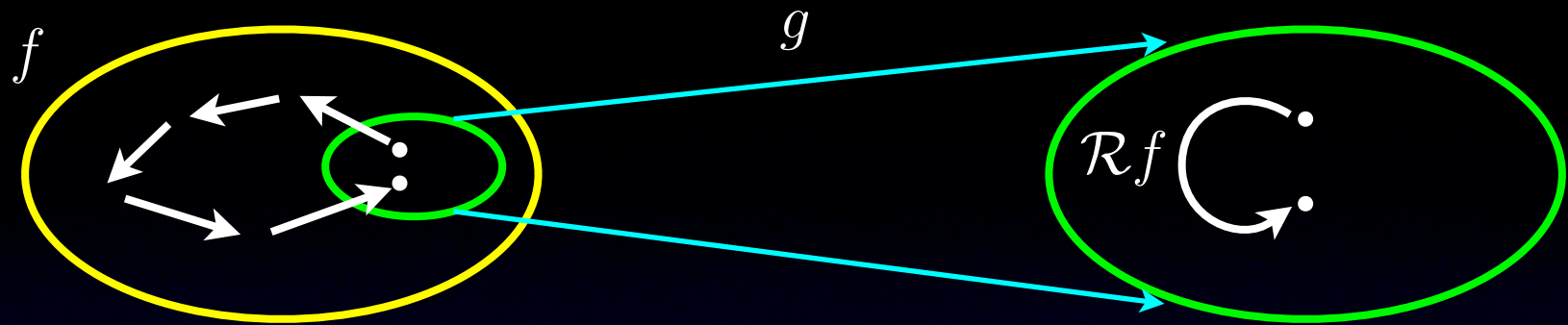


Return map and renormalization



$$\begin{aligned}\mathcal{R}f &= (\text{first return map of } f) \text{ after rescaling} \\ &= g \circ f^k \circ g^{-1} \quad (\text{if return time } \equiv k)\end{aligned}$$

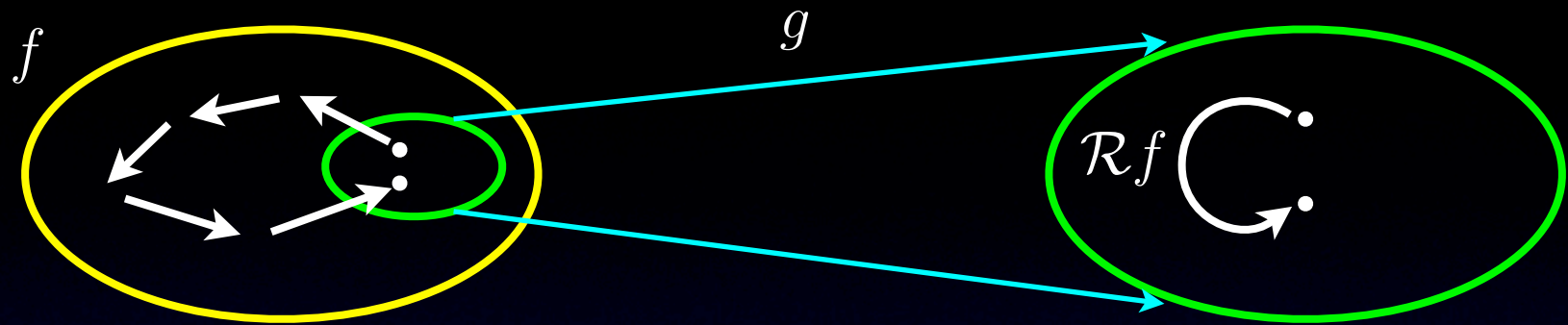
Return map and renormalization



$$\begin{aligned}\mathcal{R}f &= (\text{first return map of } f) \text{ after rescaling} \\ &= g \circ f^k \circ g^{-1} \quad (\text{if return time } \equiv k)\end{aligned}$$

Renormalization

Return map and renormalization



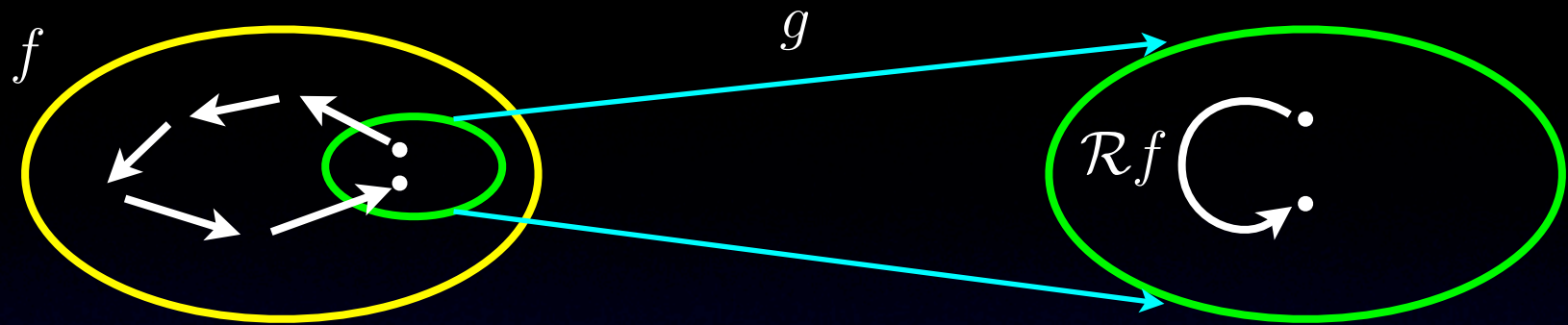
$$\begin{aligned}\mathcal{R}f &= (\text{first return map of } f) \text{ after rescaling} \\ &= g \circ f^k \circ g^{-1} \quad (\text{if return time } \equiv k)\end{aligned}$$

Renormalization

high iterates of f \longleftrightarrow fewer iterates of $\mathcal{R}f$
fine orbit structure for f \longleftrightarrow large scale orbit structure for $\mathcal{R}f$

Successive construction of $\mathcal{R}f, \mathcal{R}^2f, \dots$, helps to understand the dynamics of f (orbits, invariant sets, rigidity, bifurcation, ...)

Return map and renormalization



$$\begin{aligned} \mathcal{R}f &= (\text{first return map of } f) \text{ after rescaling} \\ &= g \circ f^k \circ g^{-1} \quad (\text{if return time } \equiv k) \end{aligned}$$

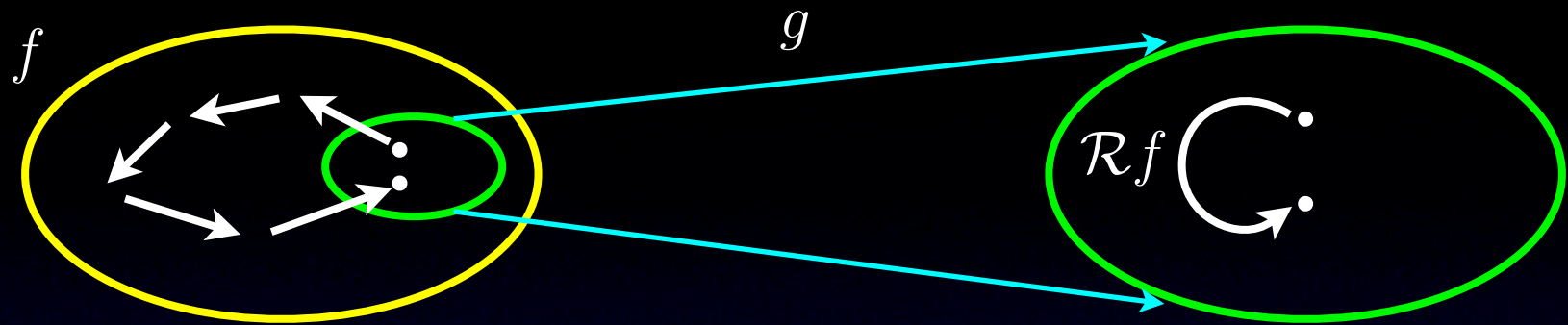
Renormalization

high iterates of f \longleftrightarrow fewer iterates of $\mathcal{R}f$
 fine orbit structure for f \longleftrightarrow large scale orbit structure for $\mathcal{R}f$

Successive construction of $\mathcal{R}f, \mathcal{R}^2f, \dots$, helps to understand the dynamics of f (orbits, invariant sets, rigidity, bifurcation, ...)

If f is a fixed point of renormalization \mathcal{R} (with return time $\equiv k$), then $g \circ f^k \circ g^{-1} = f$, i.e. $g \circ f^k = f \circ g$ (intertwining relation).

Return map and renormalization



$$\begin{aligned} \mathcal{R}f &= (\text{first return map of } f) \text{ after rescaling} \\ &= g \circ f^k \circ g^{-1} \quad (\text{if return time } \equiv k) \end{aligned}$$

Renormalization

high iterates of f \longleftrightarrow fewer iterates of $\mathcal{R}f$
 fine orbit structure for f \longleftrightarrow large scale orbit structure for $\mathcal{R}f$

Successive construction of $\mathcal{R}f, \mathcal{R}^2f, \dots$, helps to understand the dynamics of f (orbits, invariant sets, rigidity, bifurcation, ...)

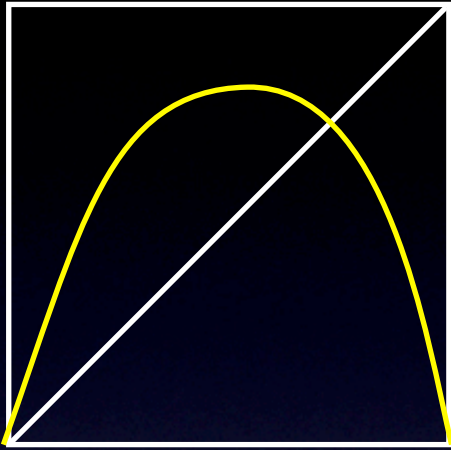
If f is a fixed point of renormalization \mathcal{R} (with return time $\equiv k$), then $g \circ f^k \circ g^{-1} = f$, i.e. $g \circ f^k = f \circ g$ (intertwining relation).

Usually f tame and g expanding (chaotic).

Feigenbaum-Coulet-Tresser for unimodal maps

Feigenbaum-Coulet-Tresser for unimodal maps

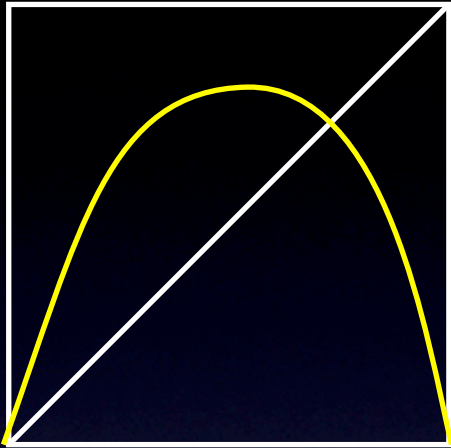
f



$$I = [0, 1]$$

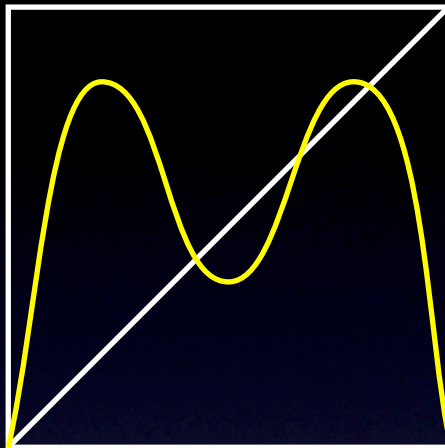
Feigenbaum-Coulet-Tresser for unimodal maps

f



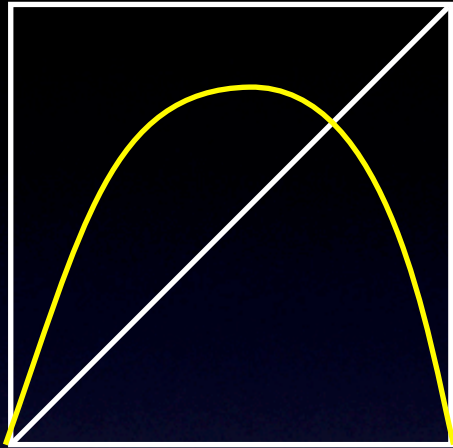
$I = [0, 1]$

f^2



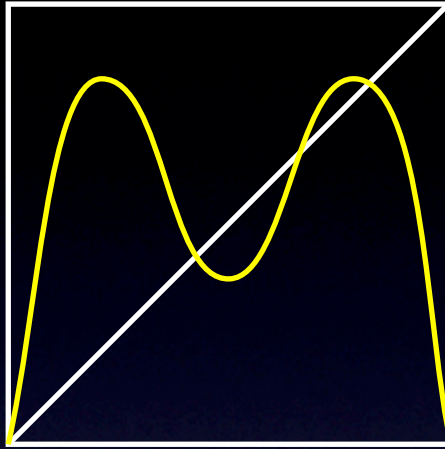
Feigenbaum-Coulet-Tresser for unimodal maps

f



$$I = [0, 1]$$

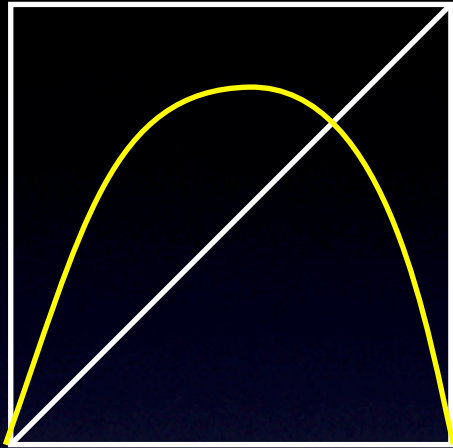
f^2



$$J \subset I \text{ s.t. } f^2(J) \subset J$$

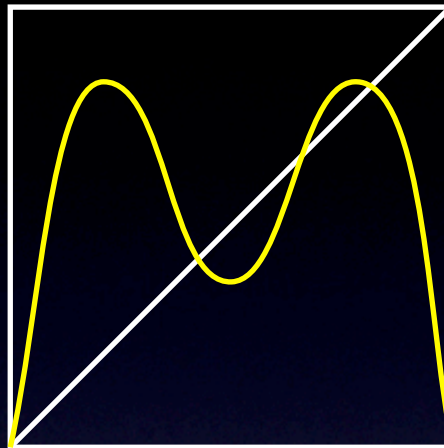
Feigenbaum-Coulet-Tresser for unimodal maps

f



$$I = [0, 1]$$

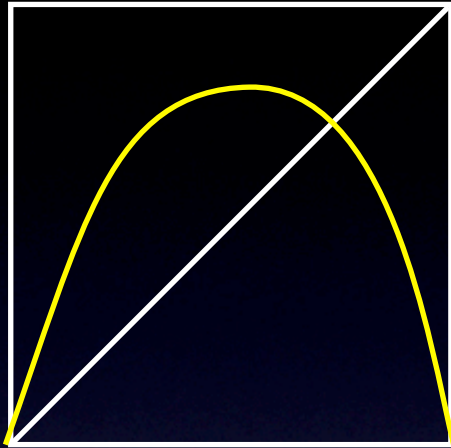
f^2



$$J \subset I \text{ s.t. } f^2(J) \subset J$$

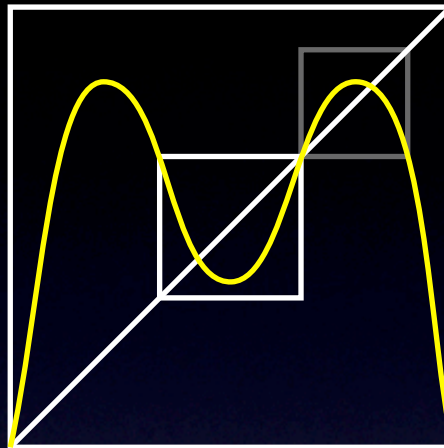
Feigenbaum-Coulet-Tresser for unimodal maps

f



$$I = [0, 1]$$

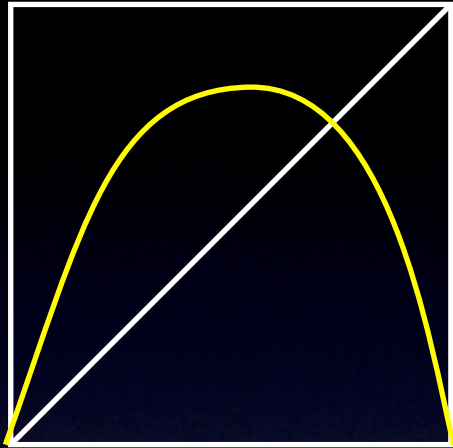
f^2



$$J \subset I \text{ s.t. } f^2(J) \subset J$$

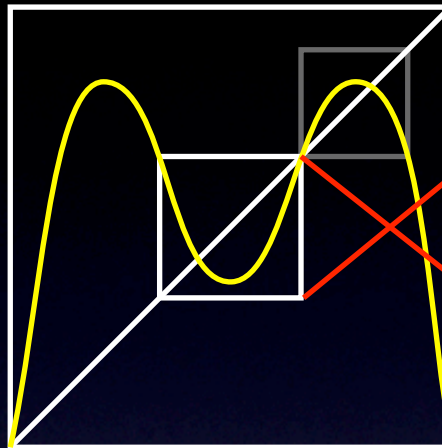
Feigenbaum-Coulet-Tresser for unimodal maps

f



$$I = [0, 1]$$

f^2



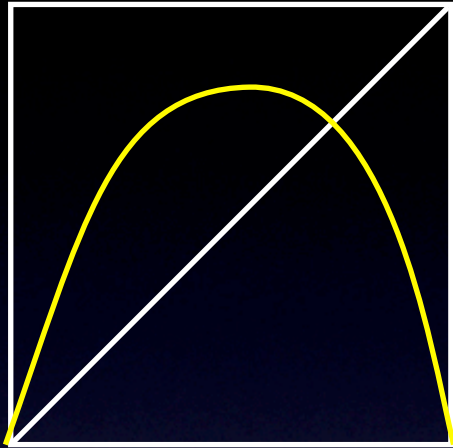
$$J \subset I \text{ s.t. } f^2(J) \subset J$$

g



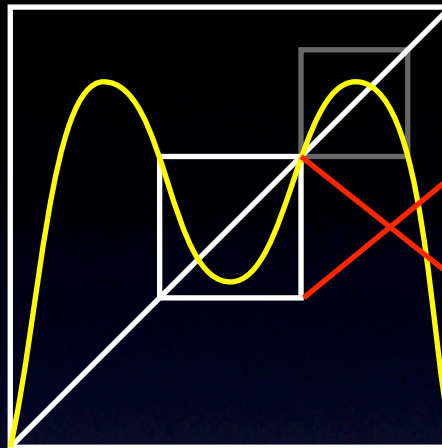
Feigenbaum-Coullet-Tresser for unimodal maps

f



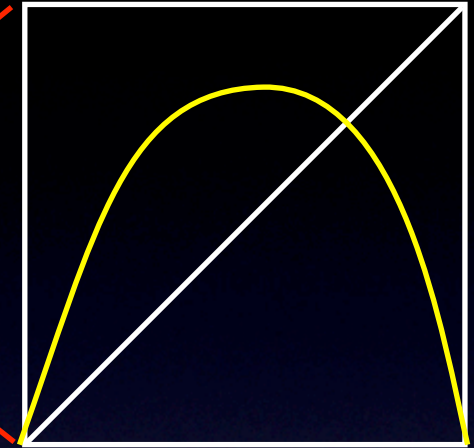
$$I = [0, 1]$$

f^2



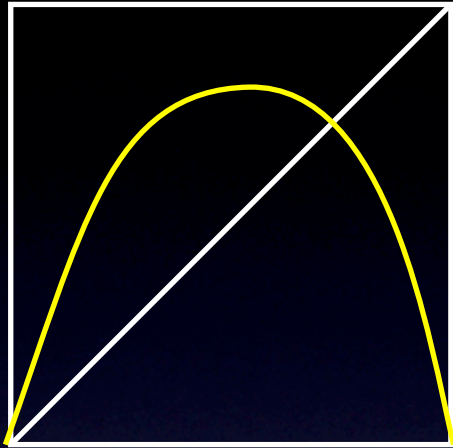
$$J \subset I \text{ s.t. } f^2(J) \subset J$$

g



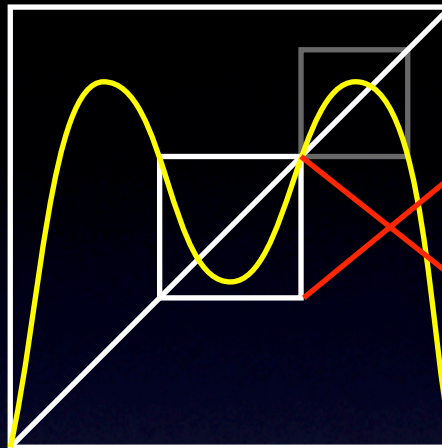
Feigenbaum-Coullet-Tresser for unimodal maps

f



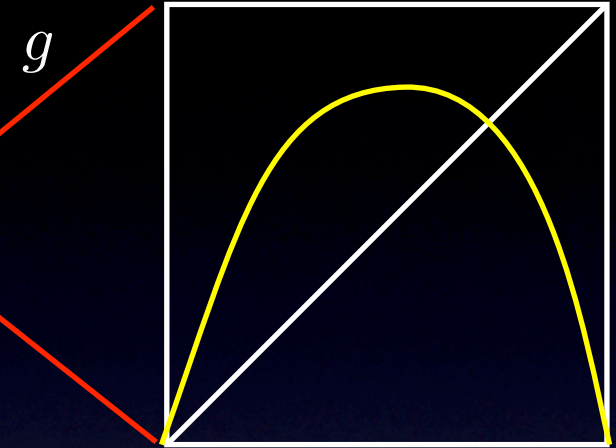
$I = [0, 1]$

f^2



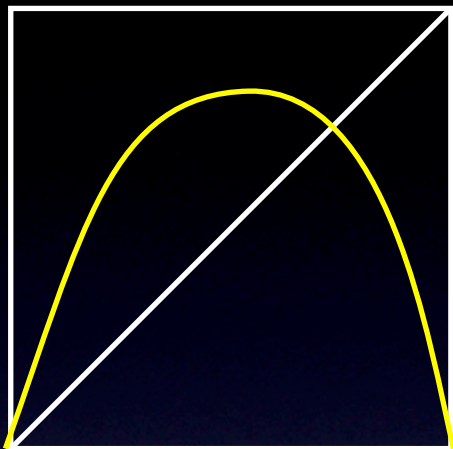
$J \subset I$ s.t. $f^2(J) \subset J$

$\mathcal{R}f = g \circ (f^2|_J) \circ g^{-1}$



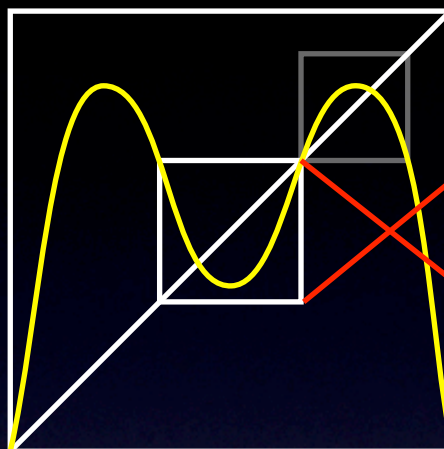
Feigenbaum-Coulet-Tresser for unimodal maps

f



$$I = [0, 1]$$

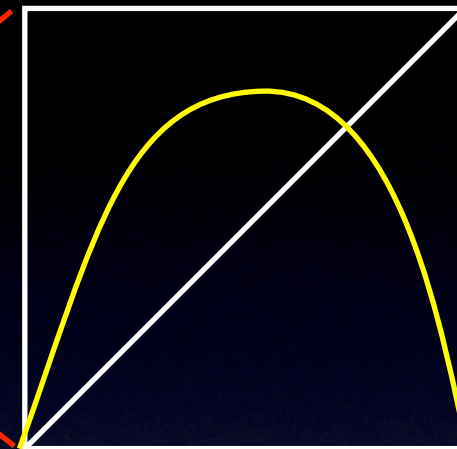
f^2



$$J \subset I \text{ s.t. } f^2(J) \subset J$$

$$\mathcal{R}f = g \circ (f^2|_J) \circ g^{-1}$$

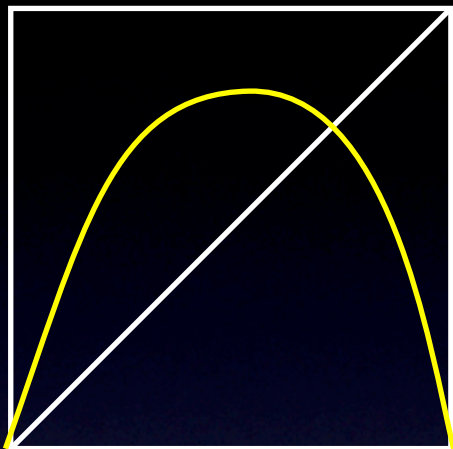
g



Renormalization: Meta-dynamics
Dynamics *on the space of certain dynamical systems*

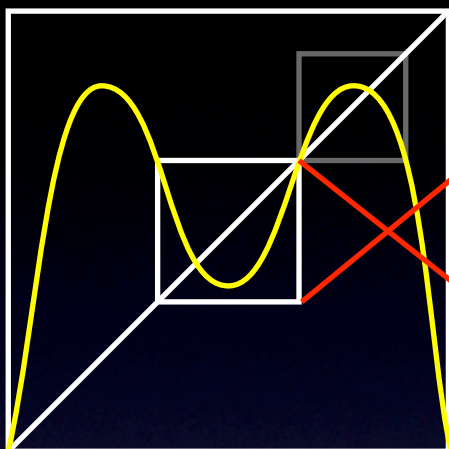
Feigenbaum-Coulet-Tresser for unimodal maps

f



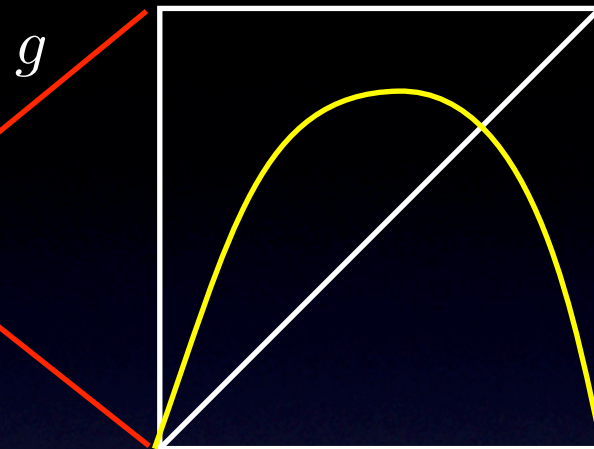
$$I = [0, 1]$$

f^2



$$J \subset I \text{ s.t. } f^2(J) \subset J$$

$$\mathcal{R}f = g \circ (f^2|_J) \circ g^{-1}$$

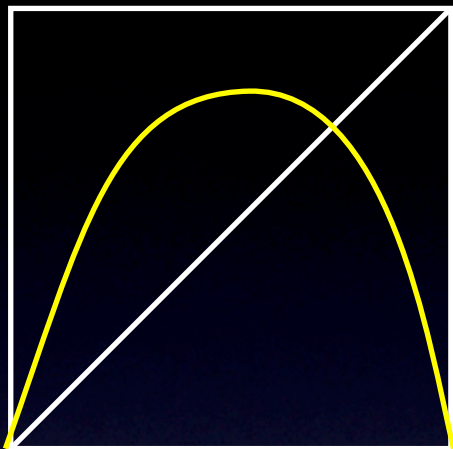


Renormalization: Meta-dynamics
Dynamics *on the space of certain dynamical systems*

Hyperbolic fixed point or hyperbolic horseshoe of the meta-dynamics imply conclusion on rigidity and structure of parameter space

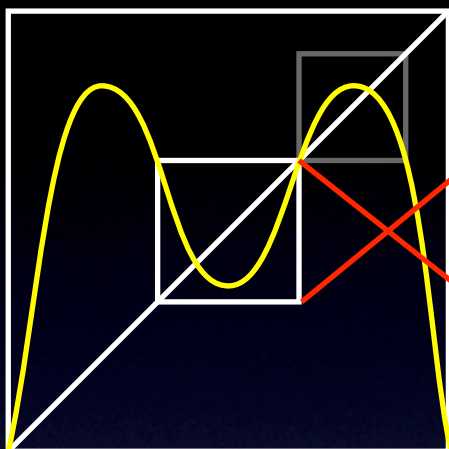
Feigenbaum-Coulet-Tresser for unimodal maps

f



$I = [0, 1]$

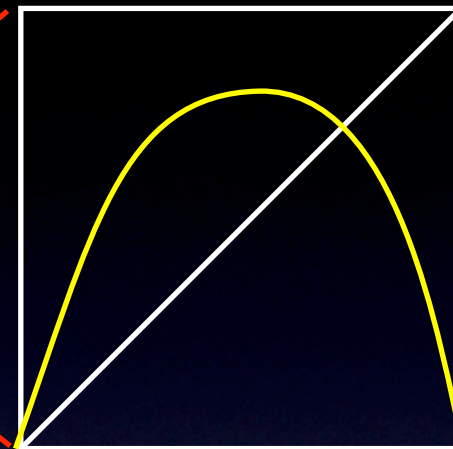
f^2



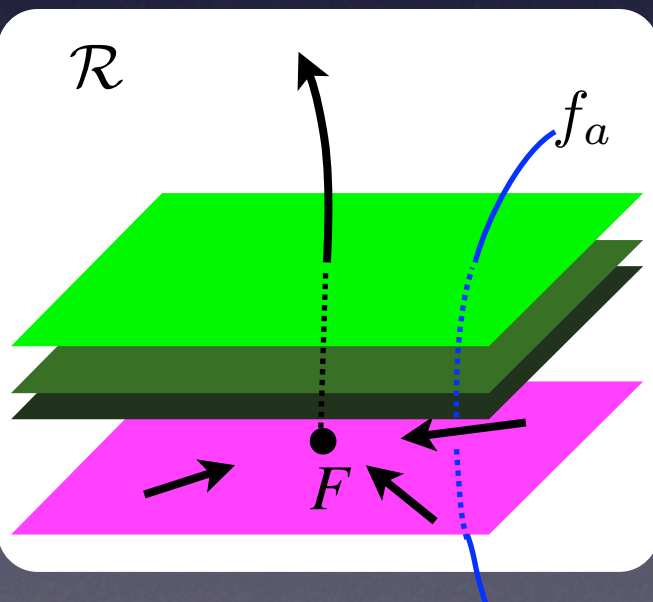
$J \subset I$ s.t. $f^2(J) \subset J$

$\mathcal{R}f = g \circ (f^2|_J) \circ g^{-1}$

g



\mathcal{R}

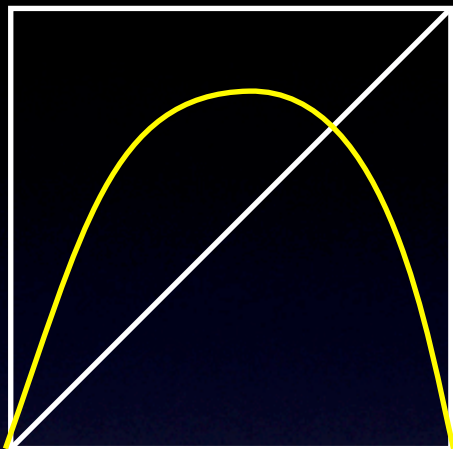


Renormalization: Meta-dynamics
Dynamics *on the space of certain dynamical systems*

Hyperbolic fixed point or hyperbolic horseshoe of the meta-dynamics imply conclusion on rigidity and structure of parameter space

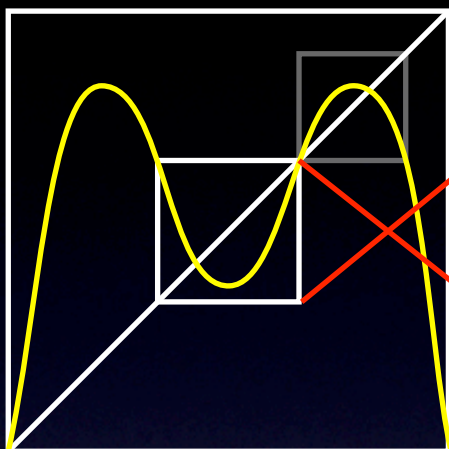
Feigenbaum-Coulet-Tresser for unimodal maps

f



$I = [0, 1]$

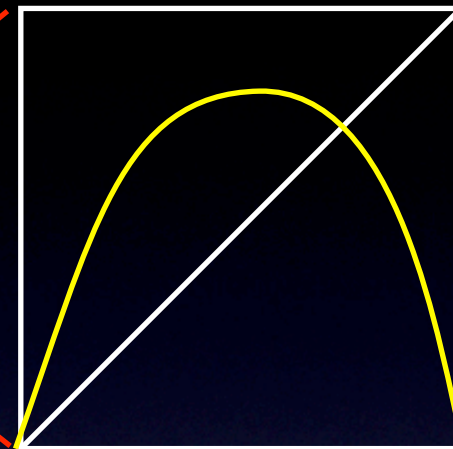
f^2



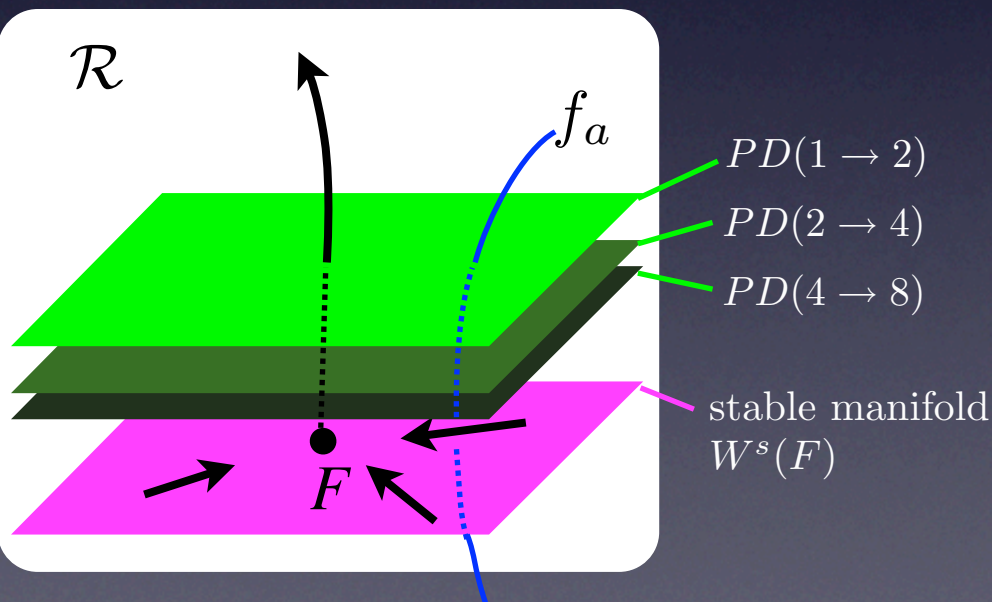
$J \subset I$ s.t. $f^2(J) \subset J$

$\mathcal{R}f = g \circ (f^2|_J) \circ g^{-1}$

g



\mathcal{R}



Renormalization: Meta-dynamics
Dynamics *on the space of certain dynamical systems*

Hyperbolic fixed point or hyperbolic horseshoe of the meta-dynamics imply conclusion on rigidity and structure of parameter space

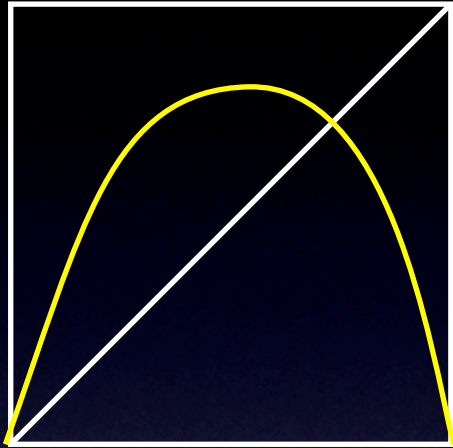
Various Renormalizations

Various Renormalizations

Feigenbaum

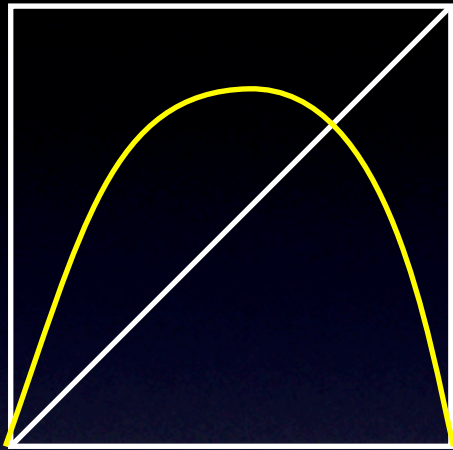
Various Renormalizations

Feigenbaum



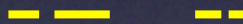
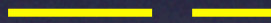
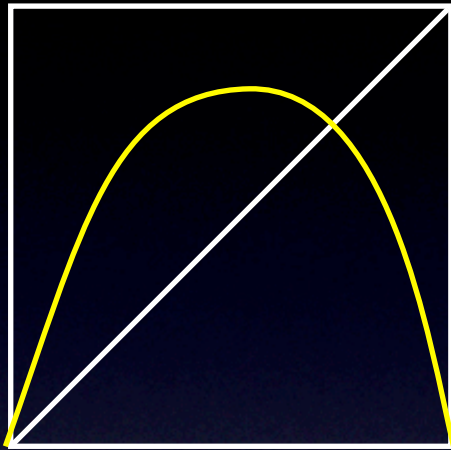
Various Renormalizations

Feigenbaum



Various Renormalizations

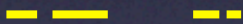
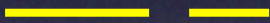
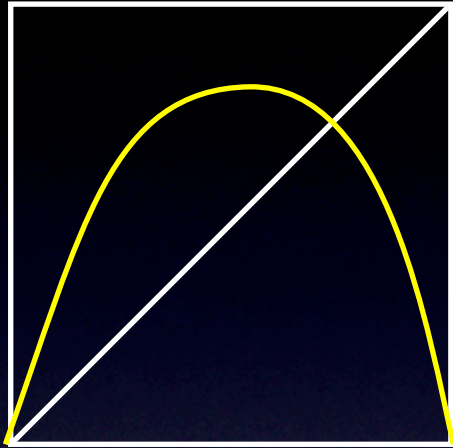
Feigenbaum



proper subintervals

Various Renormalizations

Feigenbaum

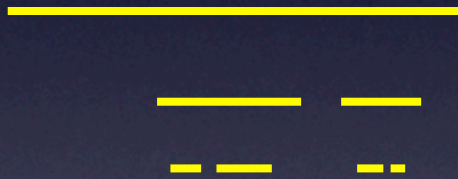
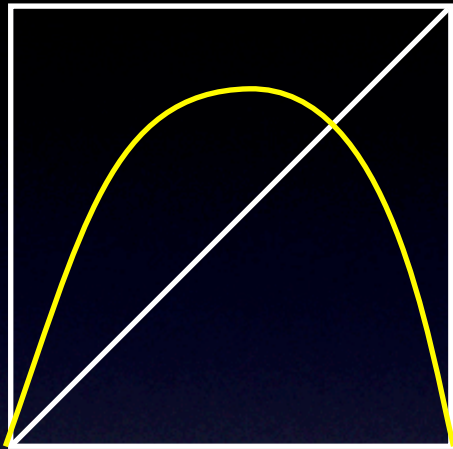


proper subintervals

-> Cantor set

Various Renormalizations

Feigenbaum



proper subintervals

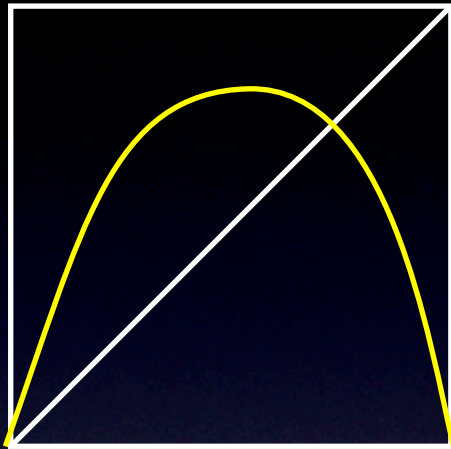
-> Cantor set

Feigenbaum, Couillet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Various Renormalizations

Feigenbaum

Circle map



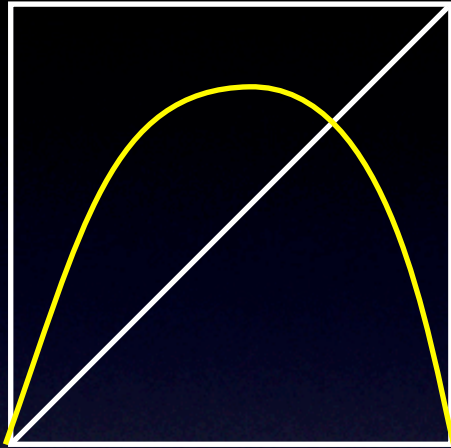
proper subintervals

-> Cantor set

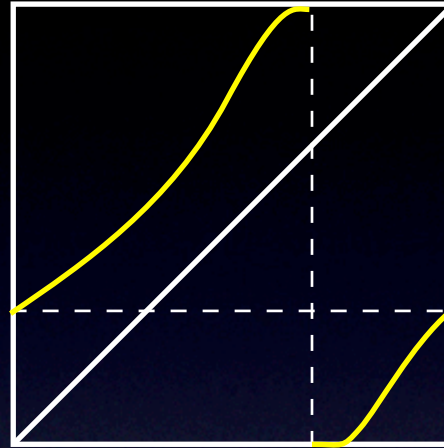
Feigenbaum, Couillet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Various Renormalizations

Feigenbaum



Circle map



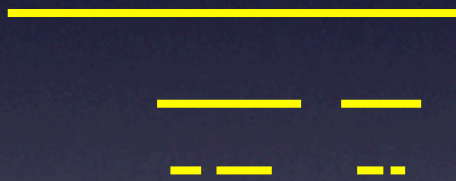
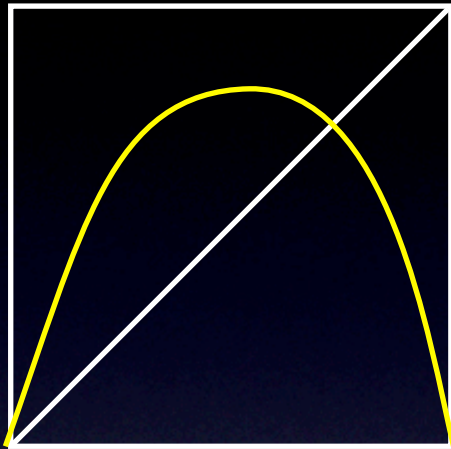
proper subintervals

-> Cantor set

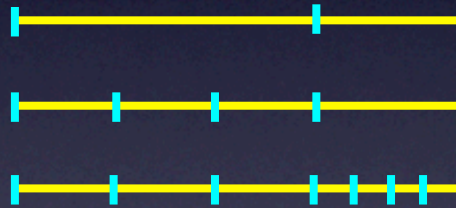
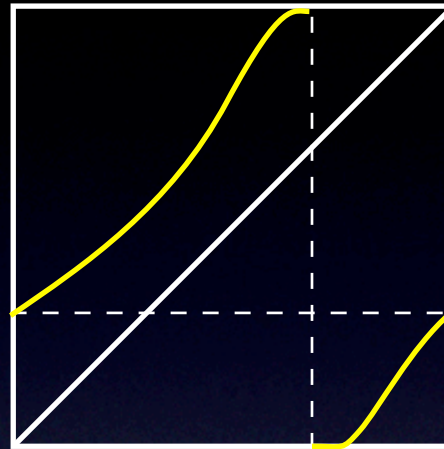
Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Various Renormalizations

Feigenbaum



Circle map



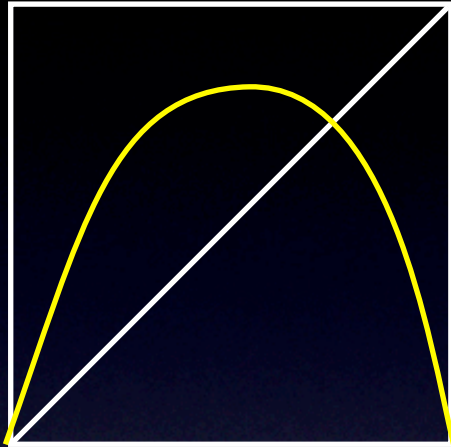
proper subintervals

-> Cantor set

Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Various Renormalizations

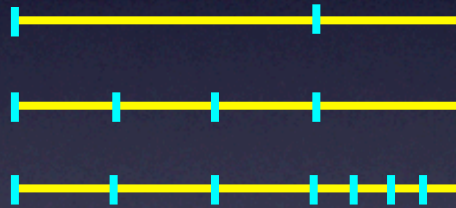
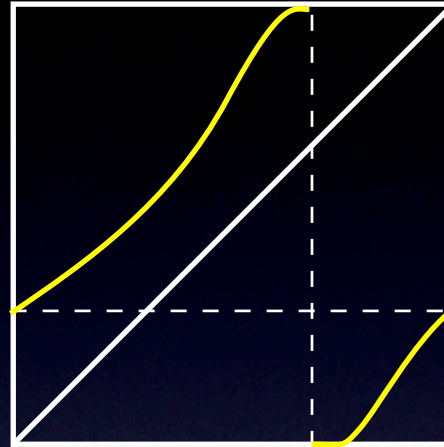
Feigenbaum



proper subintervals

-> Cantor set

Circle map

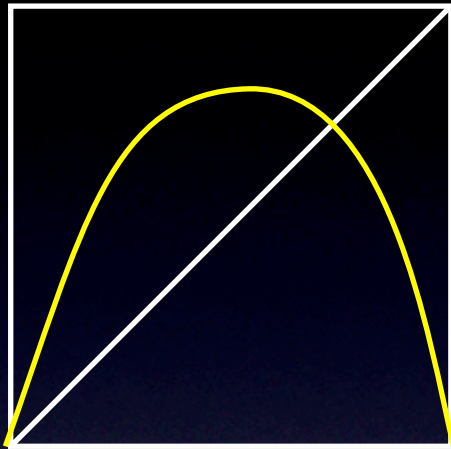


partition of interval

Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Various Renormalizations

Feigenbaum

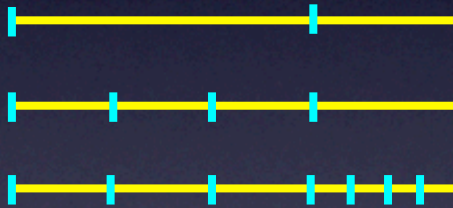
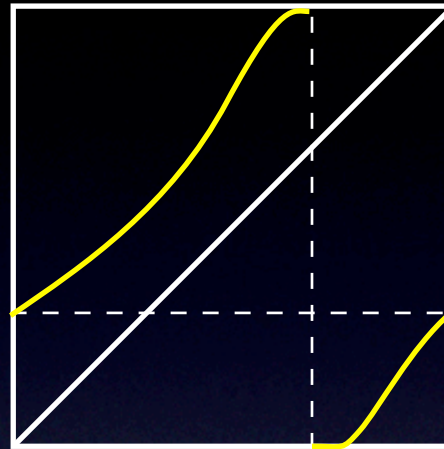


proper subintervals

-> Cantor set

Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Circle map

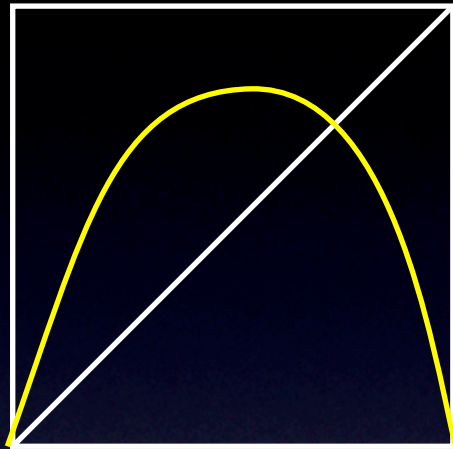


partition of interval

Rand,
Khanin-Sinai,
de Faria, Yampolsky,
A. Epstein-Yamplosky

Various Renormalizations

Feigenbaum

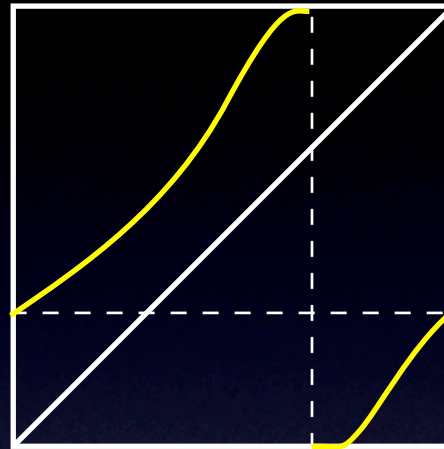


proper subintervals

-> Cantor set

Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Circle map



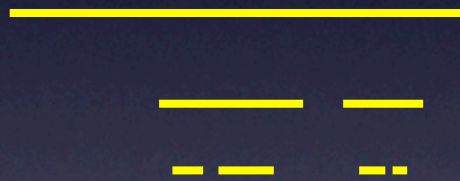
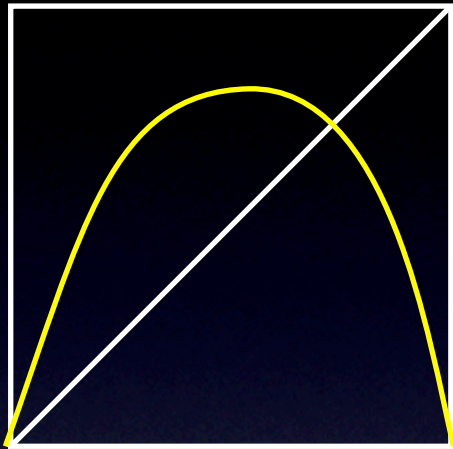
partition of interval

Rand,
Khanin-Sinai,
de Faria, Yampolsky,
A. Epstein-Yamplosky

Sector/Near-parabolic

Various Renormalizations

Feigenbaum

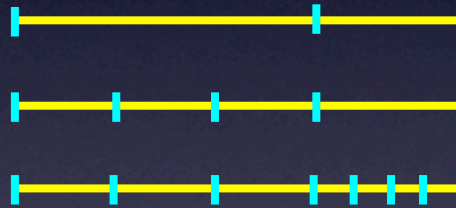
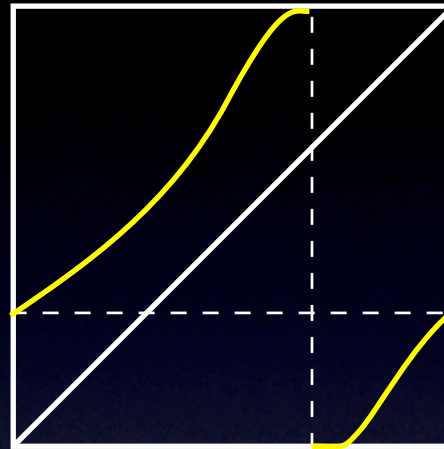


proper subintervals

-> Cantor set

Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Circle map



partition of interval

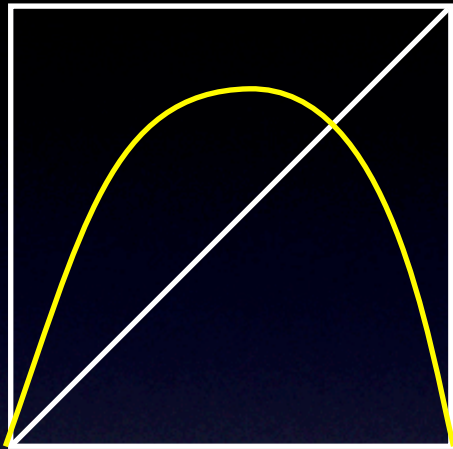
Rand,
Khanin-Sinai,
de Faria, Yampolsky,
A. Epstein-Yamplosky

Sector/Near-parabolic



Various Renormalizations

Feigenbaum

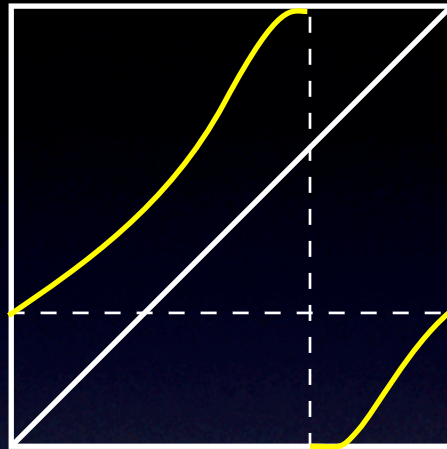


proper subintervals

-> Cantor set

Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Circle map



partition of interval

Rand,
Khanin-Sinai,
de Faria, Yampolsky,
A. Epstein-Yamplosky

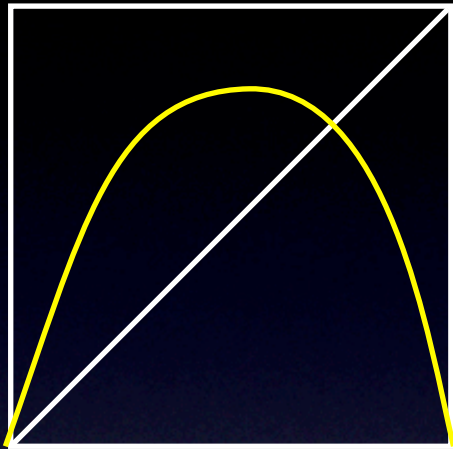
Sector/Near-parabolic



covering by sector or
croissant-like domains

Various Renormalizations

Feigenbaum

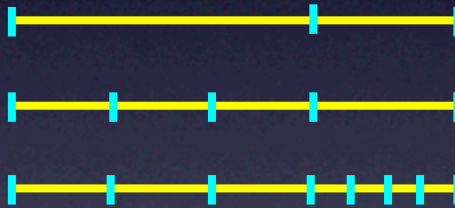
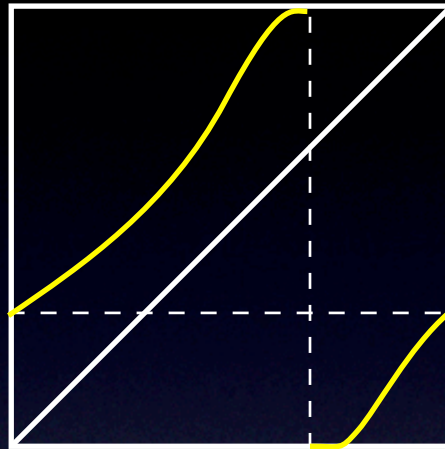


proper subintervals

-> Cantor set

Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Circle map



partition of interval

Rand,
Khanin-Sinai,
de Faria, Yampolsky,
A. Epstein-Yamplosky

Sector/Near-parabolic

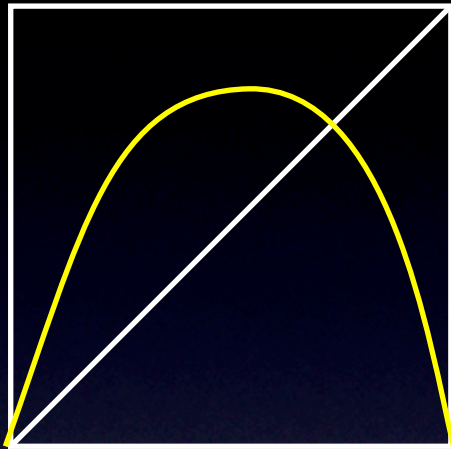


covering by sector or
croissant-like domains

gluing/identification
needed to define the
renormalization

Various Renormalizations

Feigenbaum

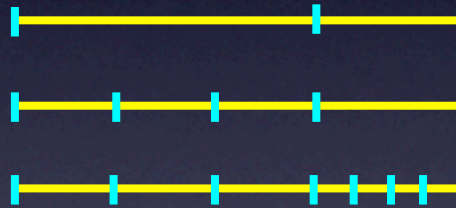
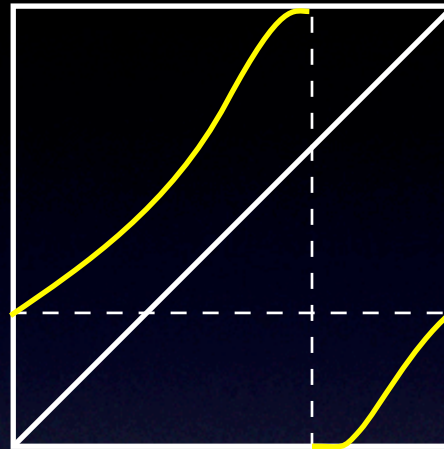


proper subintervals

-> Cantor set

Feigenbaum, Coulet-Tresser,
Lanford, H. Epstein,
Polynomial-like maps:
Douady-Hubbard, Sullivan,
McMullen, Lyubich

Circle map



partition of interval

Rand,
Khanin-Sinai,
de Faria, Yampolsky,
A. Epstein-Yamplosky

Sector/Near-parabolic



covering by sector or
croissant-like domains

gluing/identification
needed to define the
renormalization

Yoccoz, Perez-Marco,
Inou-S.

Irrationally indifferent fixed points

$$f(z) = e^{2\pi i\alpha}z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Irrationally indifferent fixed points

$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

Irrationally indifferent fixed points

$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

and beyond: boundary of linearization domain, invariant sets (hedgehogs)

Irrationally indifferent fixed points

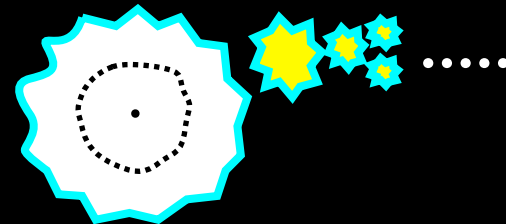
$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

and beyond: boundary of linearization domain, invariant sets (hedgehogs)
Siegel disk

Irrationally indifferent fixed points

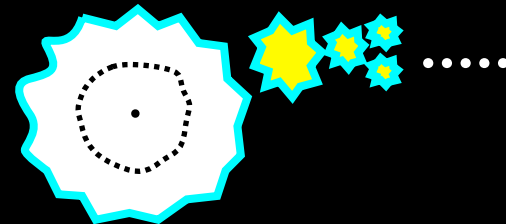
$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$



Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

and beyond: boundary of linearization domain, invariant sets (hedgehogs)
Siegel disk

Irrationally indifferent fixed points



$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

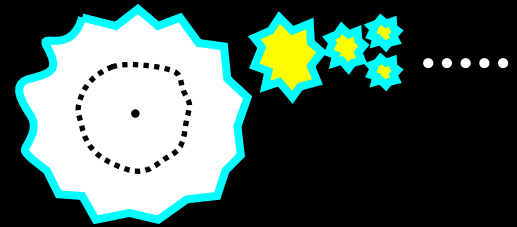
and beyond: boundary of linearization domain, invariant sets (hedgehogs)
Siegel disk

Siegel-Bruno Theorem

If α satisfies Bruno condition ($\sum \frac{\log q_{n+1}}{q_n} < \infty$ for the convergents p_n/q_n of α), then $f(z) = e^{2\pi i\alpha} z + \dots$ can be linearized.

(Yoccoz: the radius of convergence $> C \exp(-\sum \frac{\log q_{n+1}}{q_n})$ if f is univalent in $\{|z| < 1\}$.)

Irrationally indifferent fixed points



$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

and beyond: boundary of linearization domain, invariant sets (hedgehogs)
Siegel disk

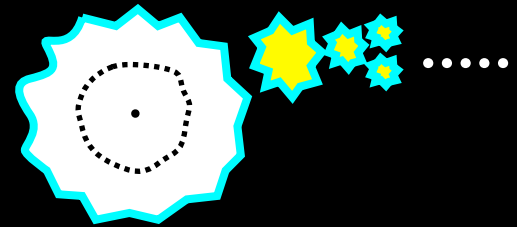
Siegel-Bruno Theorem

If α satisfies Bruno condition ($\sum \frac{\log q_{n+1}}{q_n} < \infty$ for the convergents p_n/q_n of α), then $f(z) = e^{2\pi i\alpha} z + \dots$ can be linearized.

(Yoccoz: the radius of convergence $> C \exp(-\sum \frac{\log q_{n+1}}{q_n})$ if f is univalent in $\{|z| < 1\}$.)

Proof by Yoccoz uses the renormalization

Irrationally indifferent fixed points



$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

and beyond: boundary of linearization domain, invariant sets (hedgehogs)
Siegel disk

Siegel-Bruno Theorem

If α satisfies Bruno condition ($\sum \frac{\log q_{n+1}}{q_n} < \infty$ for the convergents p_n/q_n of α), then $f(z) = e^{2\pi i\alpha} z + \dots$ can be linearized.

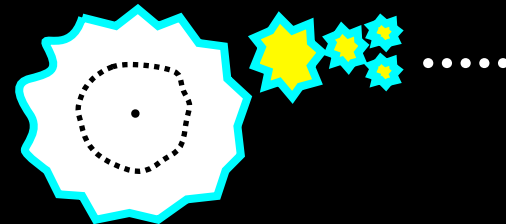
(Yoccoz: the radius of convergence $> C \exp(-\sum \frac{\log q_{n+1}}{q_n})$ if f is univalent in $\{|z| < 1\}$.)

Proof by Yoccoz uses the renormalization

Yoccoz Theorem

If α does not satisfy Bruno condition, then there exists $f(z) = e^{2\pi i\alpha} z + \dots$ which cannot be linearized. (In fact, $f(z) = e^{2\pi i\alpha} z + z^2$.)

Irrationally indifferent fixed points



$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

and beyond: boundary of linearization domain, invariant sets (hedgehogs)
Siegel disk

Siegel-Bruno Theorem

If α satisfies Bruno condition ($\sum \frac{\log q_{n+1}}{q_n} < \infty$ for the convergents p_n/q_n of α), then $f(z) = e^{2\pi i\alpha} z + \dots$ can be linearized.

(Yoccoz: the radius of convergence $> C \exp(-\sum \frac{\log q_{n+1}}{q_n})$ if f is univalent in $\{|z| < 1\}$.)

Proof by Yoccoz uses the renormalization

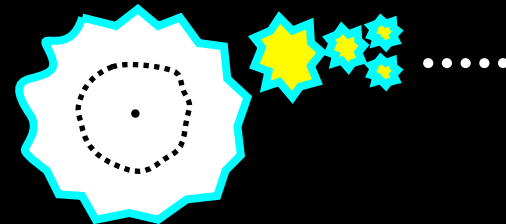
Yoccoz Theorem

If α does not satisfy Bruno condition, then there exists $f(z) = e^{2\pi i\alpha} z + \dots$ which cannot be linearized. (In fact, $f(z) = e^{2\pi i\alpha} z + z^2$.)

Yoccoz renormalization

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$

Irrationally indifferent fixed points



$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

and beyond: boundary of linearization domain, invariant sets (hedgehogs)
Siegel disk

Siegel-Bruno Theorem

If α satisfies Bruno condition ($\sum \frac{\log q_{n+1}}{q_n} < \infty$ for the convergents p_n/q_n of α), then $f(z) = e^{2\pi i\alpha} z + \dots$ can be linearized.

(Yoccoz: the radius of convergence $> C \exp(-\sum \frac{\log q_{n+1}}{q_n})$ if f is univalent in $\{|z| < 1\}$.)

Proof by Yoccoz uses the renormalization

Yoccoz Theorem

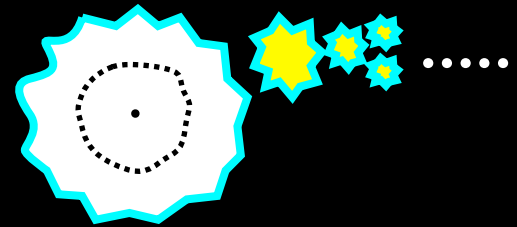
If α does not satisfy Bruno condition, then there exists $f(z) = e^{2\pi i\alpha} z + \dots$ which cannot be linearized. (In fact, $f(z) = e^{2\pi i\alpha} z + z^2$.)

Yoccoz renormalization

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$

$$\alpha_{n+1} = \text{fractional part of } \frac{1}{\alpha_n}$$

Irrationally indifferent fixed points



$$f(z) = e^{2\pi i\alpha} z + \dots, \quad \alpha \in \mathbb{R} \setminus \mathbb{Q}$$

Linearization: local conjugacy to its linear part $z \mapsto e^{2\pi i\alpha} z$

and beyond: boundary of linearization domain, invariant sets (hedgehogs)
Siegel disk

Siegel-Bruno Theorem

If α satisfies Bruno condition ($\sum \frac{\log q_{n+1}}{q_n} < \infty$ for the convergents p_n/q_n of α), then $f(z) = e^{2\pi i\alpha} z + \dots$ can be linearized.

(Yoccoz: the radius of convergence $> C \exp(-\sum \frac{\log q_{n+1}}{q_n})$ if f is univalent in $\{|z| < 1\}$.)

Proof by Yoccoz uses the renormalization

Yoccoz Theorem

If α does not satisfy Bruno condition, then there exists $f(z) = e^{2\pi i\alpha} z + \dots$ which cannot be linearized. (In fact, $f(z) = e^{2\pi i\alpha} z + z^2$.)

Yoccoz renormalization

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$

$\alpha_{n+1} =$ fractional part of $\frac{1}{\alpha_n}$ Gauss map for continued fractions

Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$

f_n



Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$

f_n



Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

f_n



Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$

f_n



Yoccoz renormalization for Siegel-Bruno Theorem

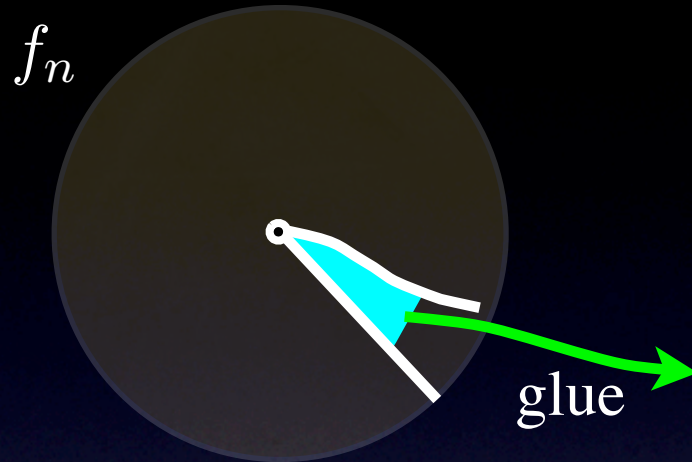
$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$

f_n



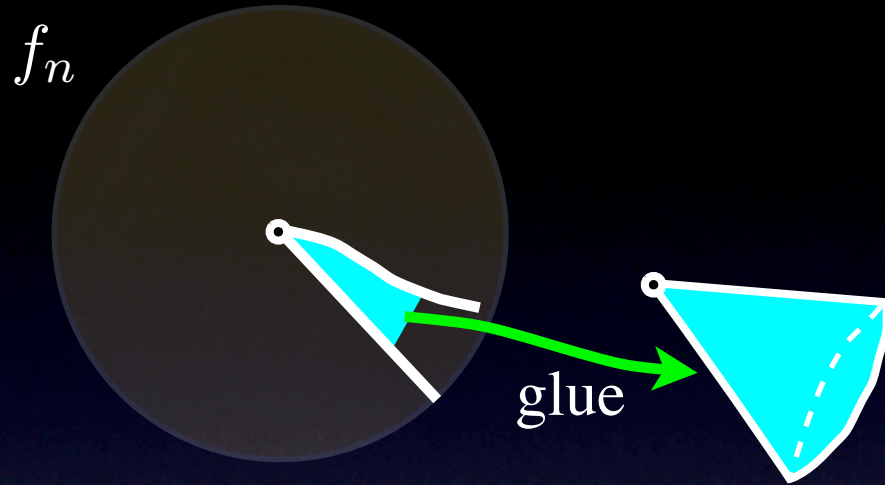
Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$



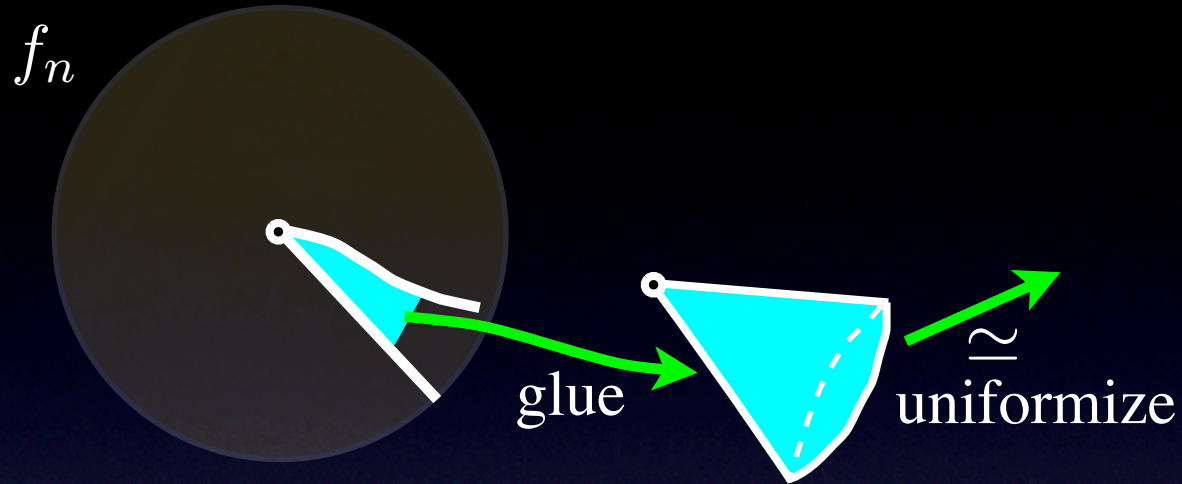
Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$



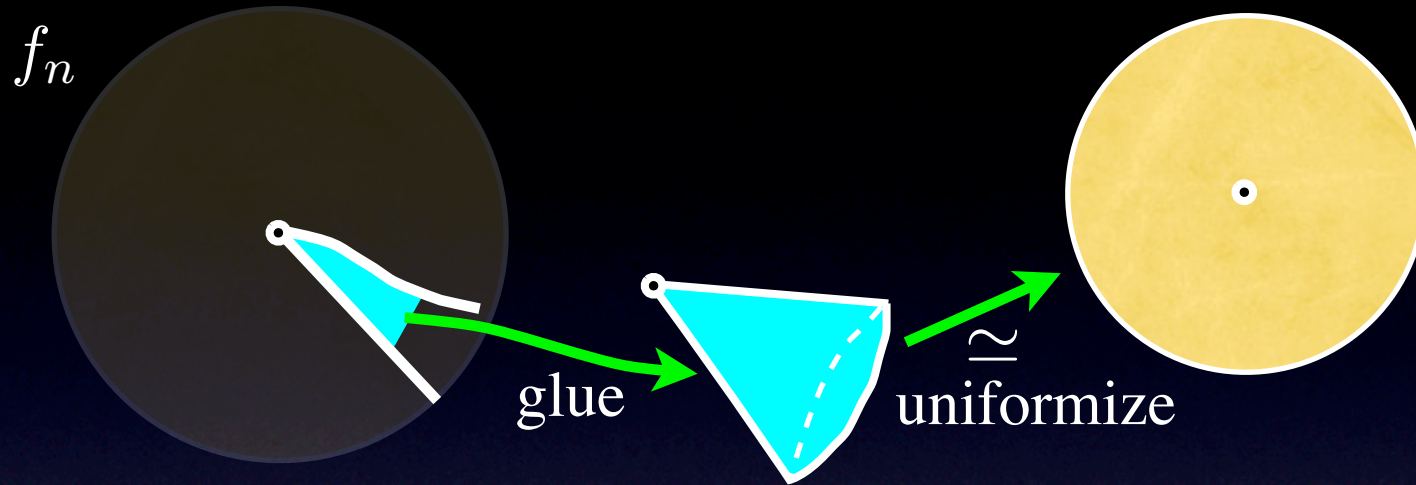
Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$



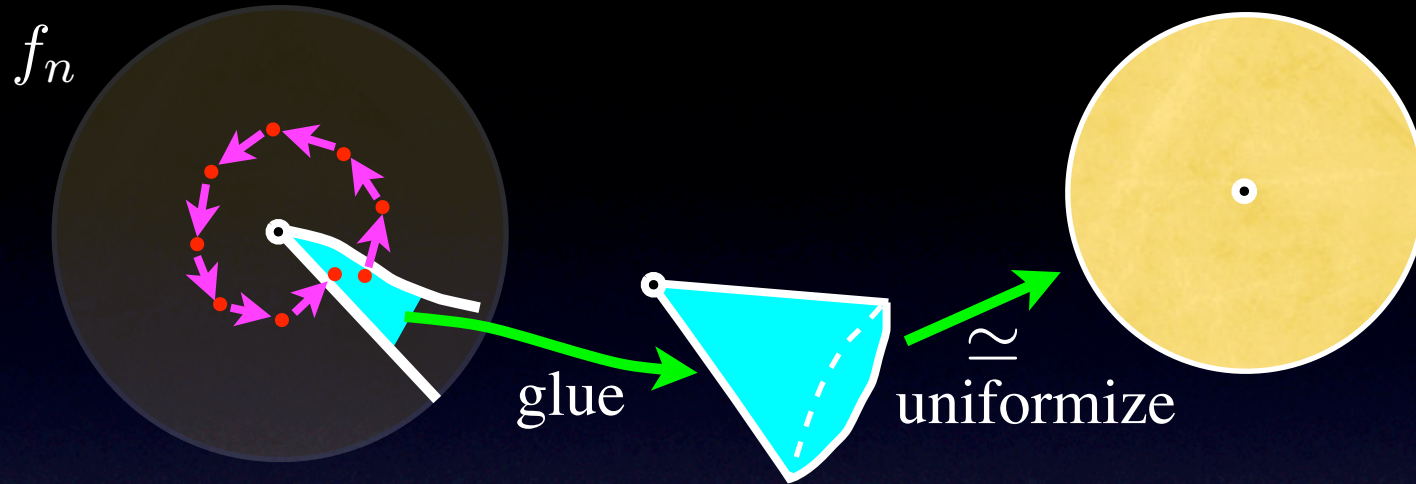
Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$



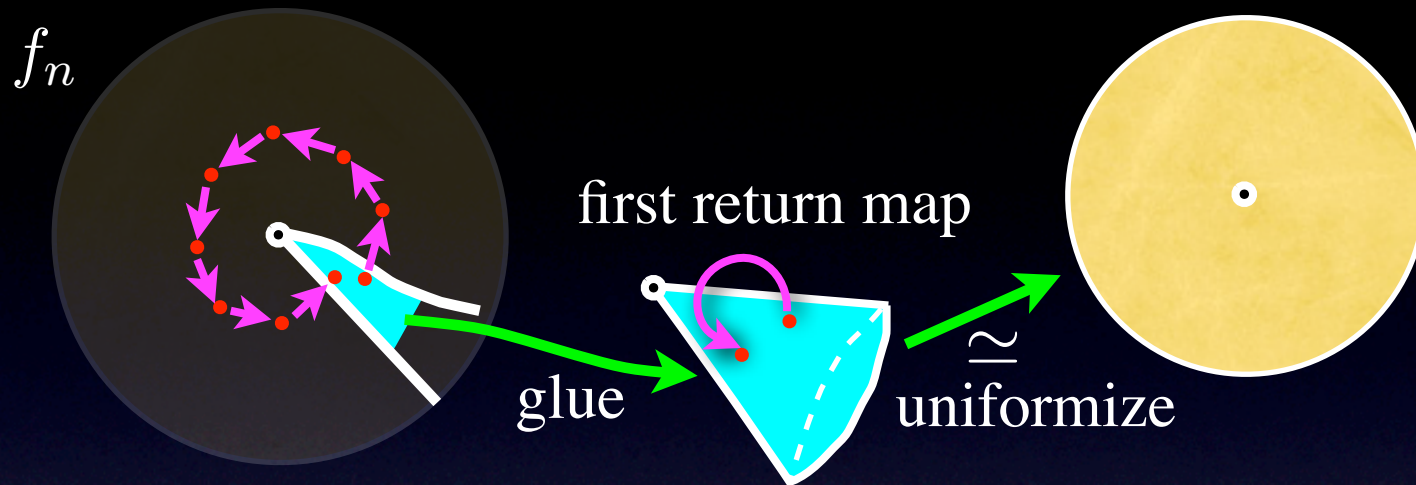
Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$



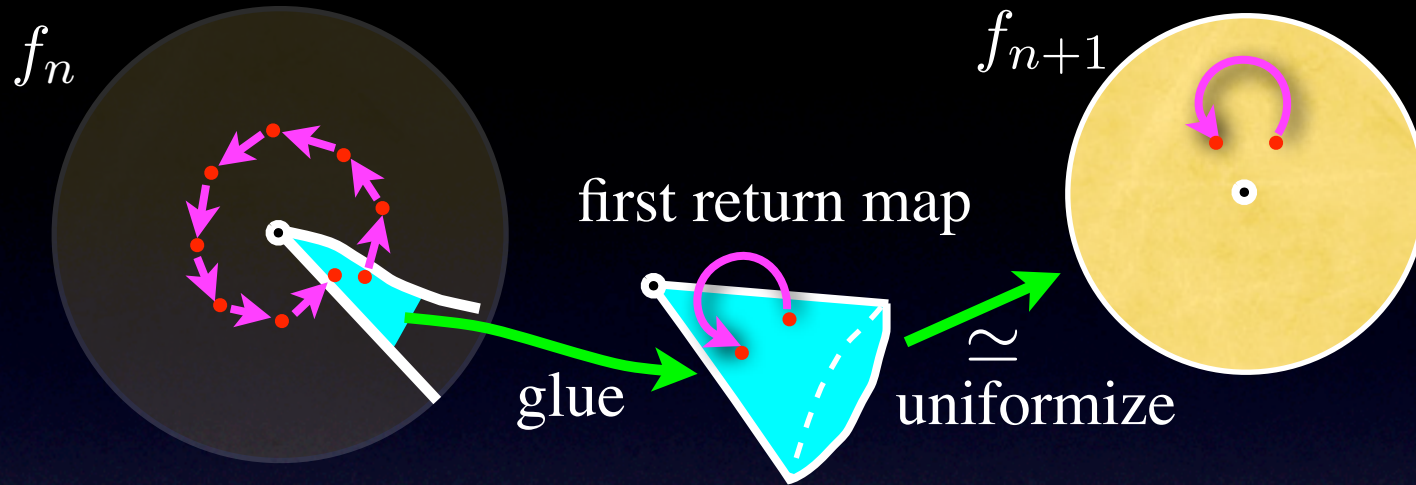
Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$



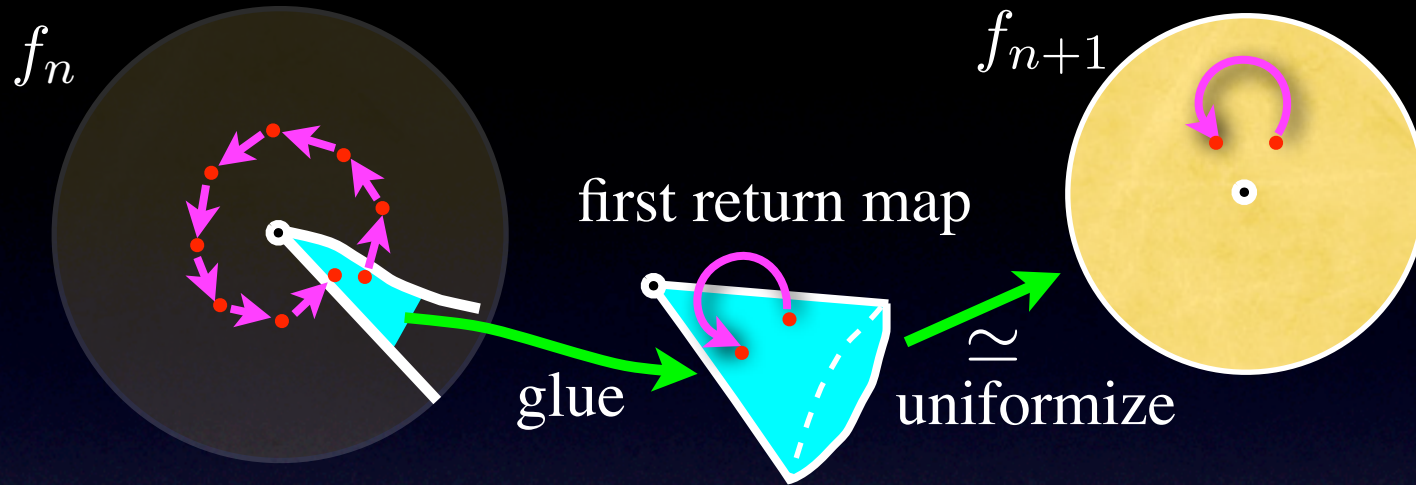
Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$



Yoccoz renormalization for Siegel-Bruno Theorem

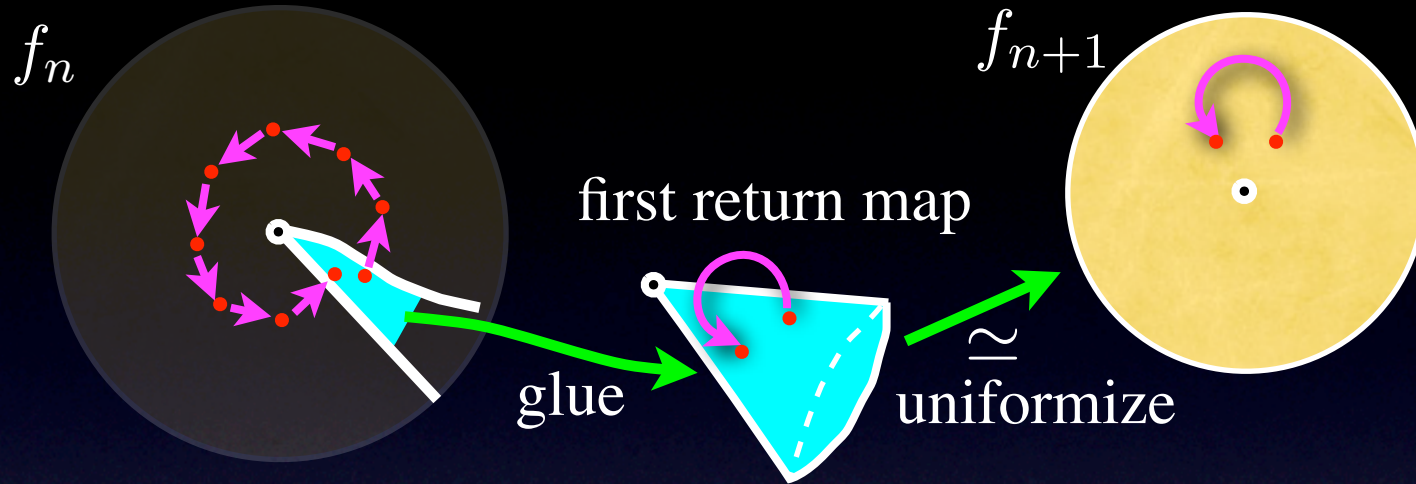
$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$



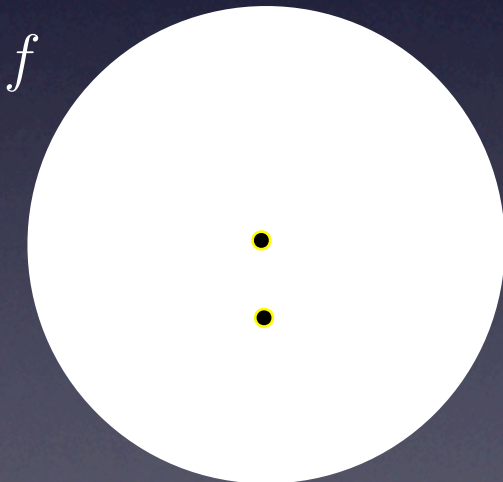
Cylinder/Near-parabolic renormalization

Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

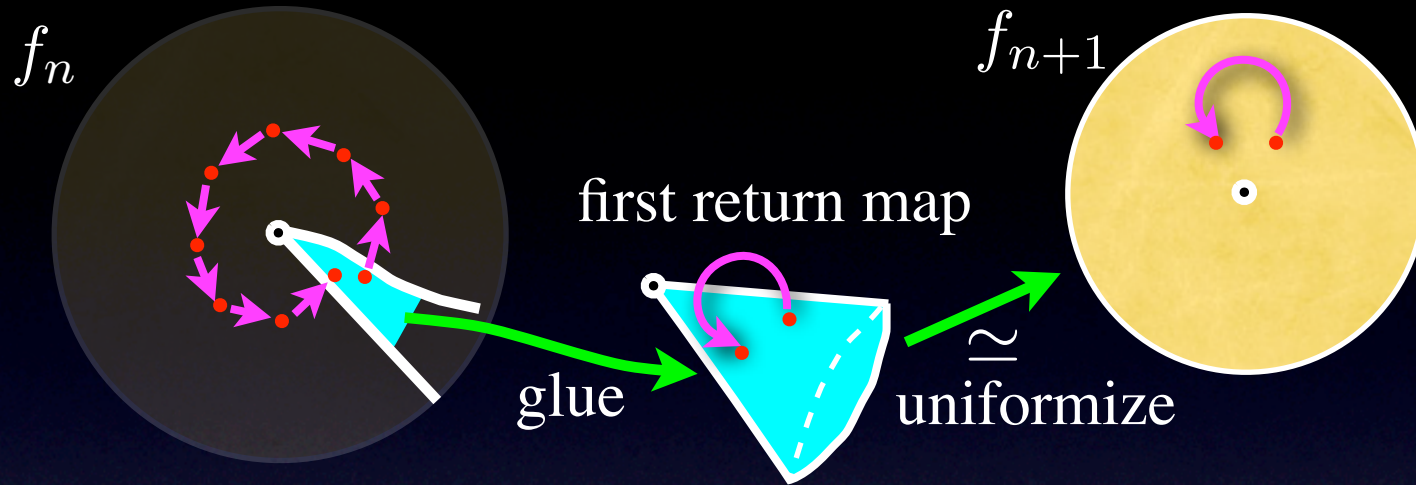


Cylinder/Near-parabolic renormalization

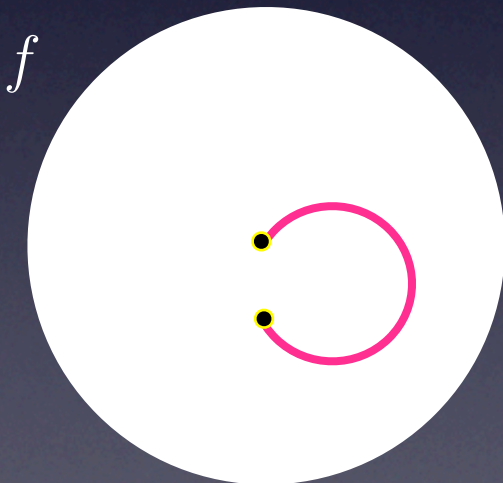


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

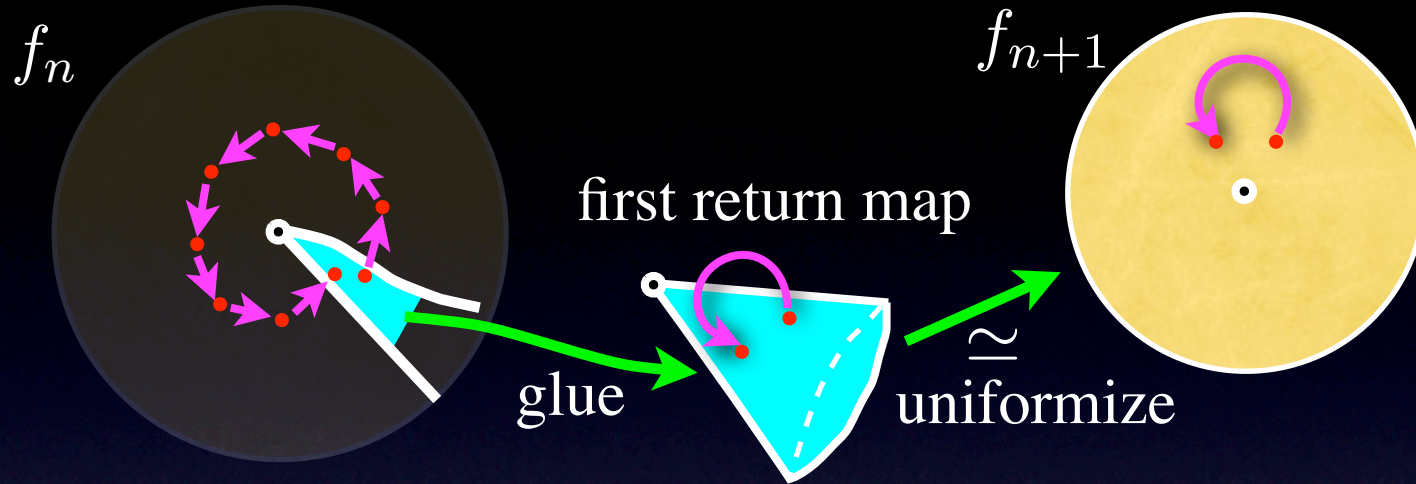


Cylinder/Near-parabolic renormalization

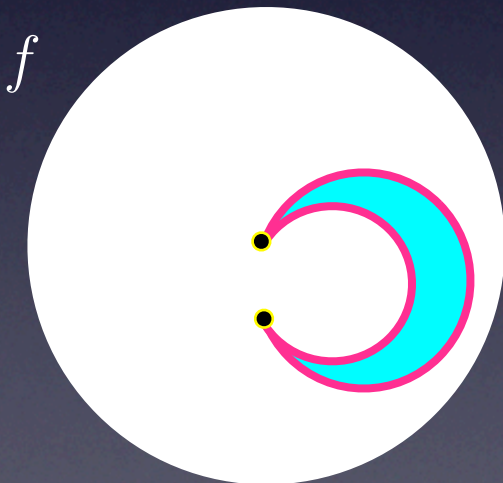


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

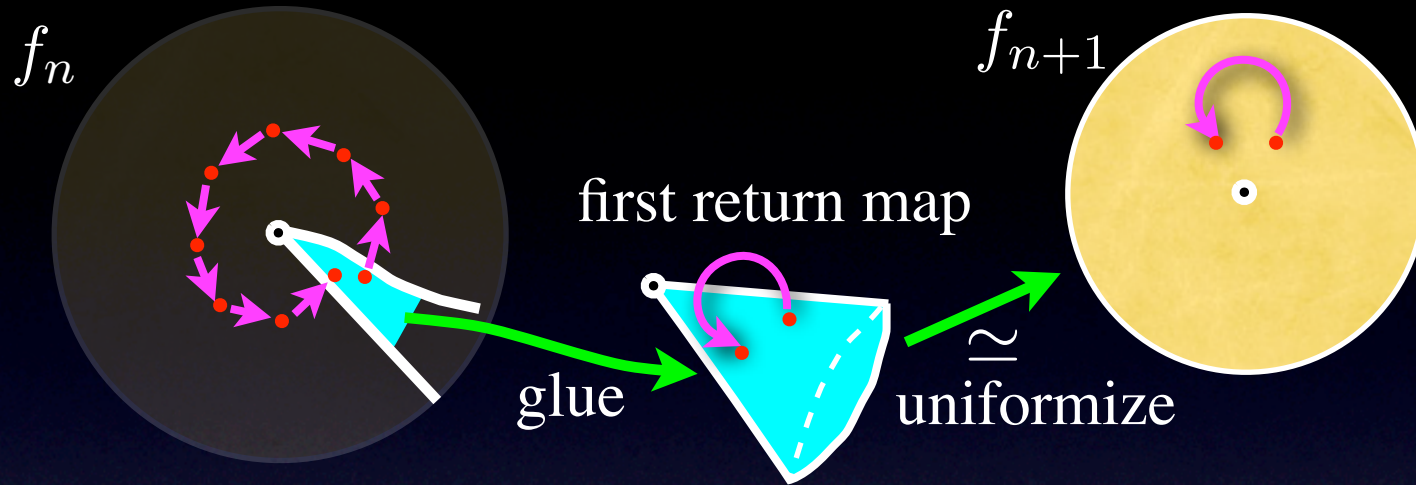


Cylinder/Near-parabolic renormalization

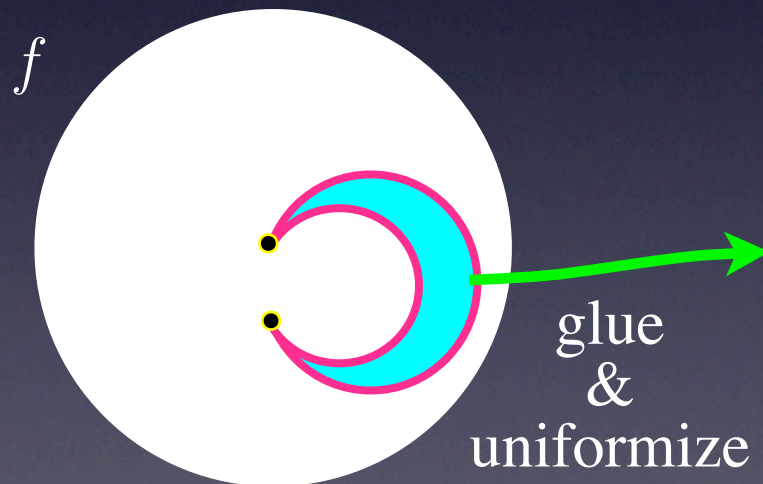


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i\alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i\alpha_{n+1}} z + \dots$$

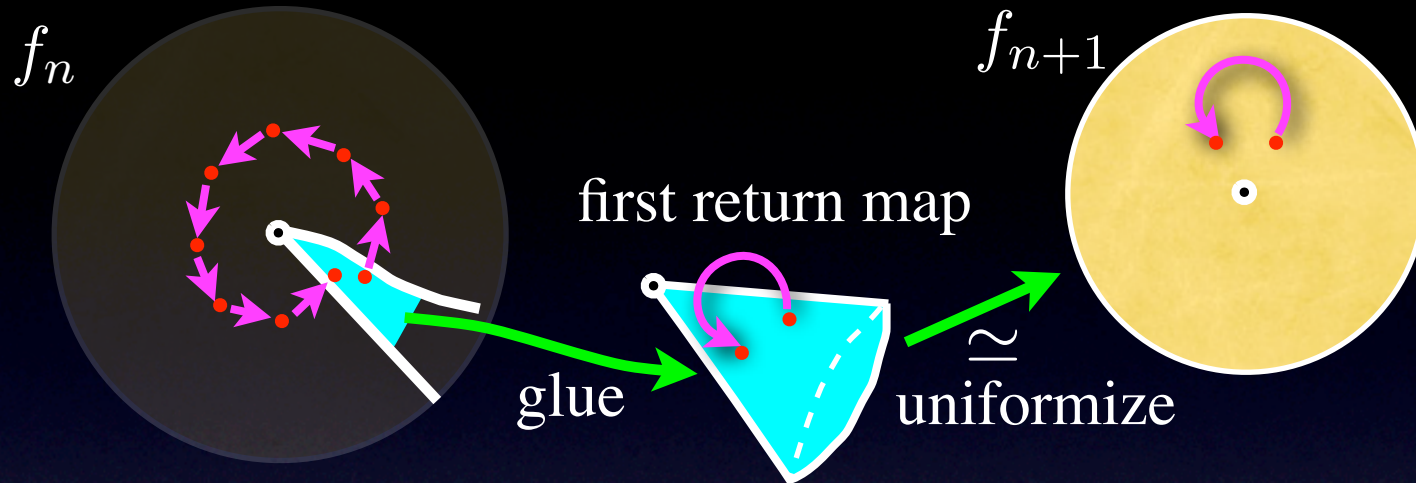


Cylinder/Near-parabolic renormalization

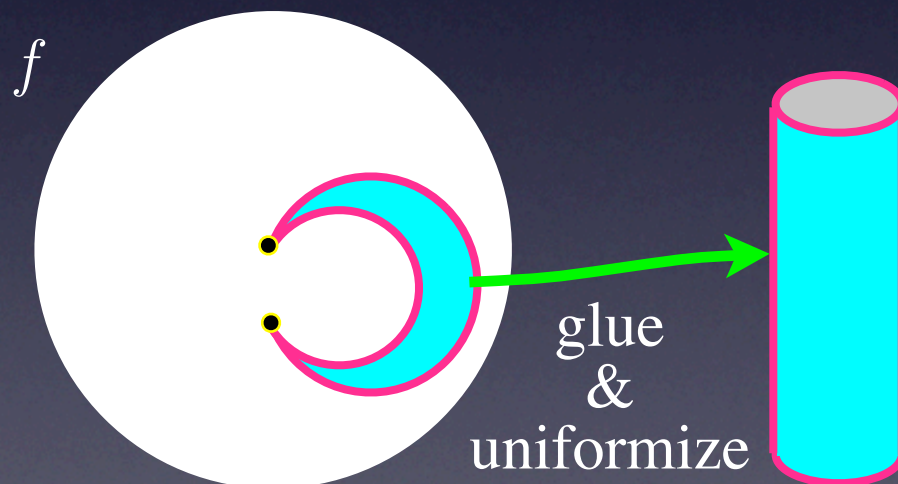


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

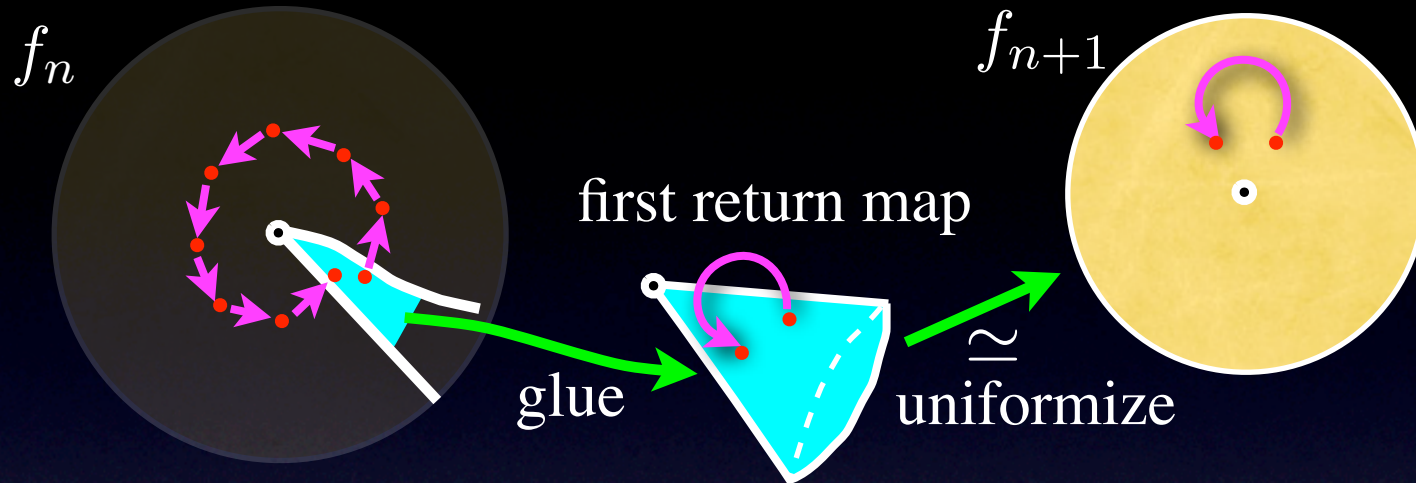


Cylinder/Near-parabolic renormalization

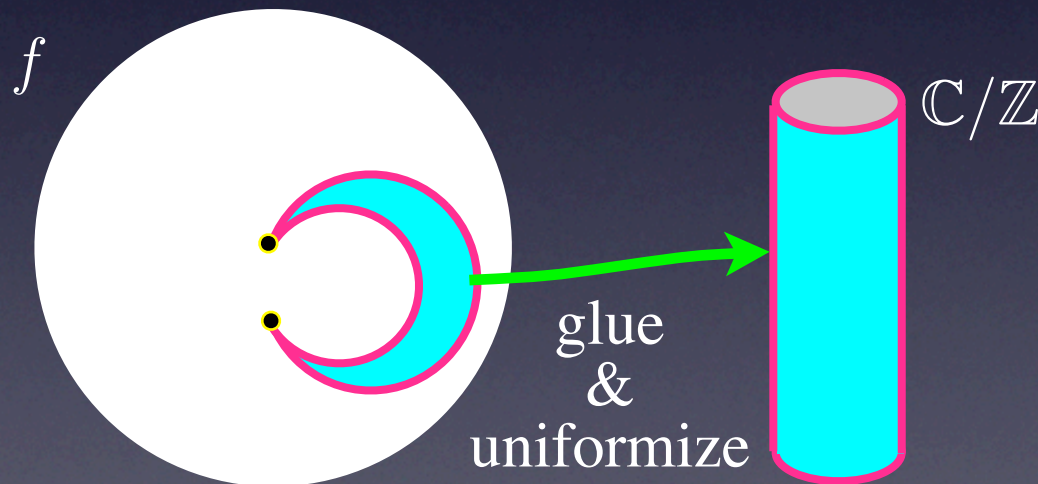


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

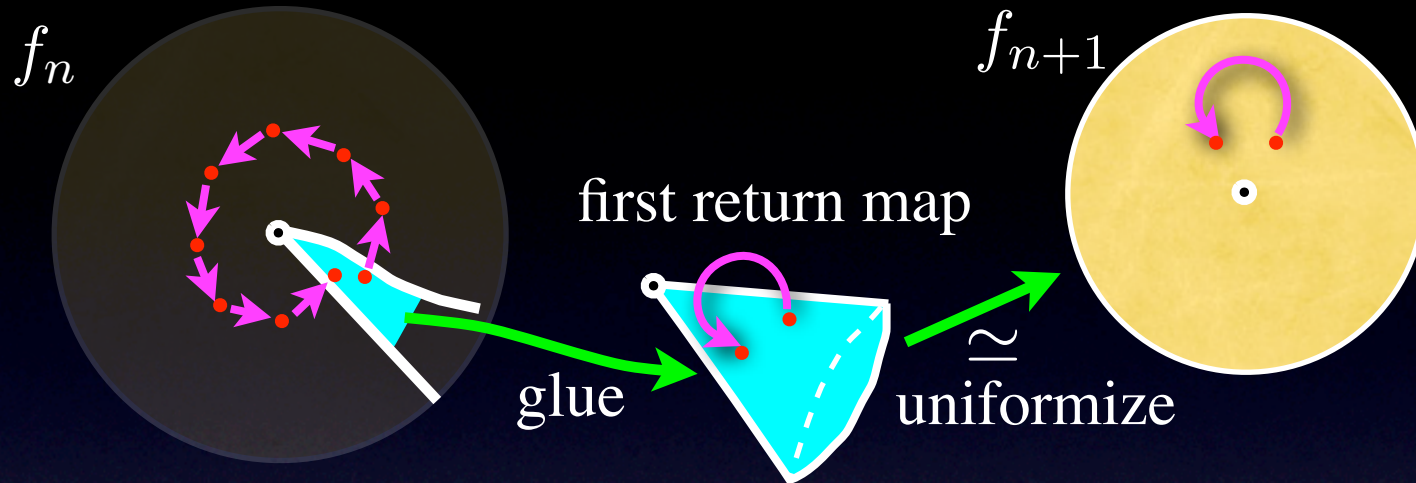


Cylinder/Near-parabolic renormalization

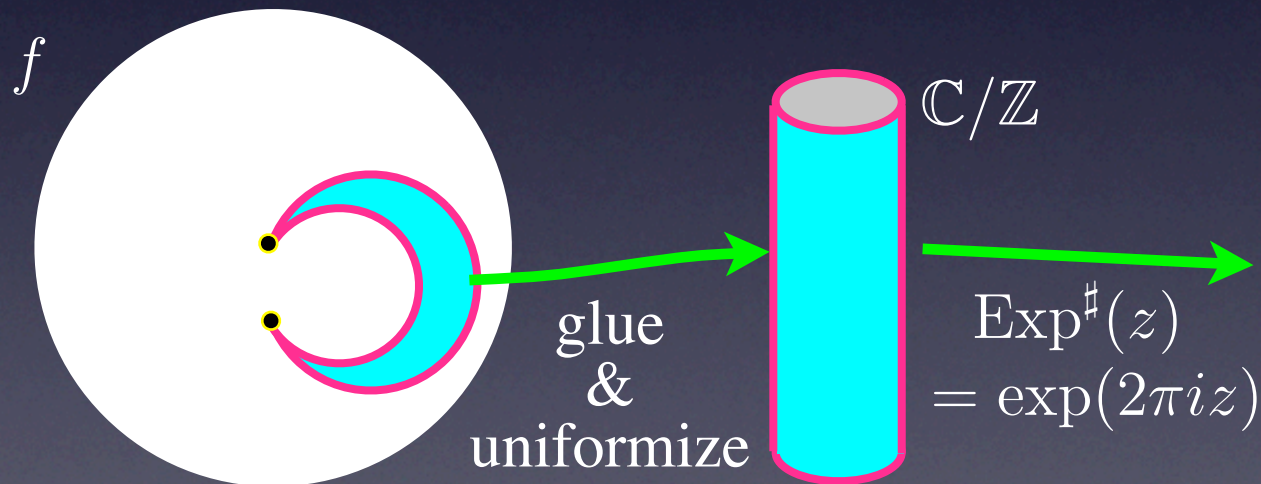


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

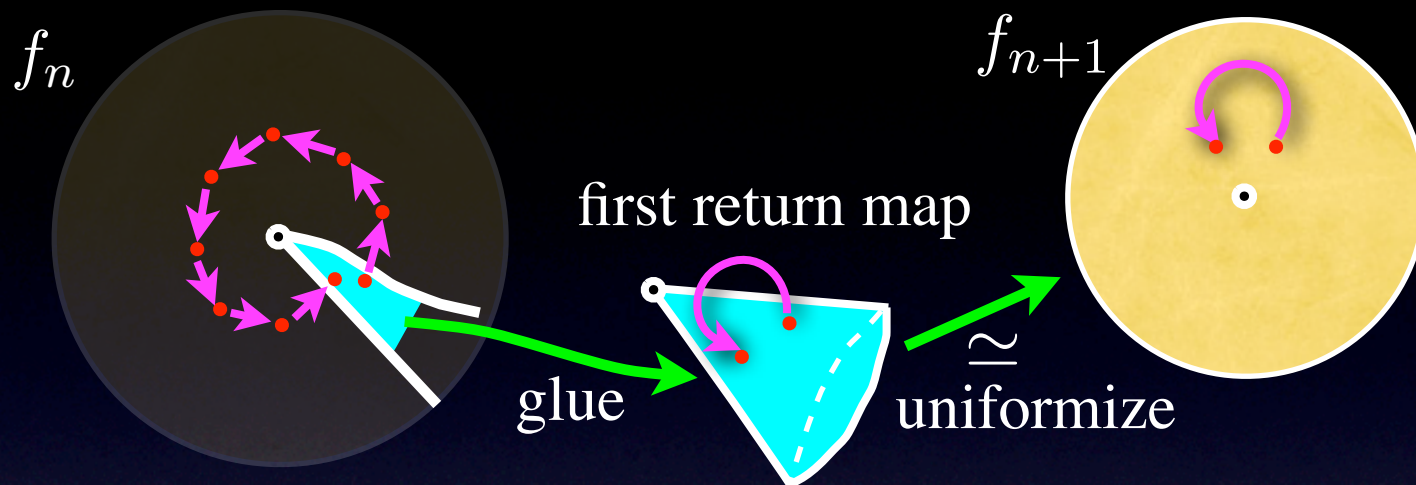


Cylinder/Near-parabolic renormalization

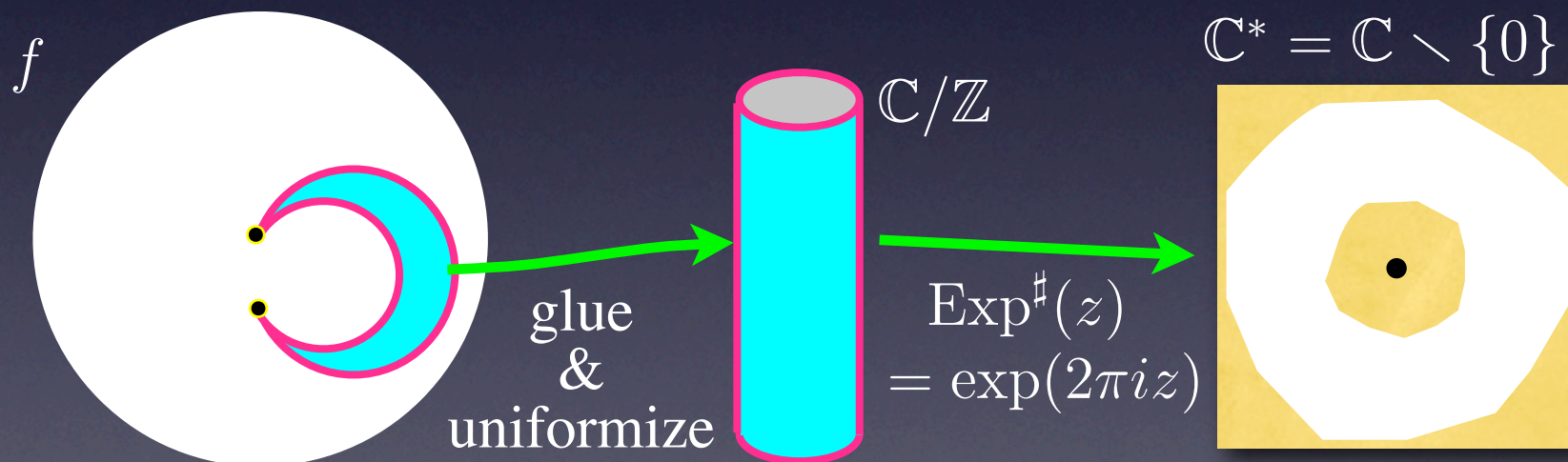


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

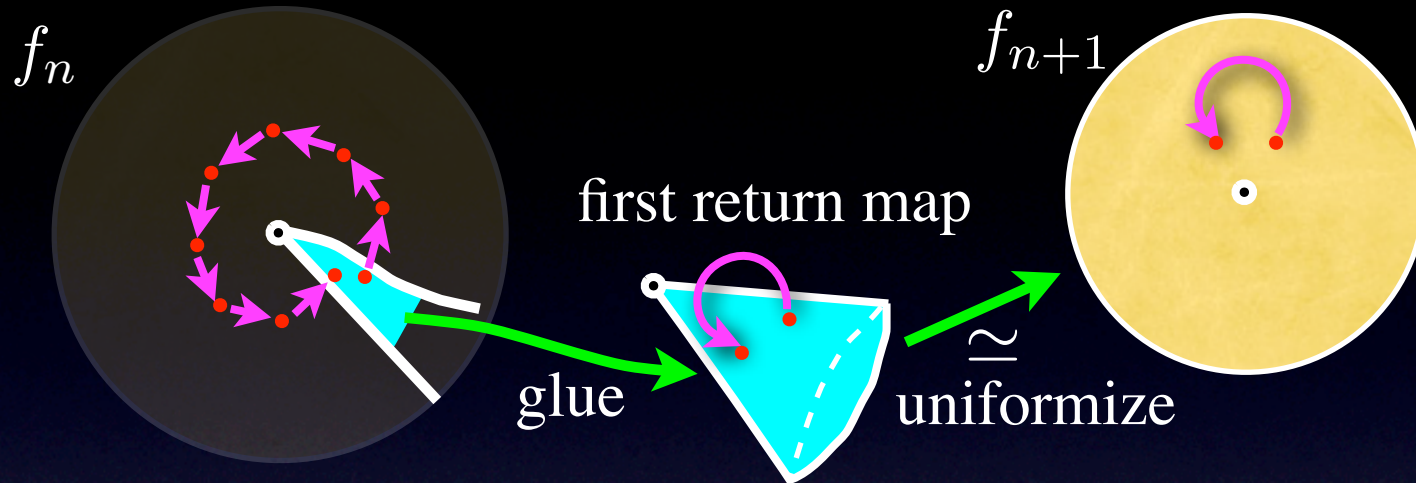


Cylinder/Near-parabolic renormalization

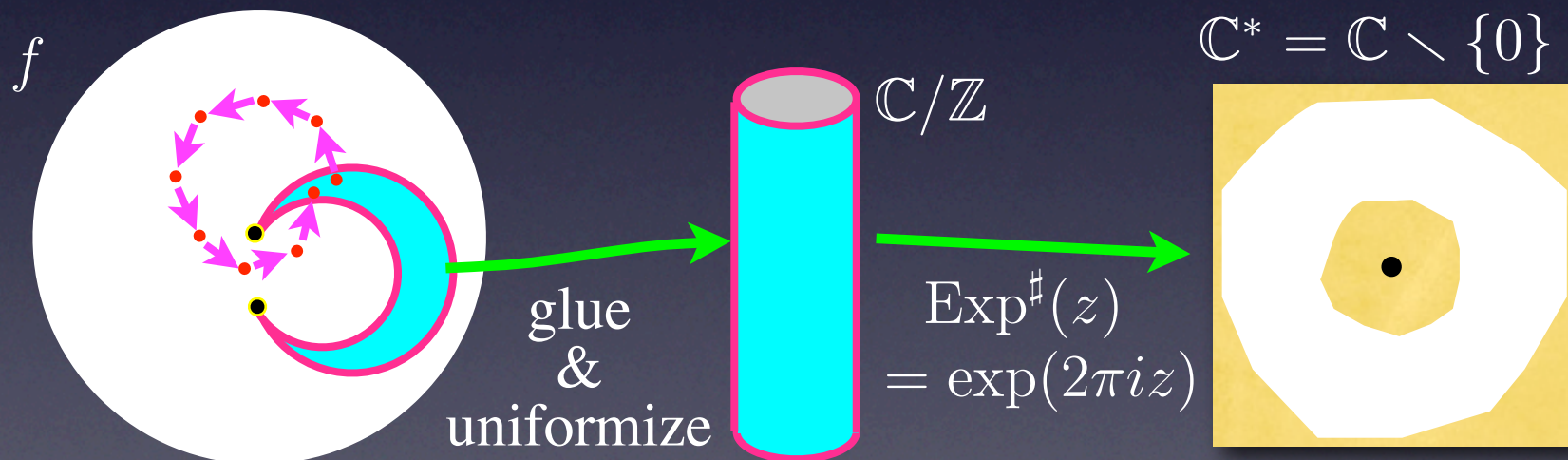


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

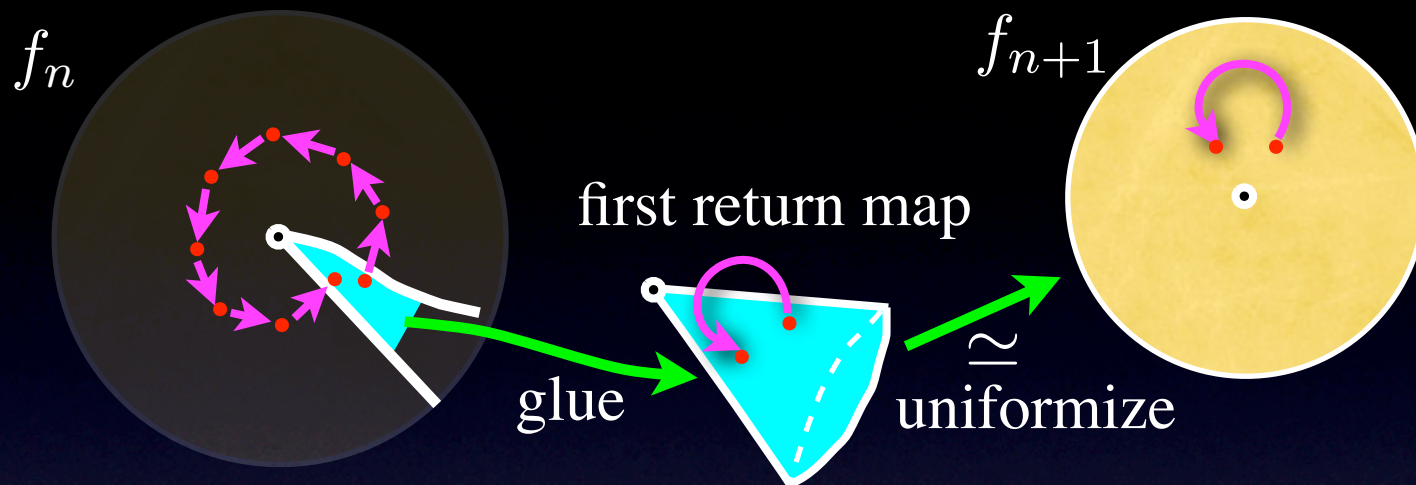


Cylinder/Near-parabolic renormalization

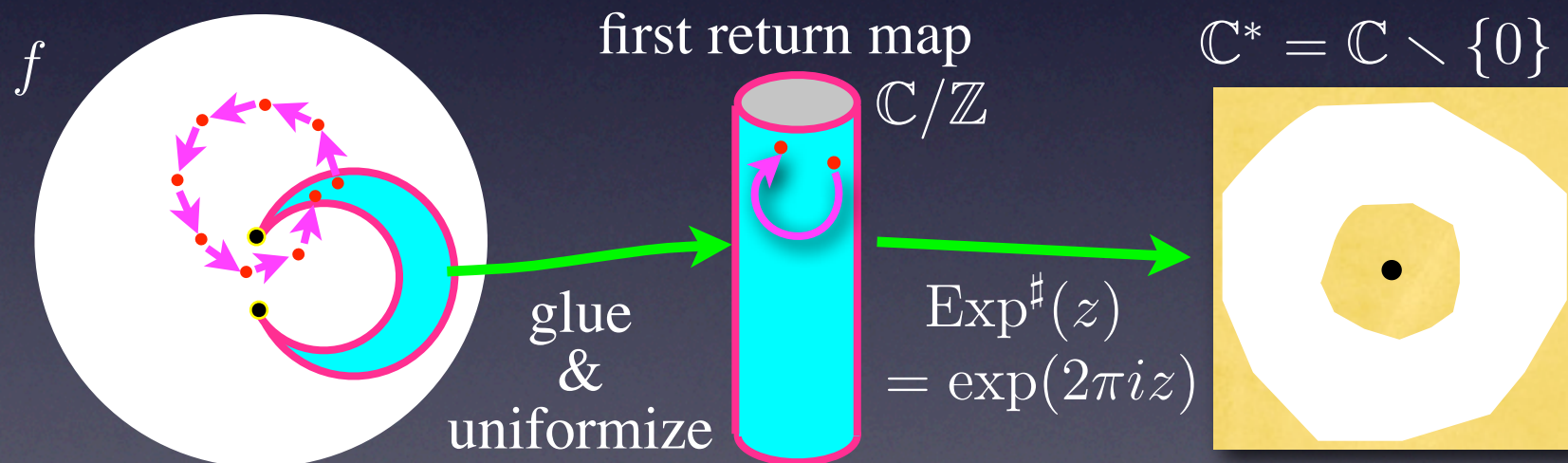


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

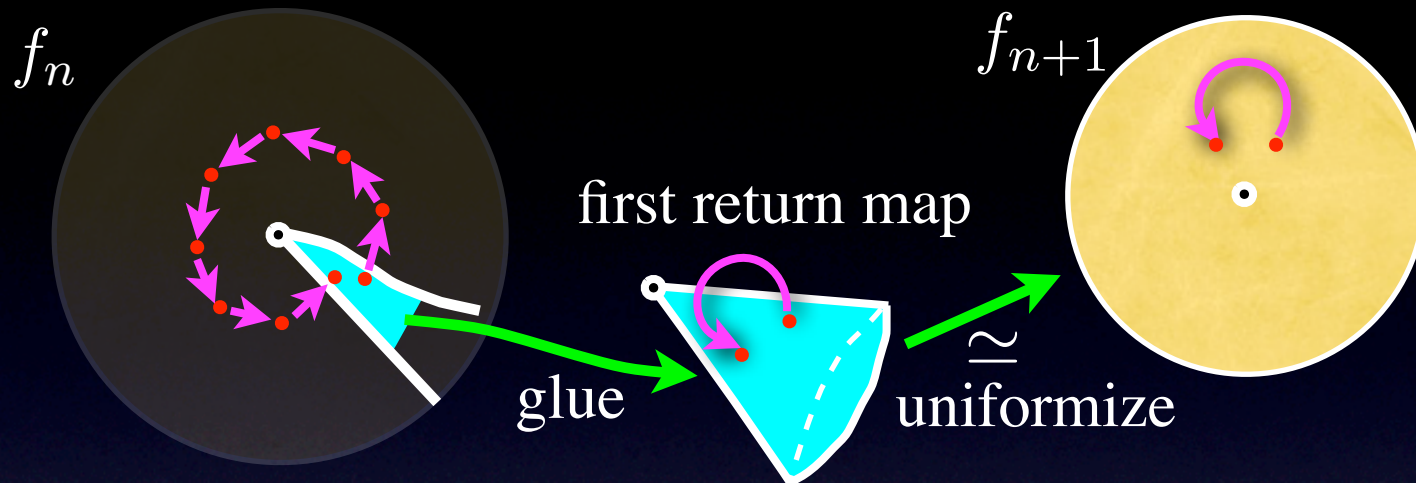


Cylinder/Near-parabolic renormalization

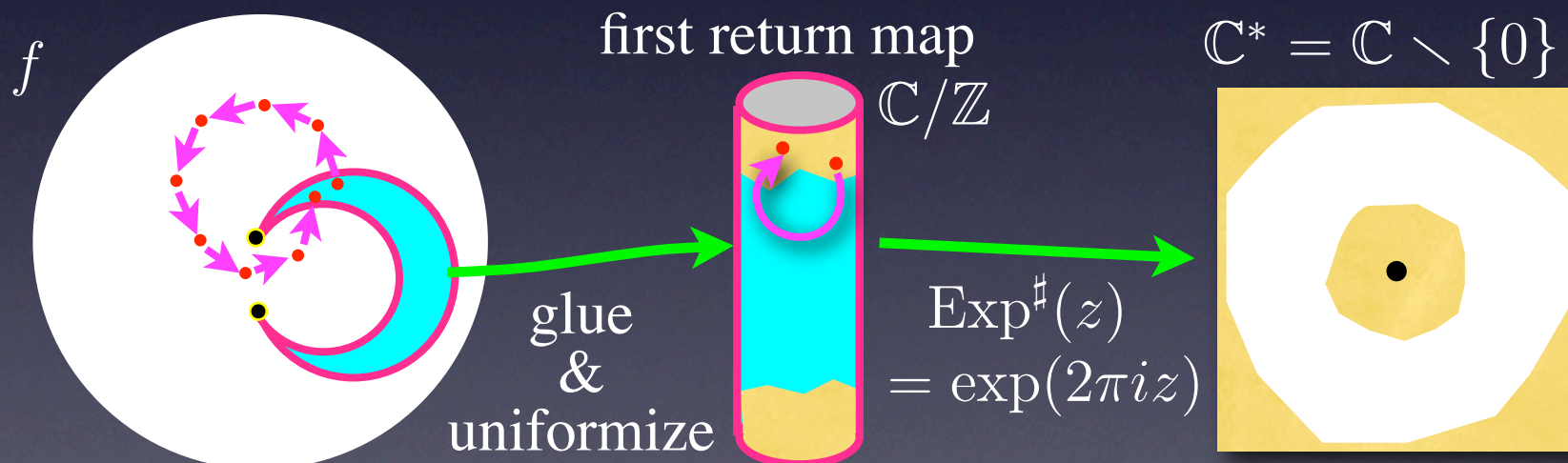


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

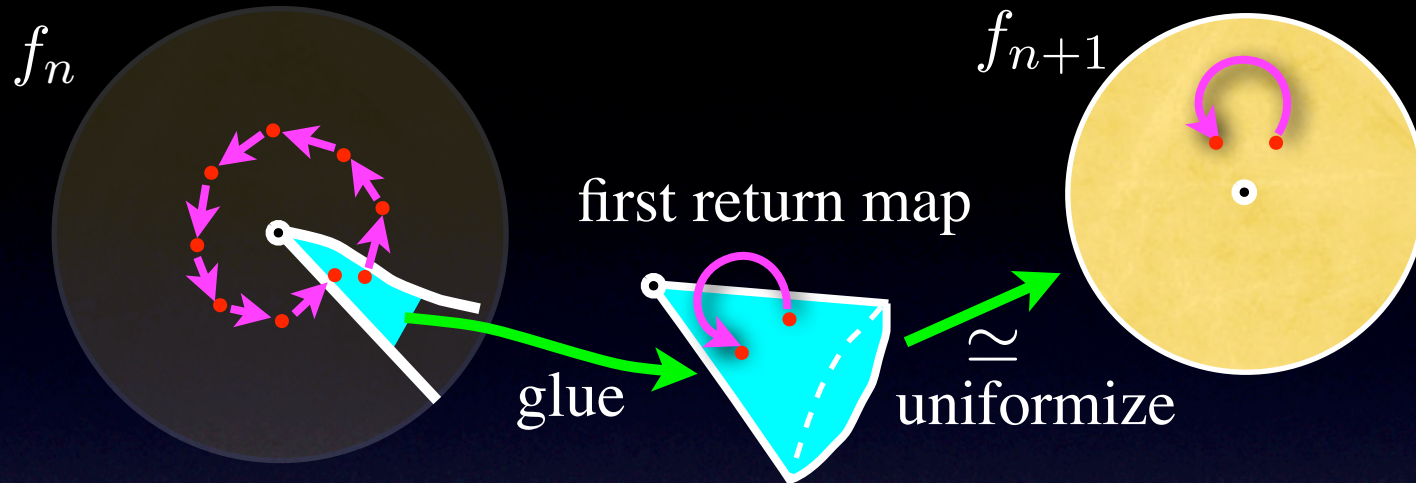


Cylinder/Near-parabolic renormalization

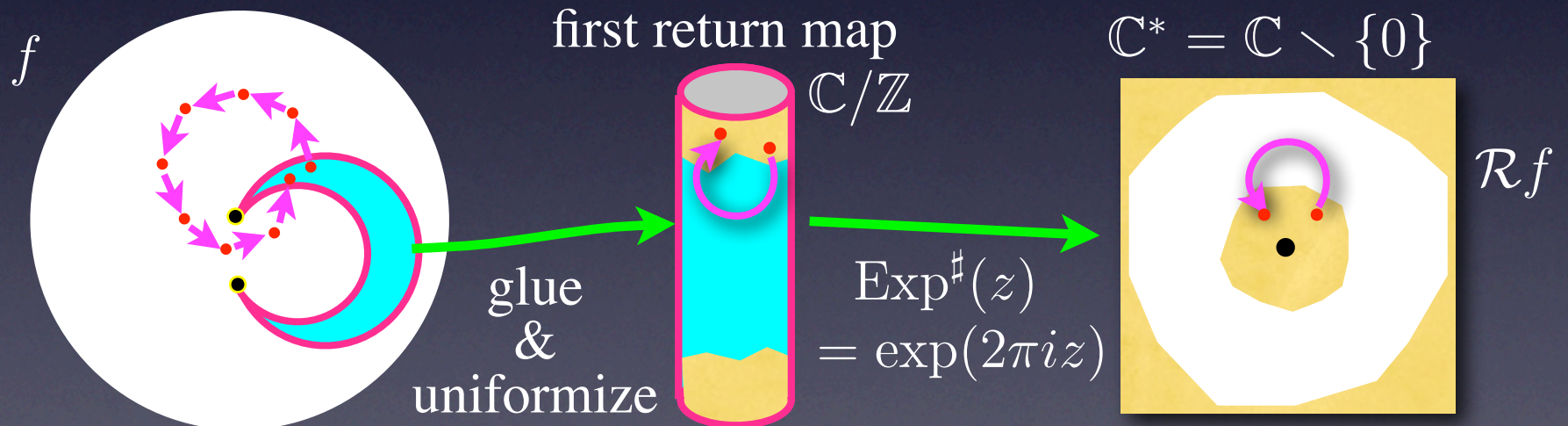


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$

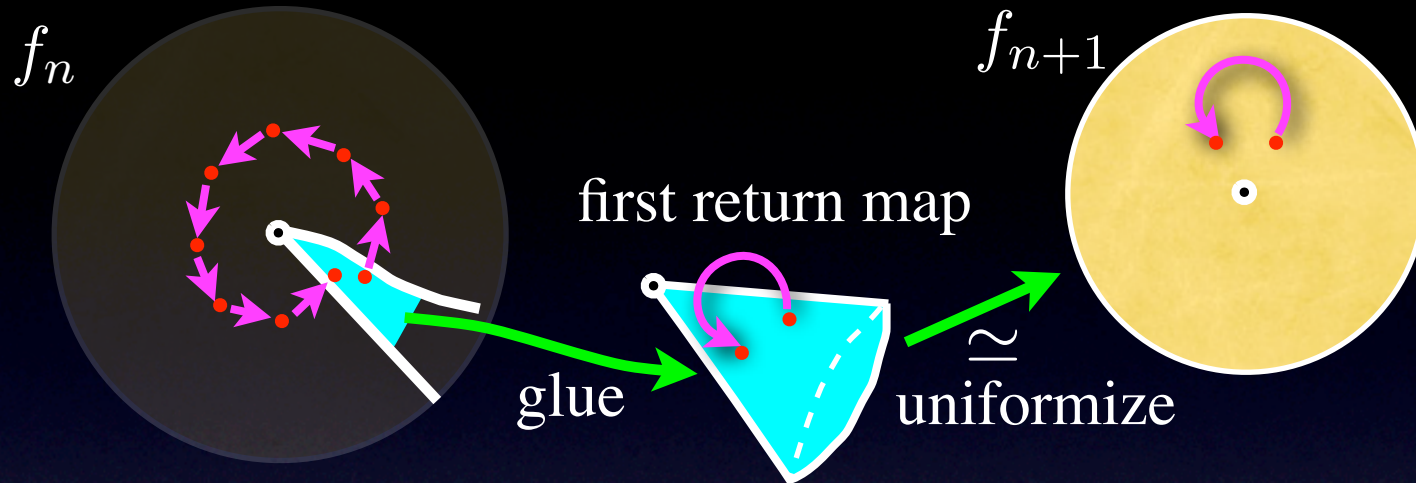


Cylinder/Near-parabolic renormalization

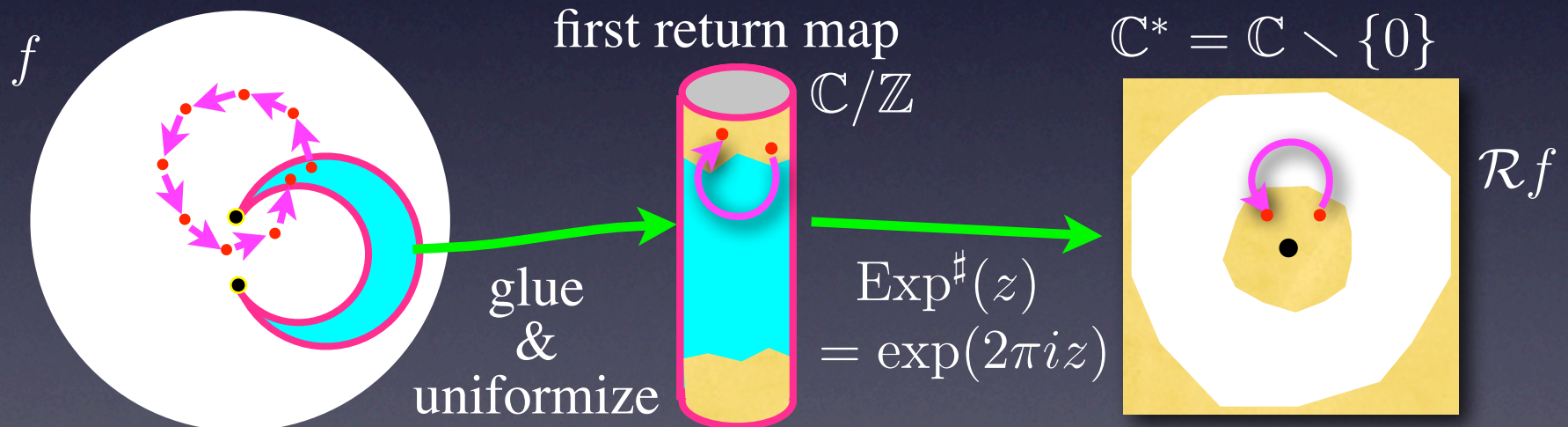


Yoccoz renormalization for Siegel-Bruno Theorem

$$f_n(z) = e^{2\pi i \alpha_n} z + \dots \rightsquigarrow f_{n+1}(z) = e^{2\pi i \alpha_{n+1}} z + \dots$$



Cylinder/Near-parabolic renormalization



$\mathcal{R}f$ can be defined when $f(z) = e^{2\pi i \alpha} z + \dots$ is a small perturbation of $z + a_2 z^2 + \dots$ ($a_2 \neq 0$) and $|\arg \alpha| < \pi/4$.

For $N \in \mathbb{N}$, let $Irrat_N$ be the set of irrational number of high type:

$$Irrat_N \ni \alpha = \pm \frac{1}{a_1 \pm \frac{1}{a_2 \pm \frac{1}{\ddots}}} \quad \text{where } a_i \in \mathbb{N} \text{ and } a_i \geq N,$$

For $N \in \mathbb{N}$, let $Irrat_N$ be the set of irrational number of high type:

$$Irrat_N \ni \alpha = \pm \frac{1}{a_1 \pm \frac{1}{a_2 \pm \frac{1}{\ddots}}} \quad \text{where } a_i \in \mathbb{N} \text{ and } a_i \geq N,$$

For a neighborhood V of 0, define $P(z) = z(1+z)^2$ and

$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \left| \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right. \right\}$$

For $N \in \mathbb{N}$, let $Irrat_N$ be the set of irrational number of high type:

$$Irrat_N \ni \alpha = \pm \frac{1}{a_1 \pm \frac{1}{a_2 \pm \frac{1}{\ddots}}} \quad \text{where } a_i \in \mathbb{N} \text{ and } a_i \geq N,$$

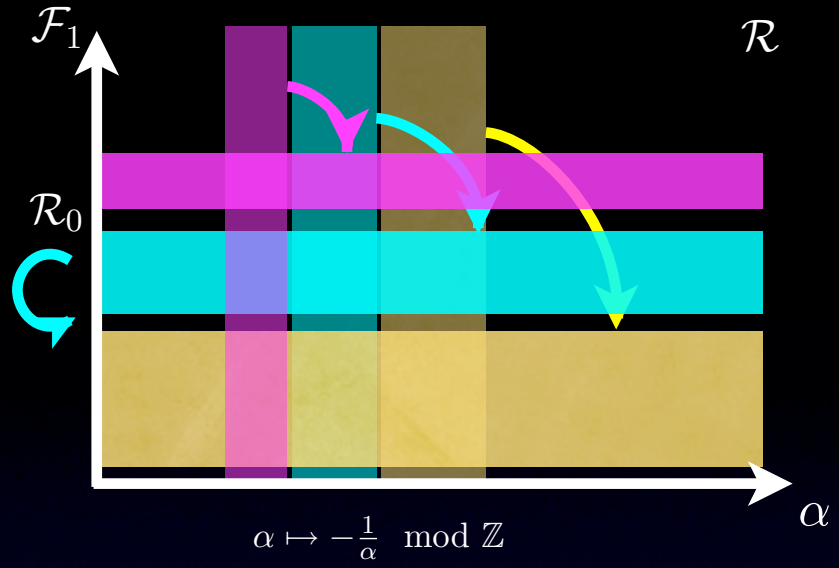
For a neighborhood V of 0, define $P(z) = z(1+z)^2$ and

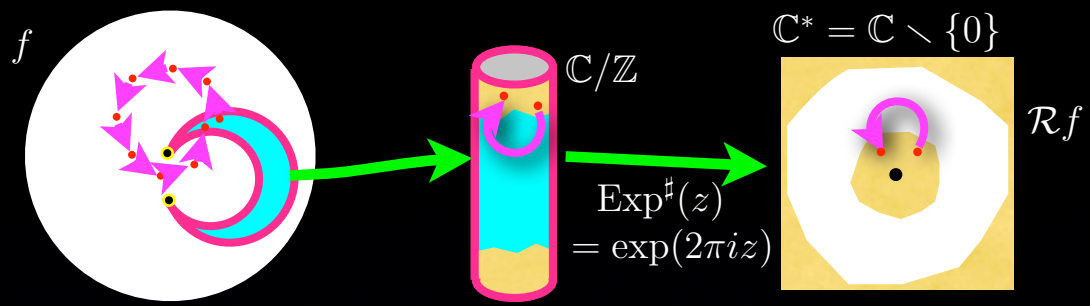
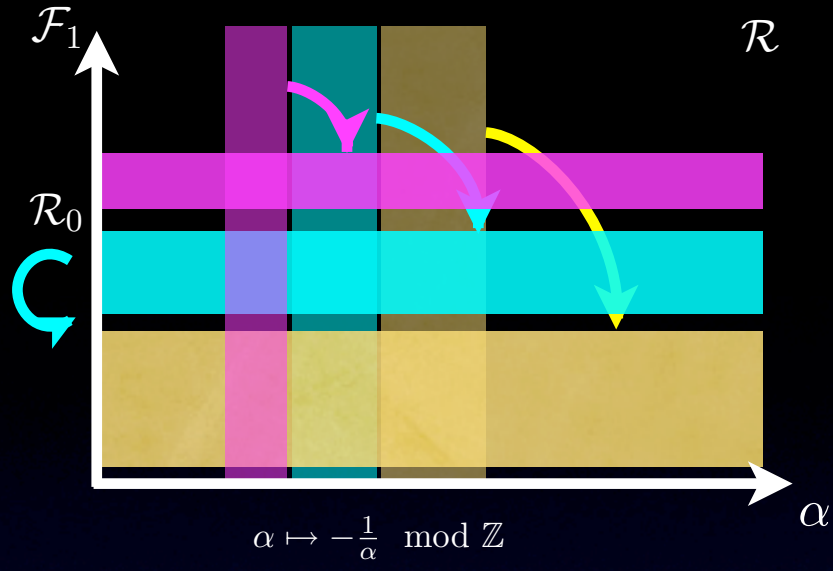
$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \left| \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right. \right\}$$

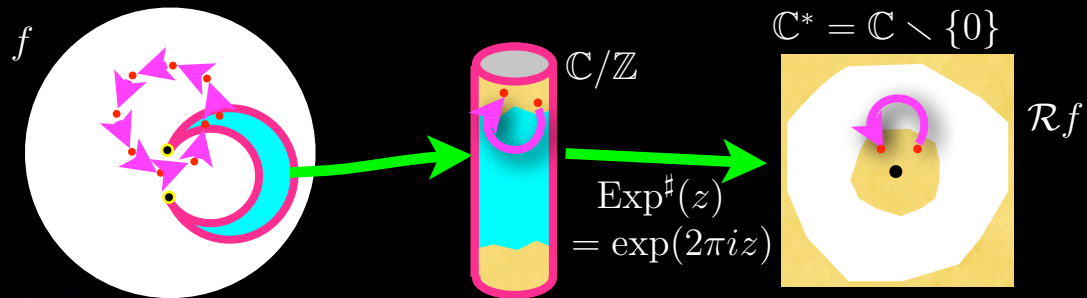
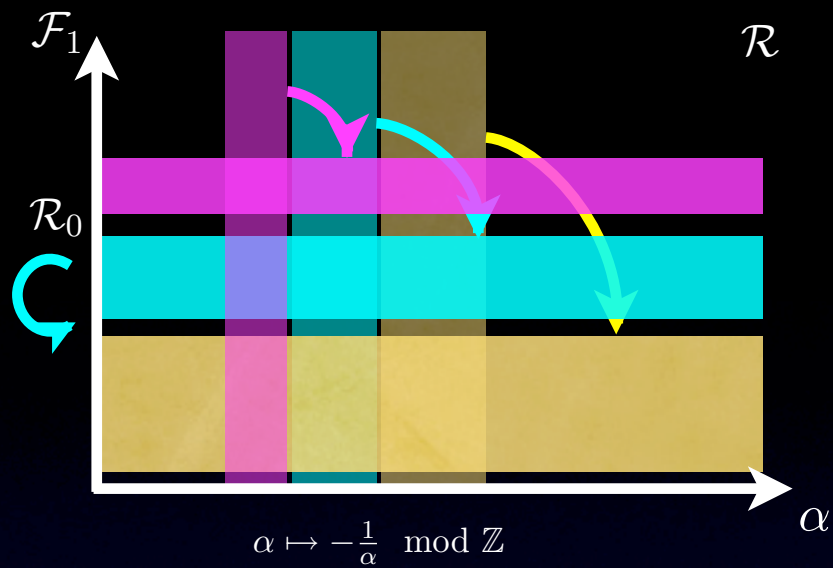
Theorem (Inou & S.): For some V and N , the near-parabolic renormalization \mathcal{R} from

$$\{e^{2\pi i\alpha} f : \alpha \in Irrat_N, f \in \mathcal{F}_1\} = Irrat_N \times \mathcal{F}_1$$

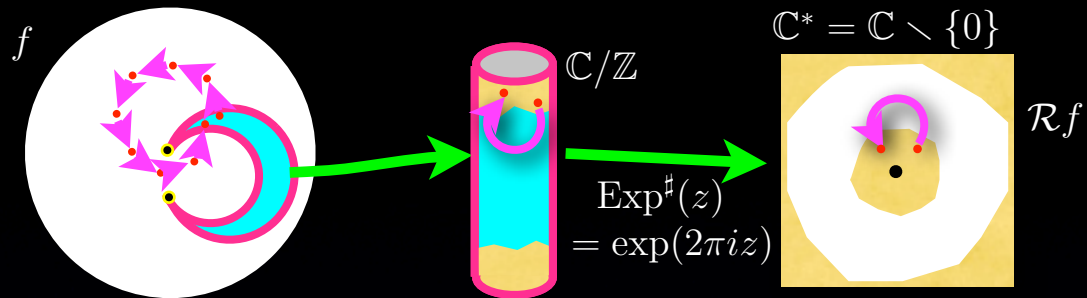
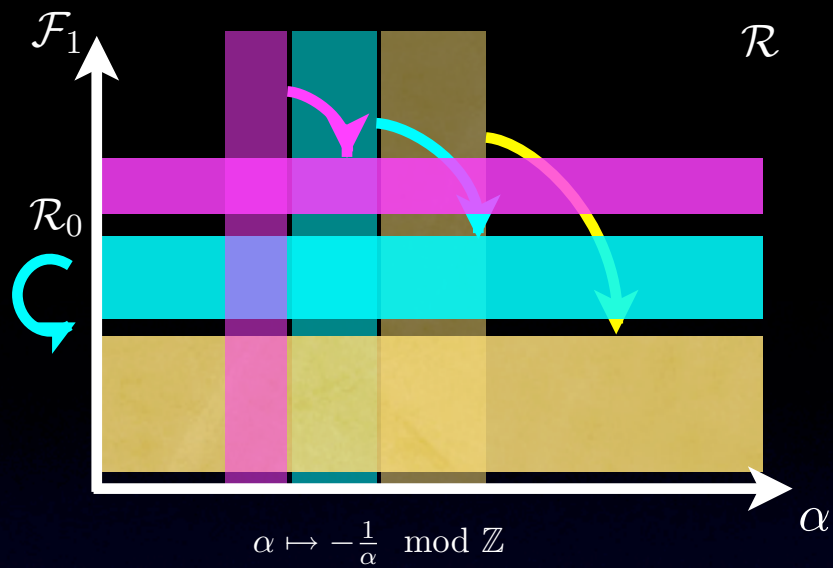
is well defined and expanding along α direction and uniformly contracting along \mathcal{F}_1 direction. Moreover $\mathcal{R}(e^{2\pi i\alpha} z + z^2)$ belong to the above set for $\alpha \in Irrat_N$.





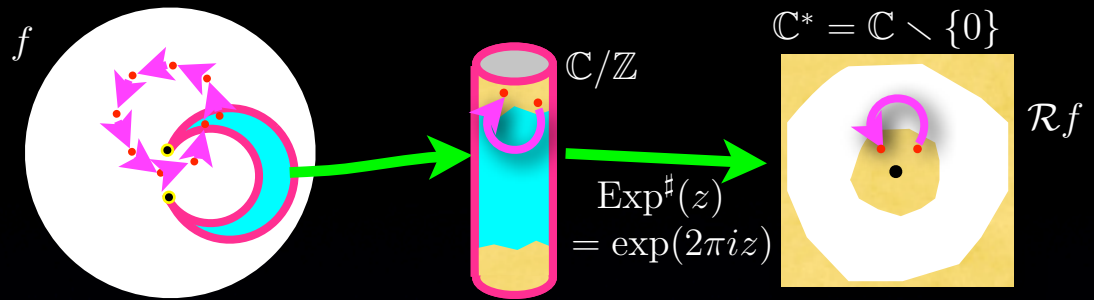
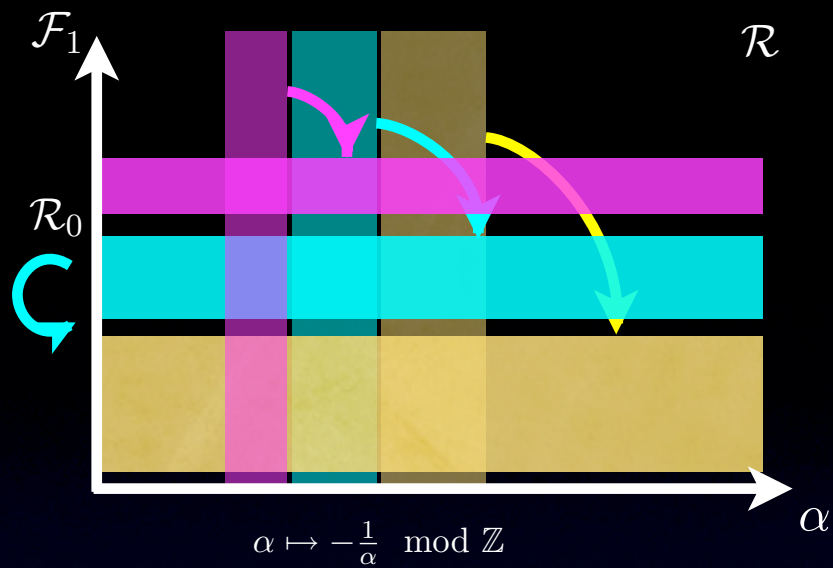


$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$



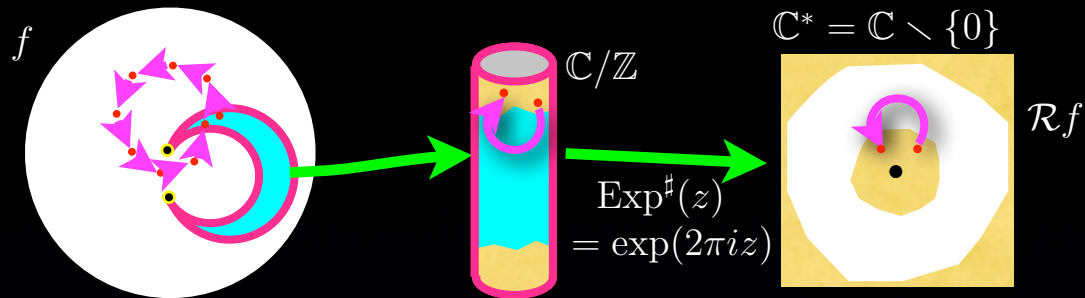
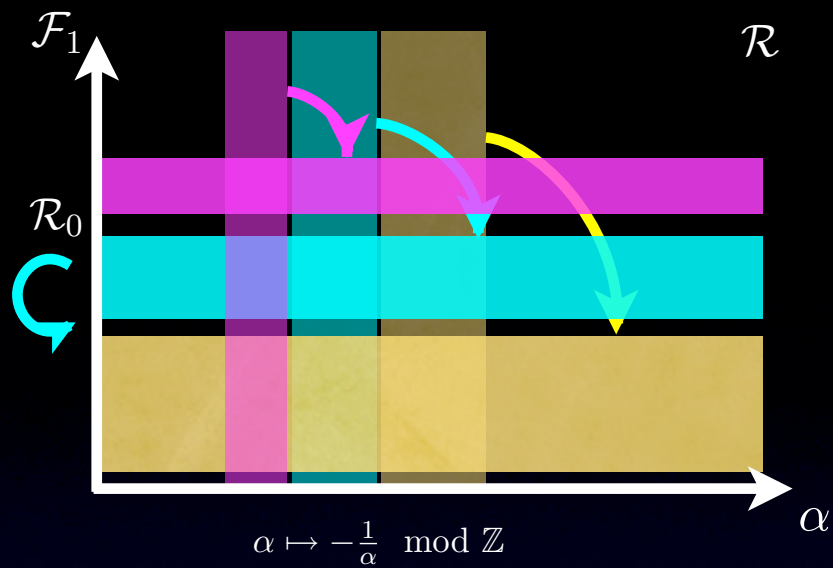
$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$

A priori bound:



$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$

A priori bound: $\mathcal{R} : e^{2\pi i \alpha} P \circ \varphi^{-1} \mapsto e^{-2\pi i / \alpha} P \circ \psi^{-1}$

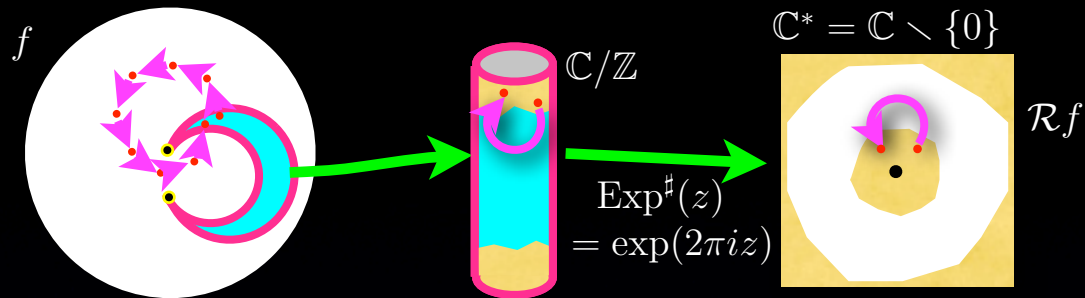
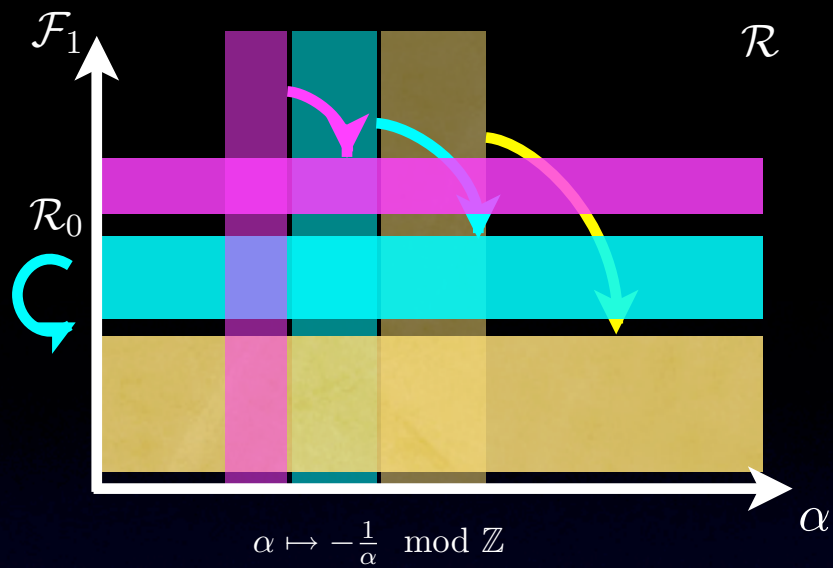


$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$

A priori bound: $\mathcal{R} : e^{2\pi i \alpha} P \circ \varphi^{-1} \mapsto e^{-2\pi i / \alpha} P \circ \psi^{-1}$

ψ satisfies the same condition as φ .

In fact, it extends to a fixed domain containing \overline{V} .



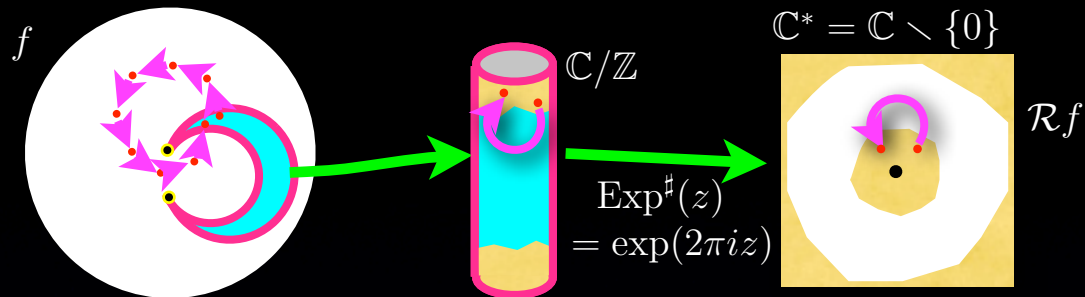
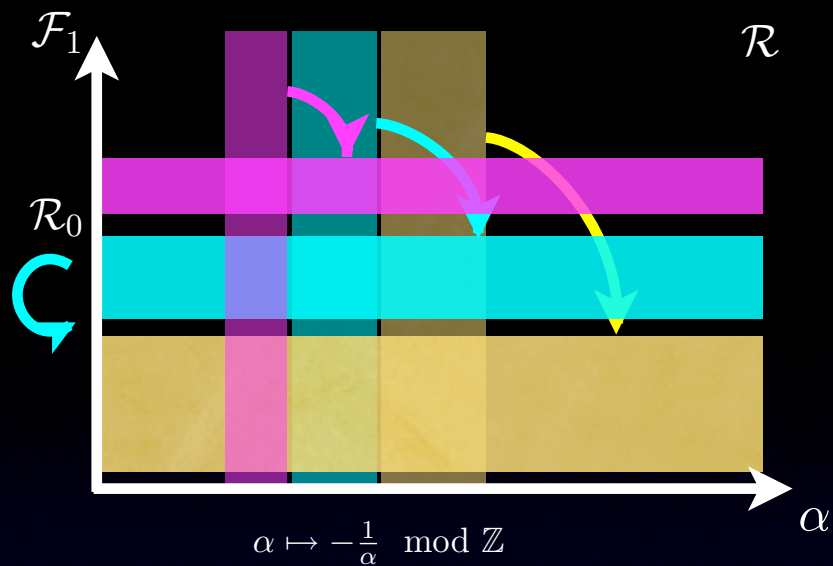
$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$

A priori bound: $\mathcal{R} : e^{2\pi i \alpha} P \circ \varphi^{-1} \mapsto e^{-2\pi i / \alpha} P \circ \psi^{-1}$

ψ satisfies the same condition as φ .

In fact, it extends to a fixed domain containing \overline{V} .

Hyperbolicity:



$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$

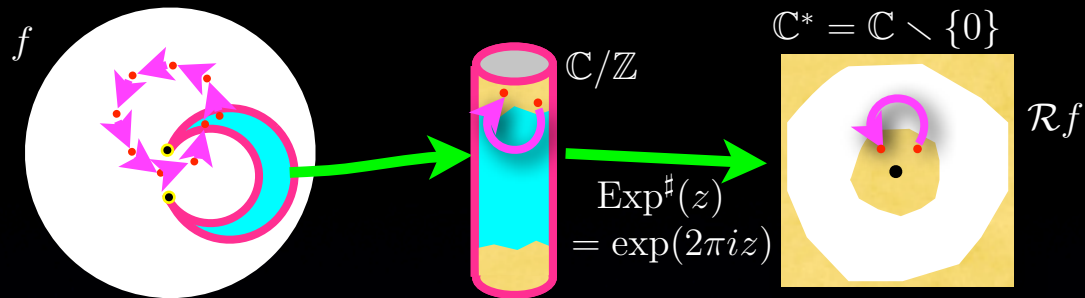
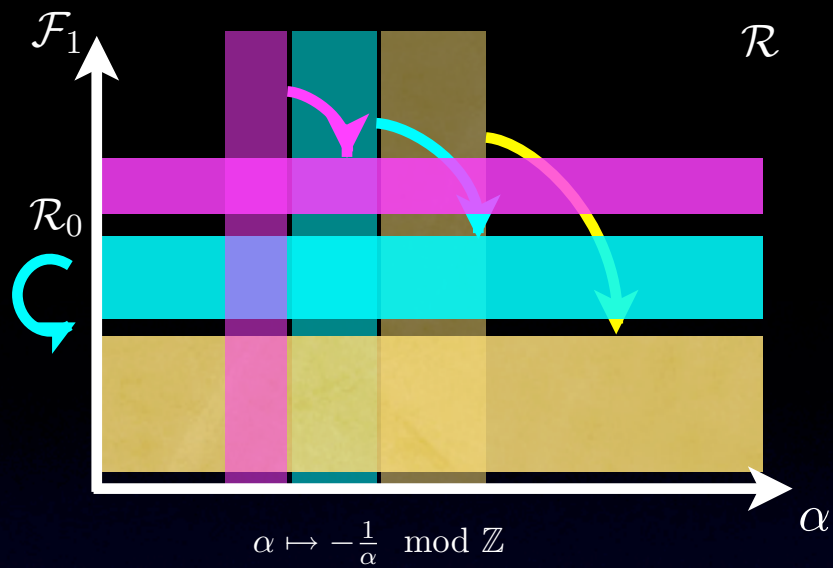
A priori bound: $\mathcal{R} : e^{2\pi i \alpha} P \circ \varphi^{-1} \mapsto e^{-2\pi i / \alpha} P \circ \psi^{-1}$

ψ satisfies the same condition as φ .

In fact, it extends to a fixed domain containing \overline{V} .

Hyperbolicity: $P \circ \varphi^{-1} \longleftrightarrow \varphi \longleftrightarrow [\tilde{\varphi}] \in \text{Teich}(\mathbb{C} \setminus V)$

where $\tilde{\varphi}$ is a qc-extension of φ to $\mathbb{C} \setminus V$



$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$

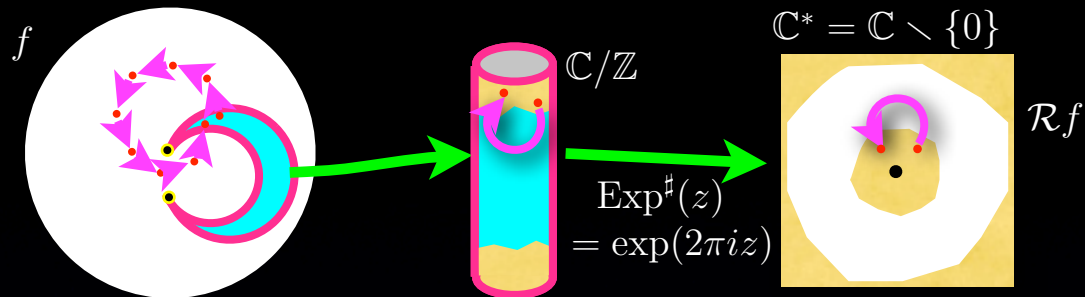
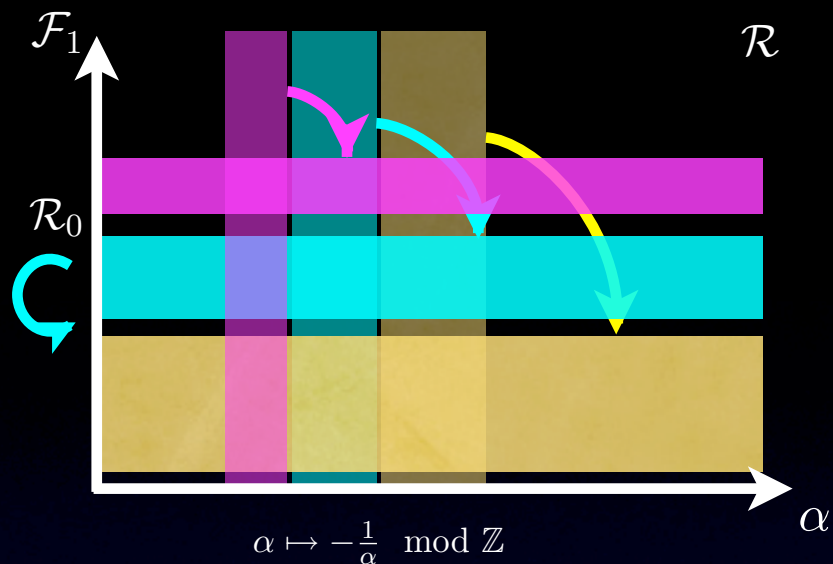
A priori bound: $\mathcal{R} : e^{2\pi i \alpha} P \circ \varphi^{-1} \mapsto e^{-2\pi i / \alpha} P \circ \psi^{-1}$

ψ satisfies the same condition as φ .

In fact, it extends to a fixed domain containing \overline{V} .

Hyperbolicity: $P \circ \varphi^{-1} \longleftrightarrow \varphi \longleftrightarrow [\tilde{\varphi}] \in \text{Teich}(\mathbb{C} \setminus V)$

where $\tilde{\varphi}$ is a qc-extension of φ to $\mathbb{C} \setminus V$
Royden-Gardiner: Holomorphic map between Teichmüller spaces are weakly contracting.



$$\mathcal{F}_1 = \left\{ f = P \circ \varphi^{-1} \mid \begin{array}{l} \varphi : V \rightarrow \mathbb{C} \text{ is univalent (with qc extension)} \\ \varphi(0) = 0, \varphi'(0) = 1 \end{array} \right\}$$

A priori bound: $\mathcal{R} : e^{2\pi i \alpha} P \circ \varphi^{-1} \mapsto e^{-2\pi i / \alpha} P \circ \psi^{-1}$

ψ satisfies the same condition as φ .

In fact, it extends to a fixed domain containing \overline{V} .

Hyperbolicity: $P \circ \varphi^{-1} \longleftrightarrow \varphi \longleftrightarrow [\tilde{\varphi}] \in \text{Teich}(\mathbb{C} \setminus V)$

where $\tilde{\varphi}$ is a qc-extension of φ to $\mathbb{C} \setminus V$
Royden-Gardiner: Holomorphic map between Teichmüller spaces are weakly contracting.

$\tilde{\psi}$ conformal near $\partial V \implies$ strict contraction.

More on a priori bound

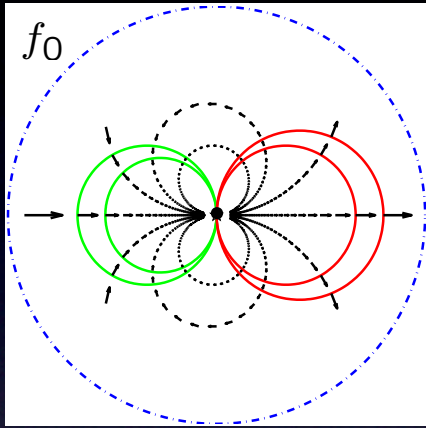
More on a priori bound

$\alpha = 0$
parabolic fixed point

⋮

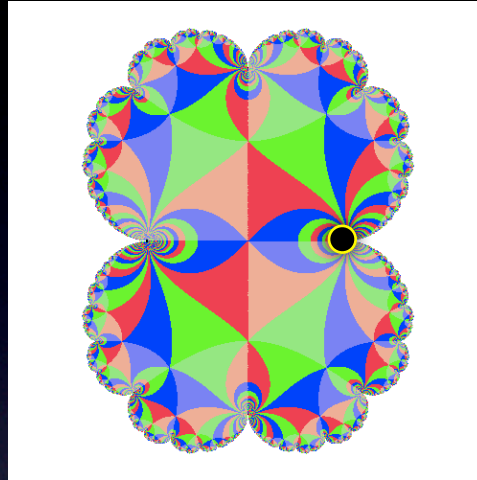
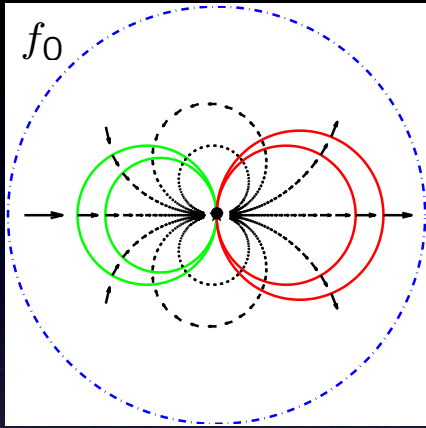
More on a priori bound

$\alpha = 0$
parabolic fixed point



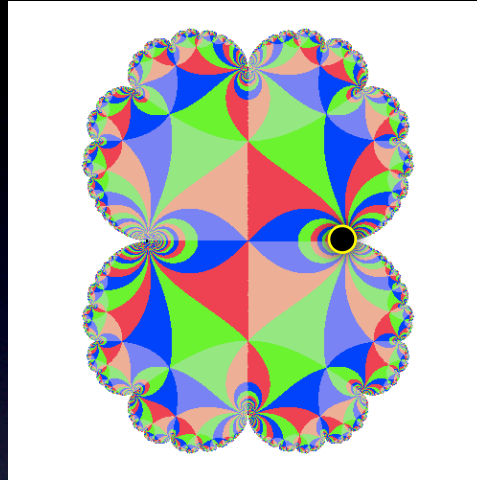
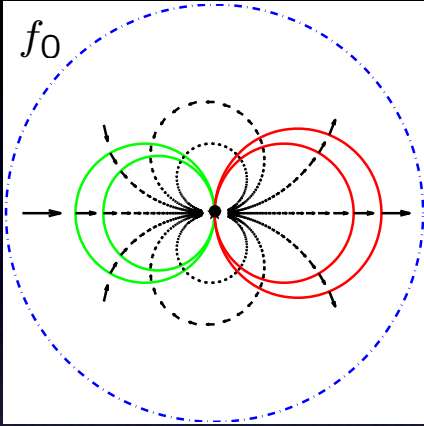
More on a priori bound

$\alpha = 0$
parabolic fixed point

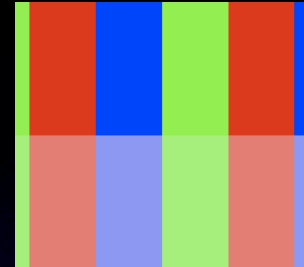


More on a priori bound

$\alpha = 0$
parabolic fixed point

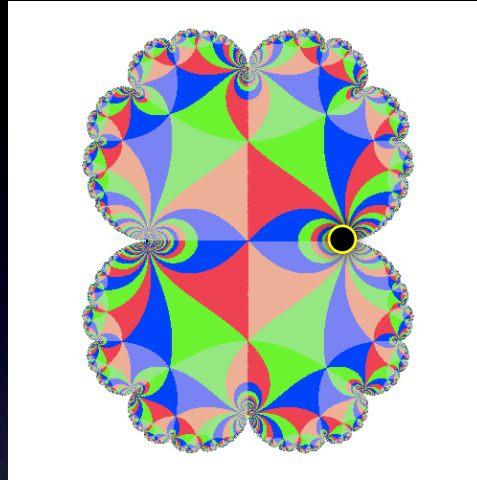
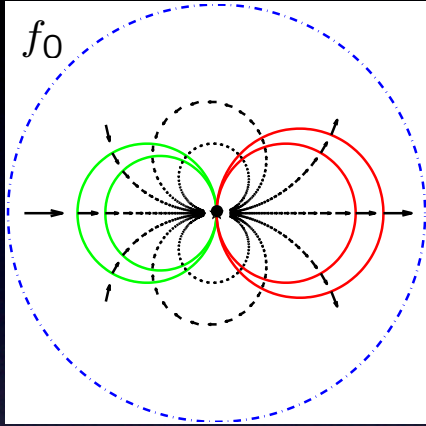


Φ_{attr}

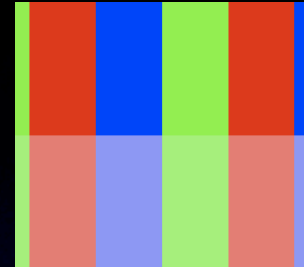


More on a priori bound

$\alpha = 0$
parabolic fixed point



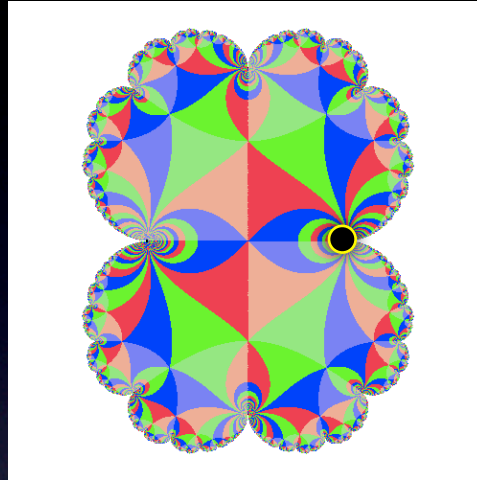
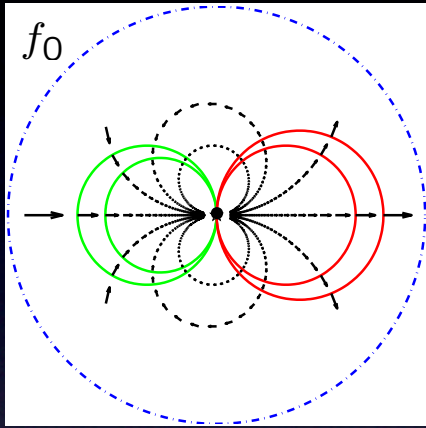
Φ_{attr}



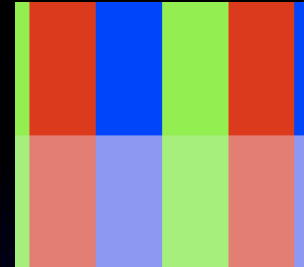
F_{can} ($\infty =$ fixed pt)

More on a priori bound

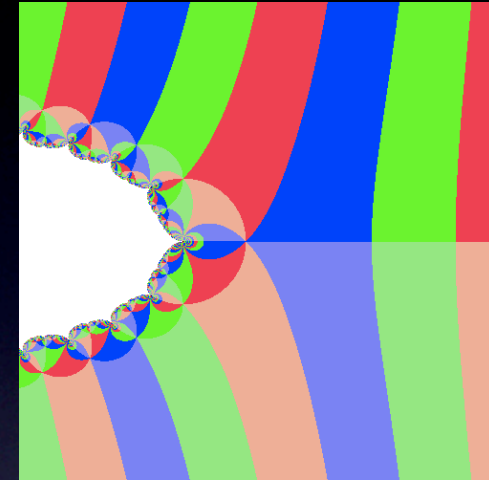
$\alpha = 0$
parabolic fixed point



Φ_{attr}

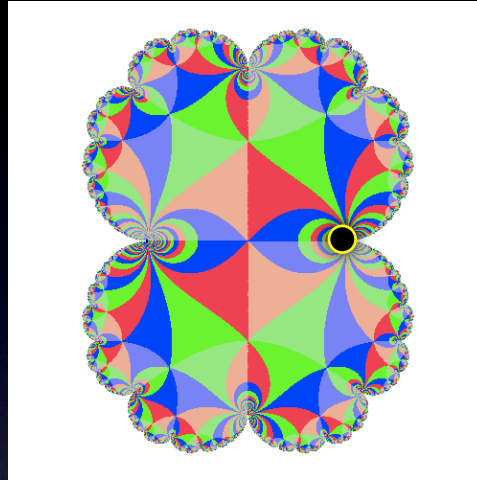
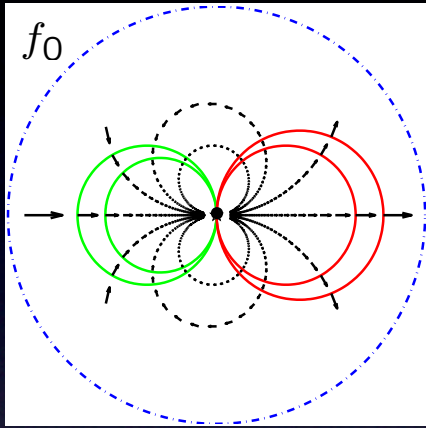


F_{can} ($\infty =$ fixed pt)

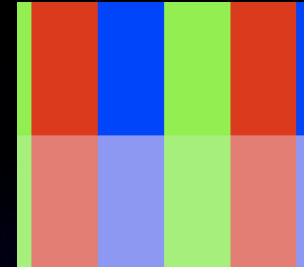


More on a priori bound

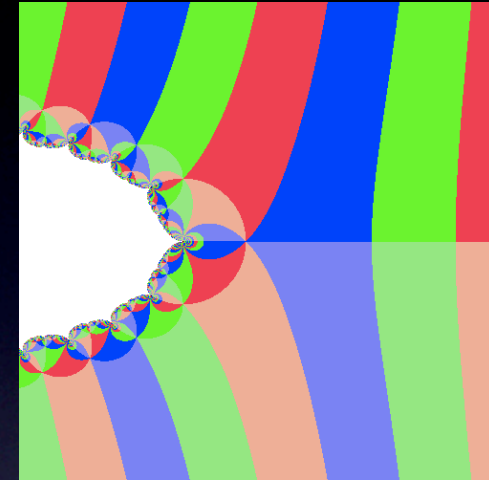
$\alpha = 0$
parabolic fixed point



Φ_{attr}



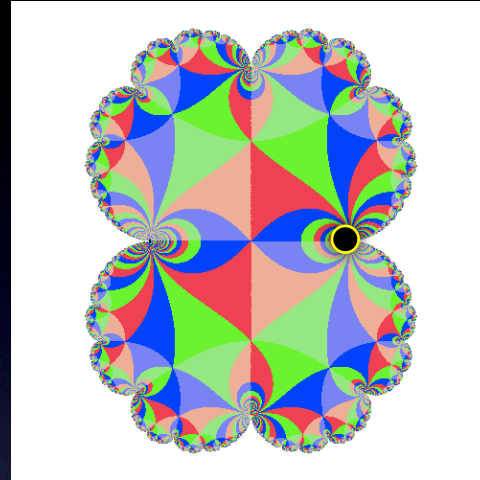
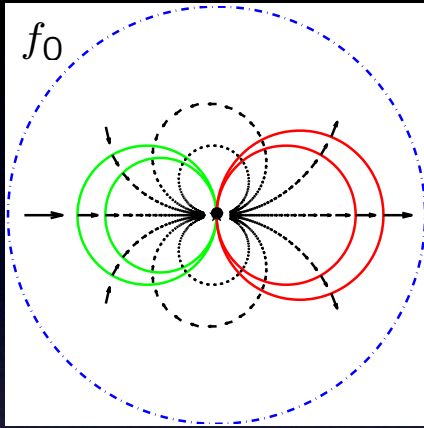
F_{can} ($\infty =$ fixed pt)



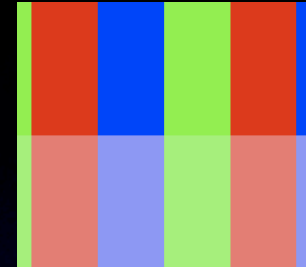
irrat. indiff. (near-parabolic)

More on a priori bound

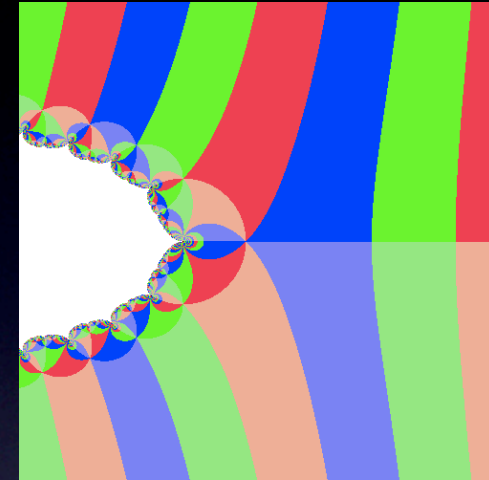
$\alpha = 0$
parabolic fixed point



Φ_{attr}



F_{can} ($\infty =$ fixed pt)

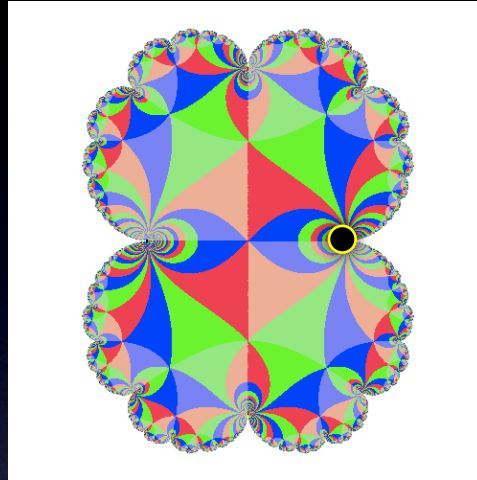
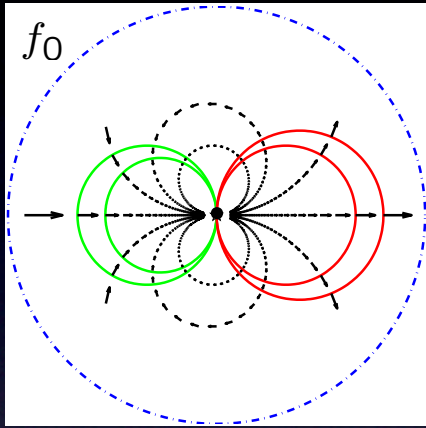


irrat. indiff. (near-parabolic)

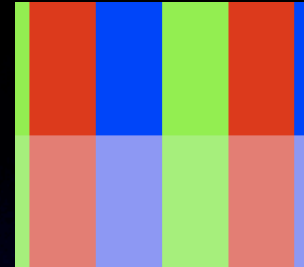
use Douady-Hubbard-Lavaurs theory of parabolic implosion

More on a priori bound

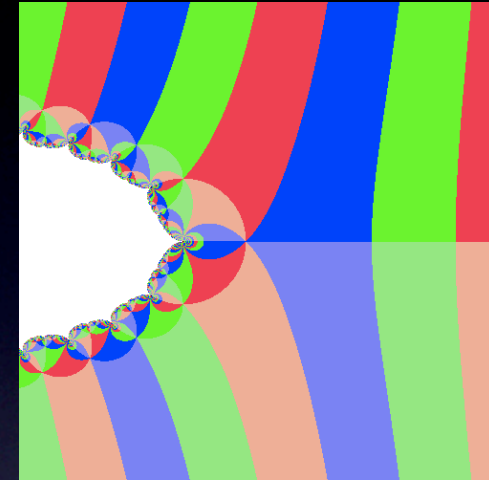
$\alpha = 0$
parabolic fixed point



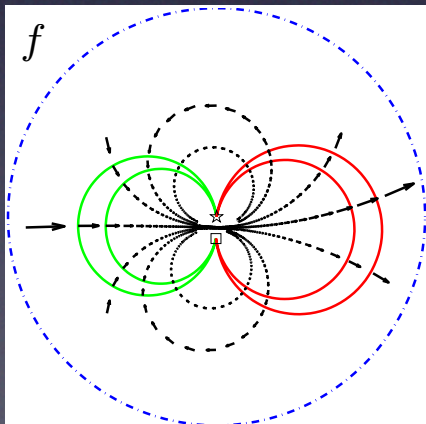
Φ_{attr}



F_{can} ($\infty = \text{fixed pt}$)



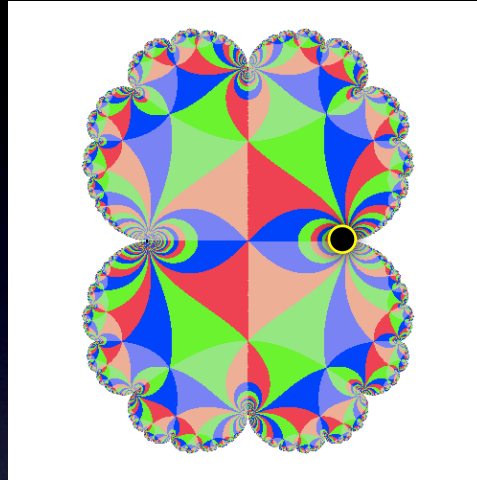
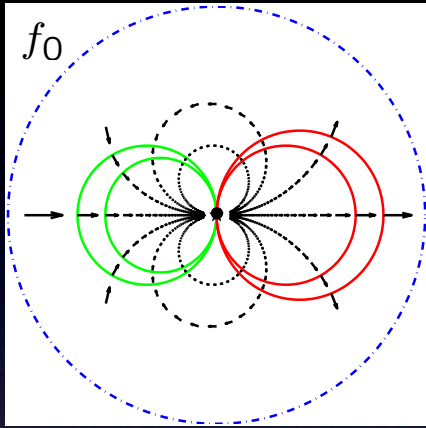
irrat. indiff. (near-parabolic)



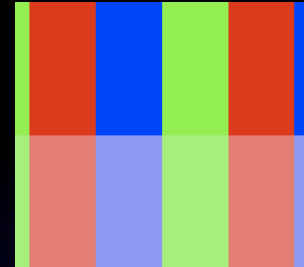
use Douady-Hubbard-Lavaurs theory of parabolic implosion

More on a priori bound

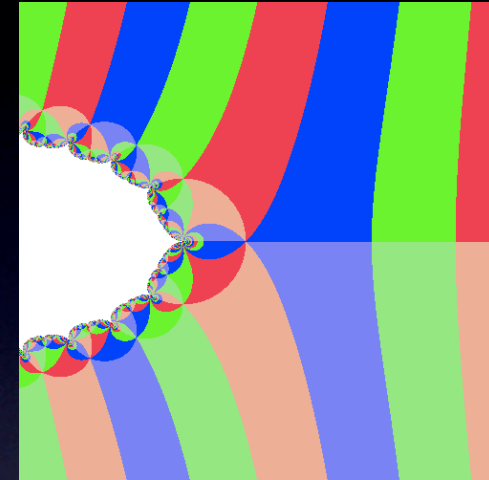
$\alpha = 0$
parabolic fixed point



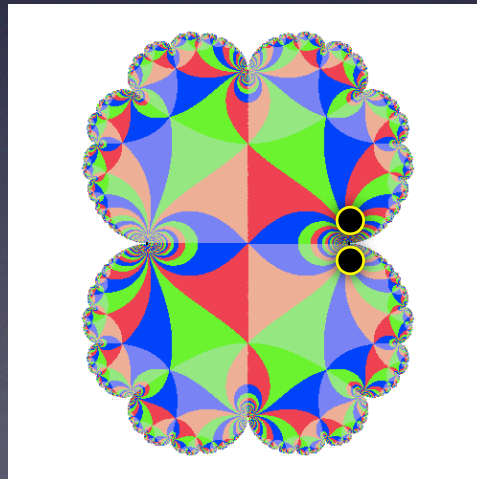
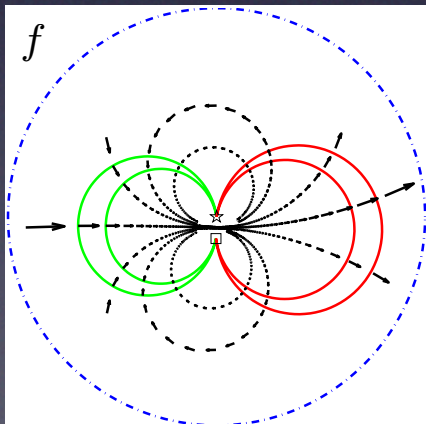
Φ_{attr}



F_{can} ($\infty =$ fixed pt)



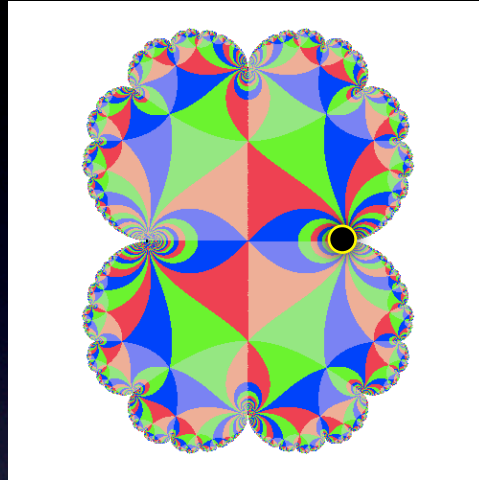
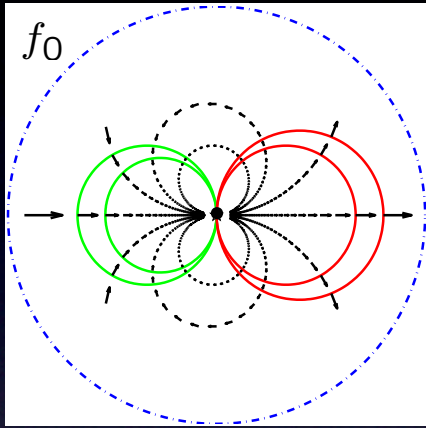
irrat. indiff. (near-parabolic)



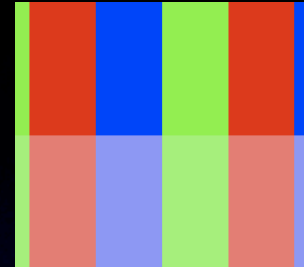
use Douady-Hubbard-Lavaurs theory of parabolic implosion

More on a priori bound

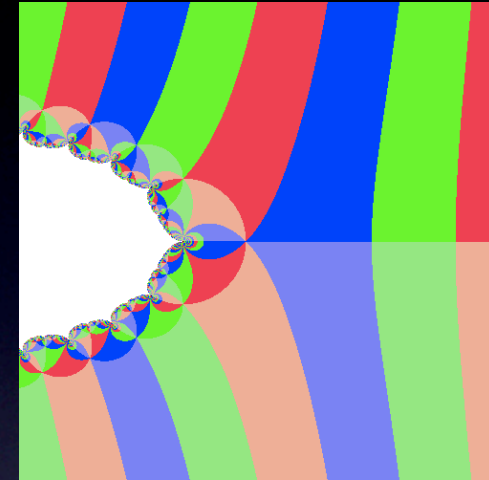
$\alpha = 0$
parabolic fixed point



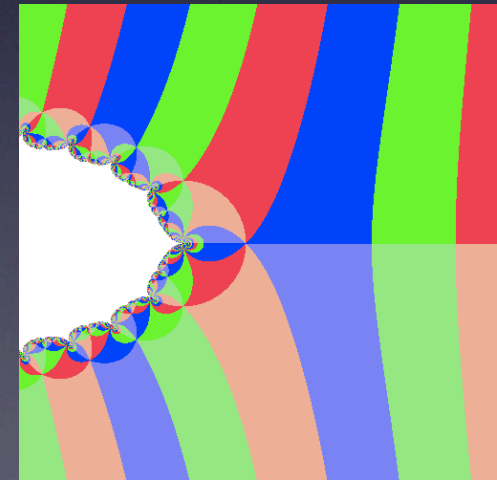
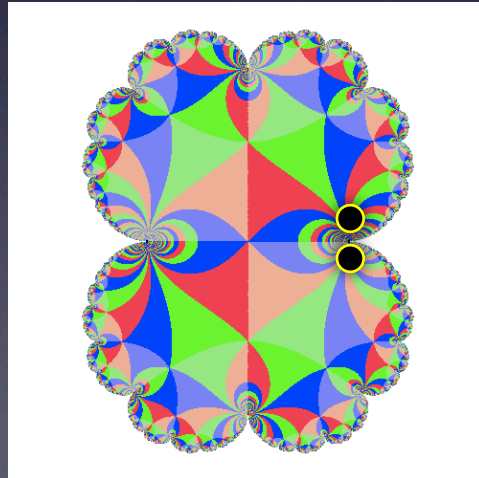
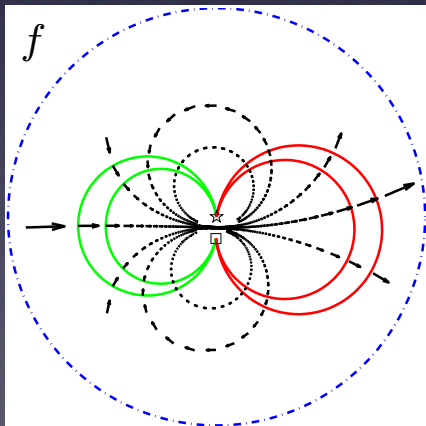
Φ_{attr}



F_{can} ($\infty = \text{fixed pt}$)



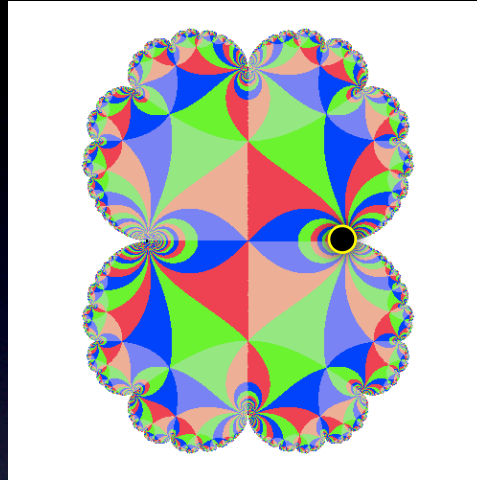
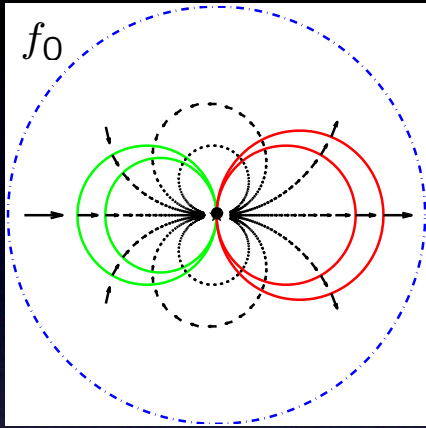
irrat. indiff. (near-parabolic)



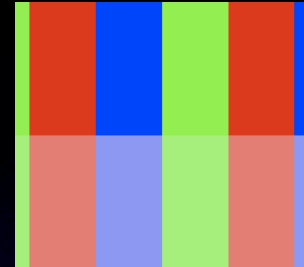
use Douady-Hubbard-Lavaurs theory of parabolic implosion

More on a priori bound

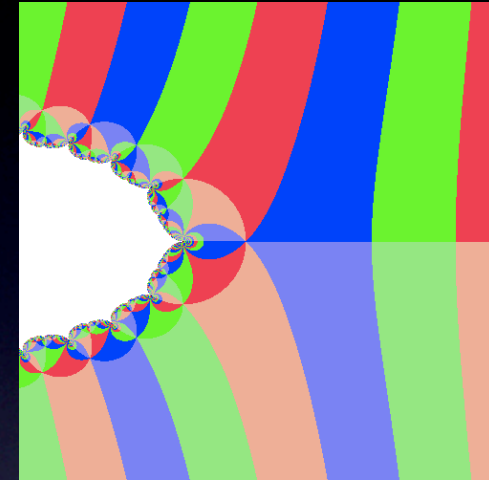
$\alpha = 0$
parabolic fixed point



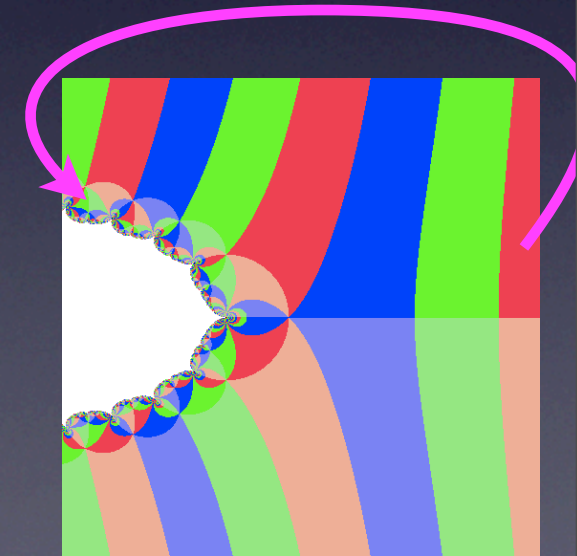
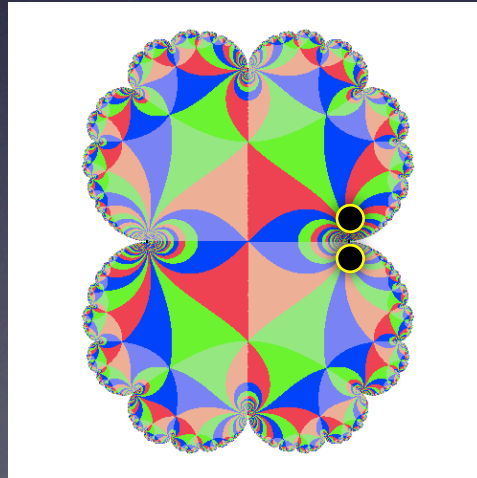
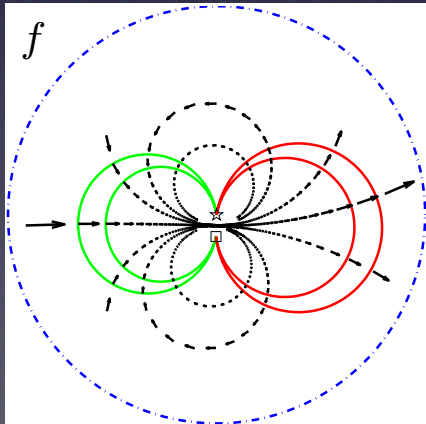
Φ_{attr}



F_{can} ($\infty = \text{fixed pt}$)



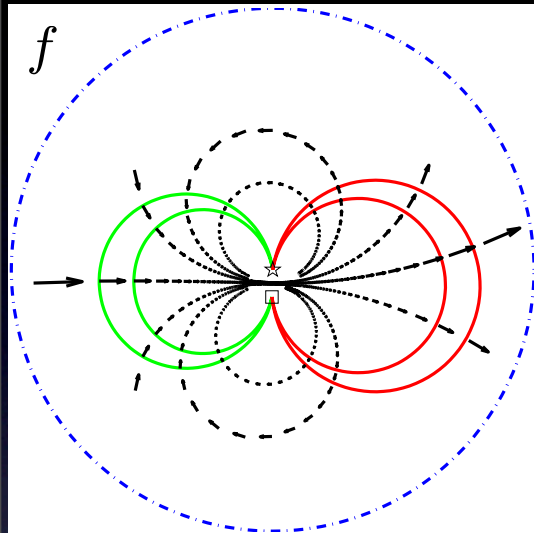
irrat. indiff. (near-parabolic)



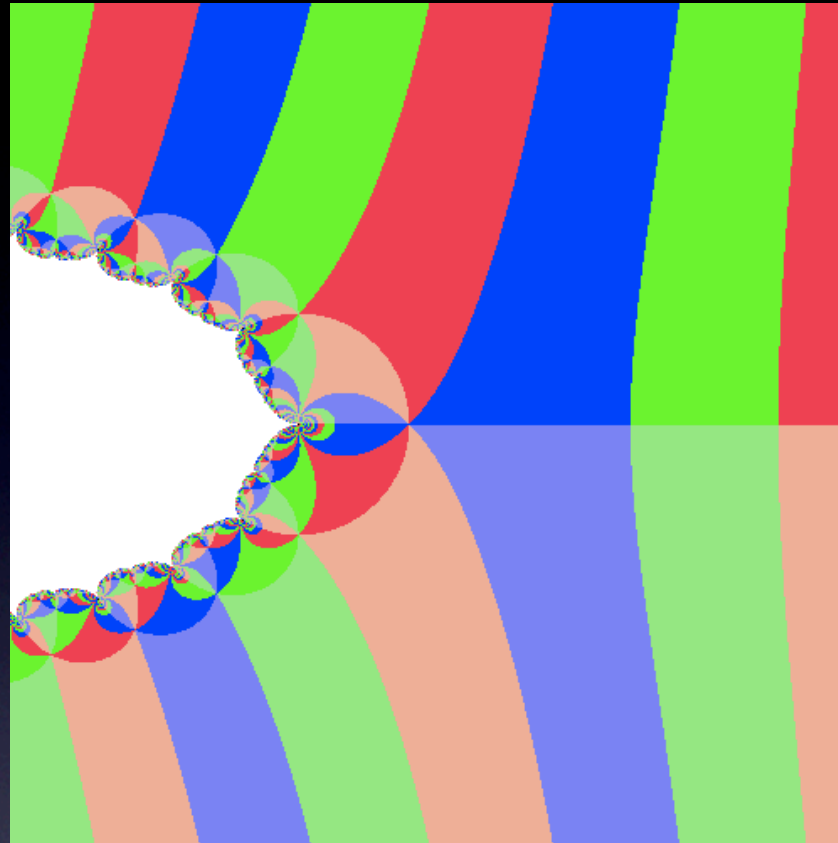
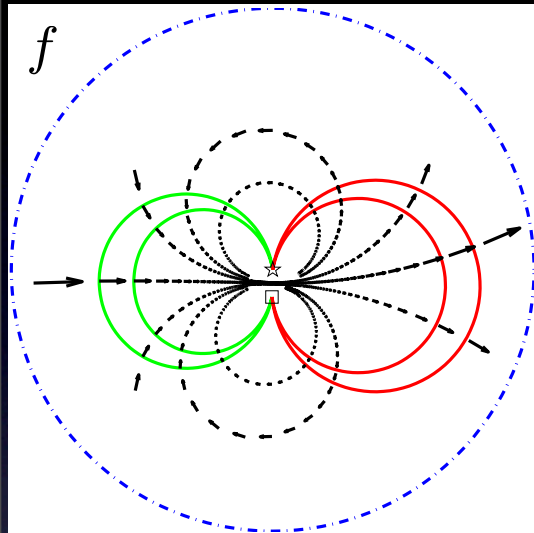
use Douady-Hubbard-Lavaurs theory of parabolic implosion

For small $\alpha \neq 0$, not every detail of the pattern is preserved.

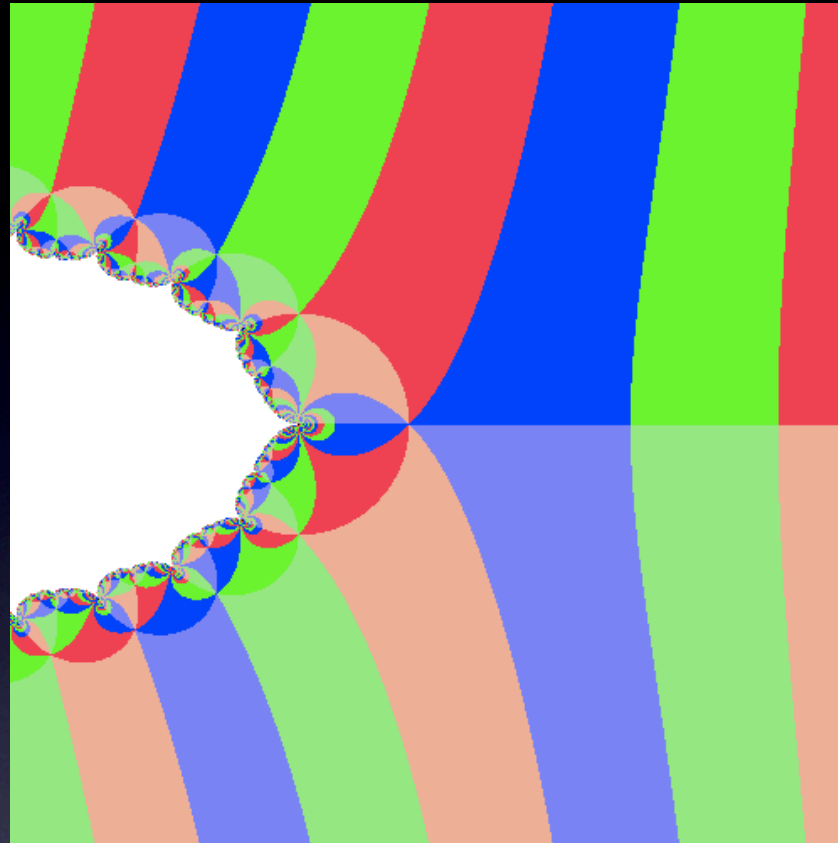
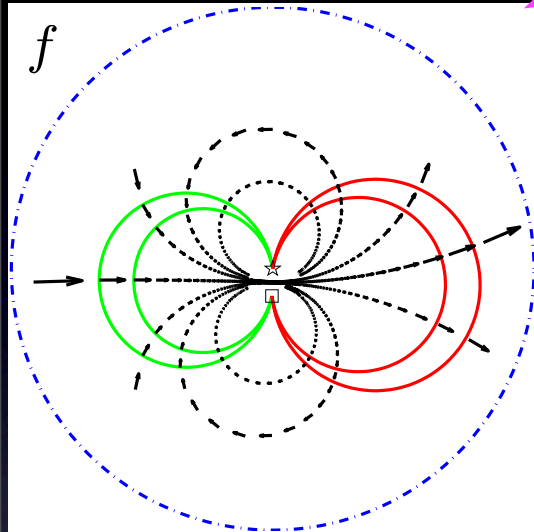
For small $\alpha \neq 0$, not every detail of the pattern is preserved.



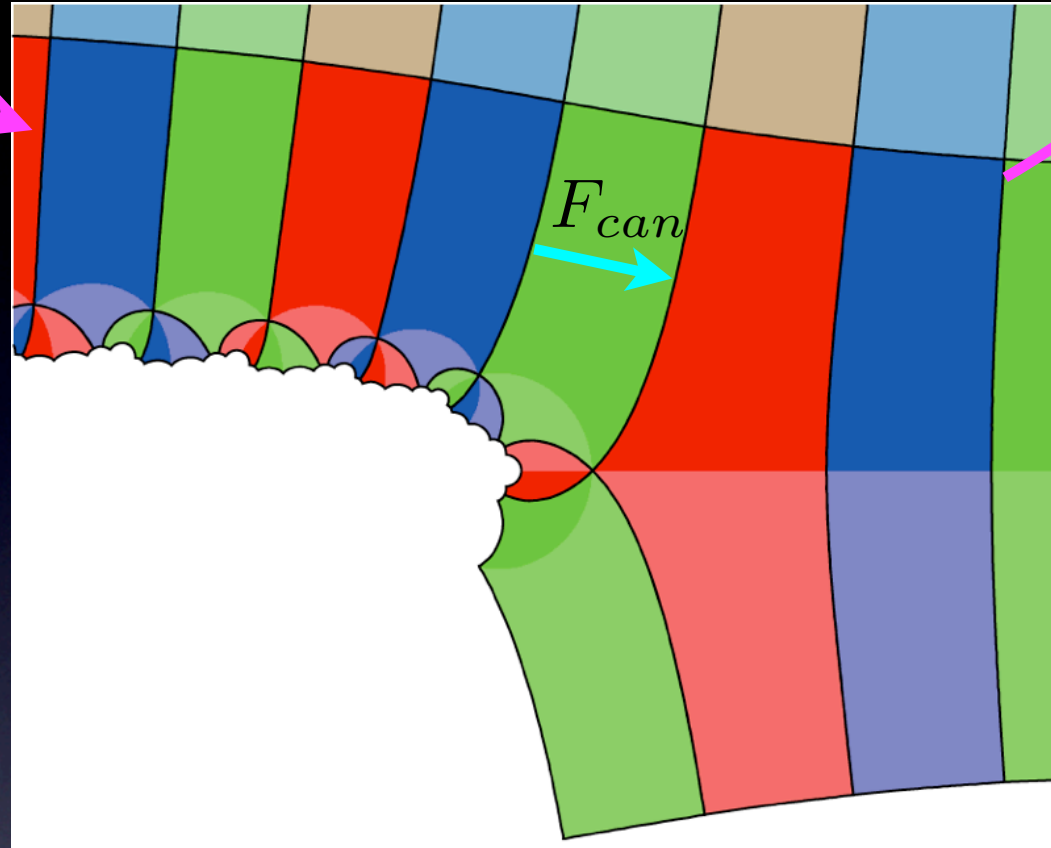
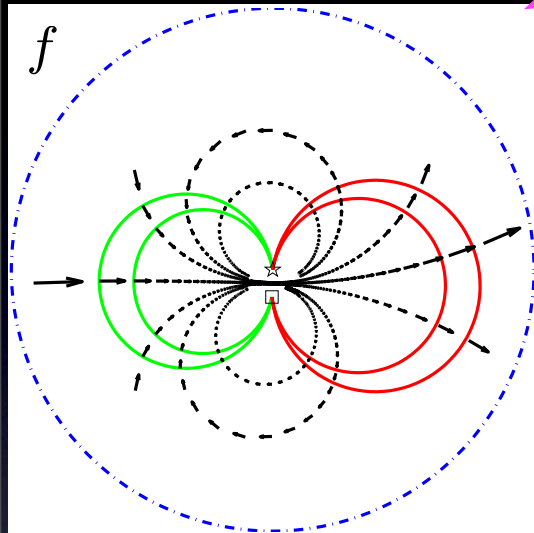
For small $\alpha \neq 0$, not every detail of the pattern is preserved.



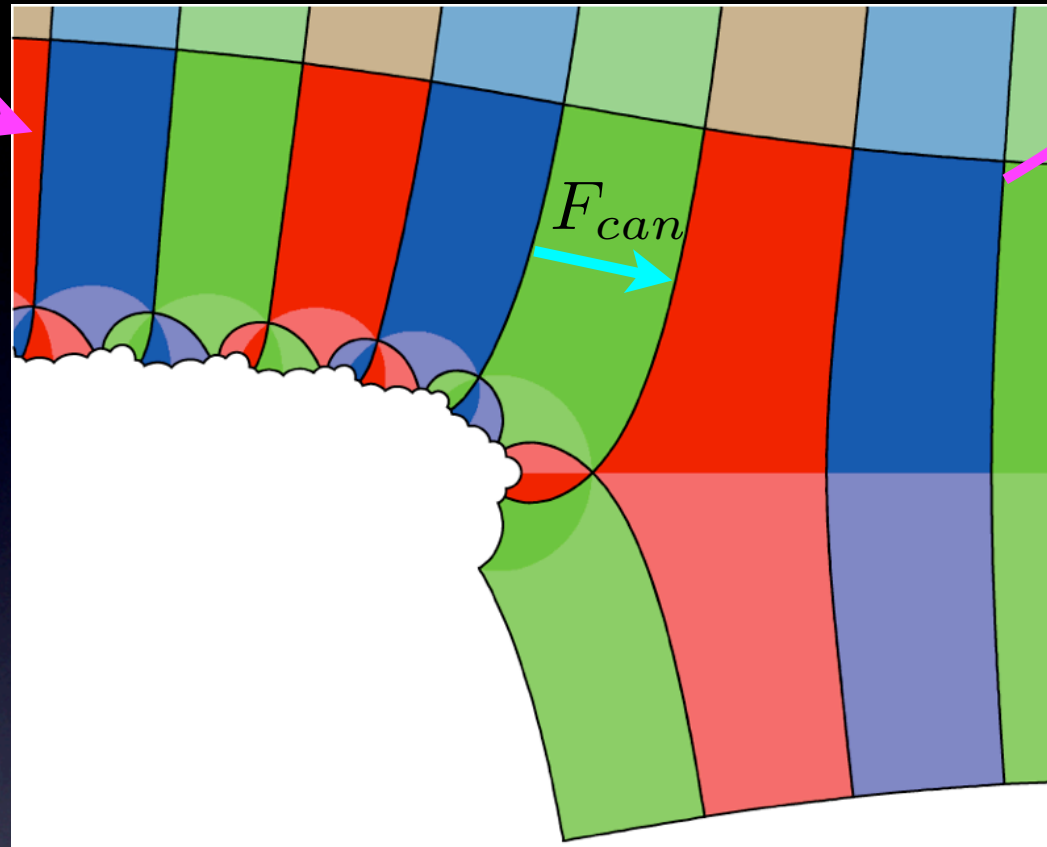
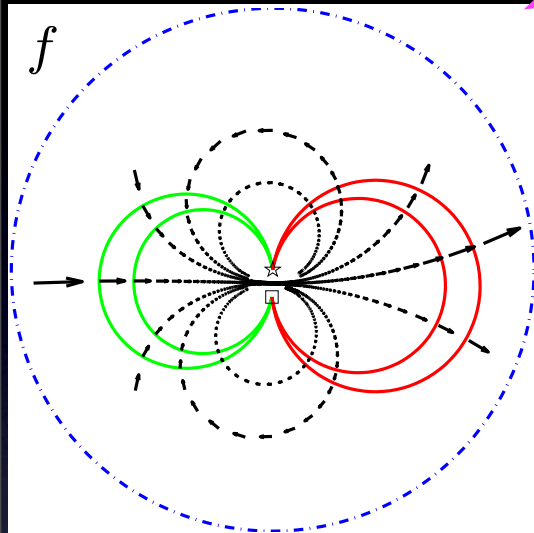
For small $\alpha \neq 0$, not every detail of the pattern is preserved.



For small $\alpha \neq 0$, not every detail of the pattern is preserved.

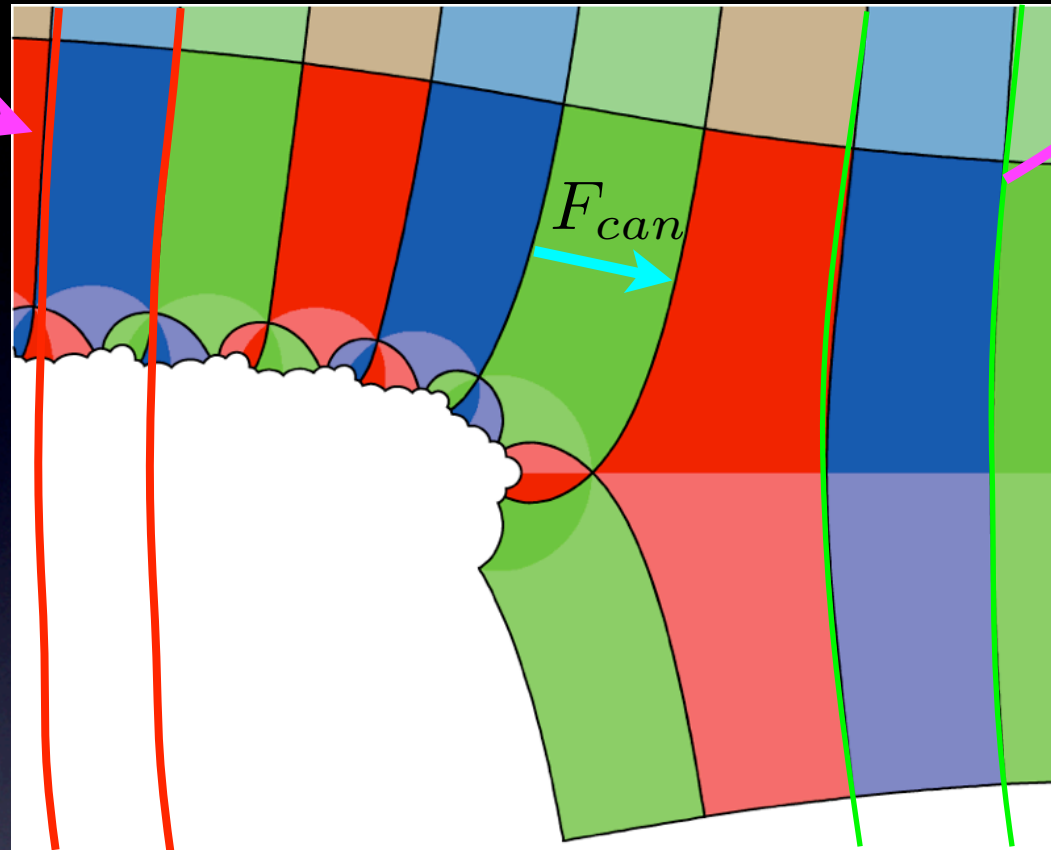
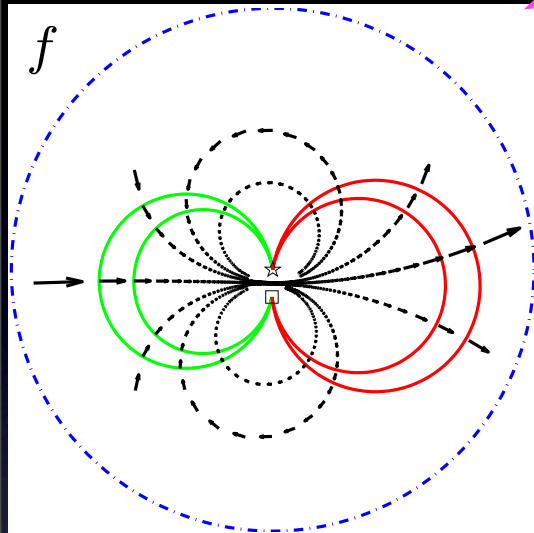


For small $\alpha \neq 0$, not every detail of the pattern is preserved.



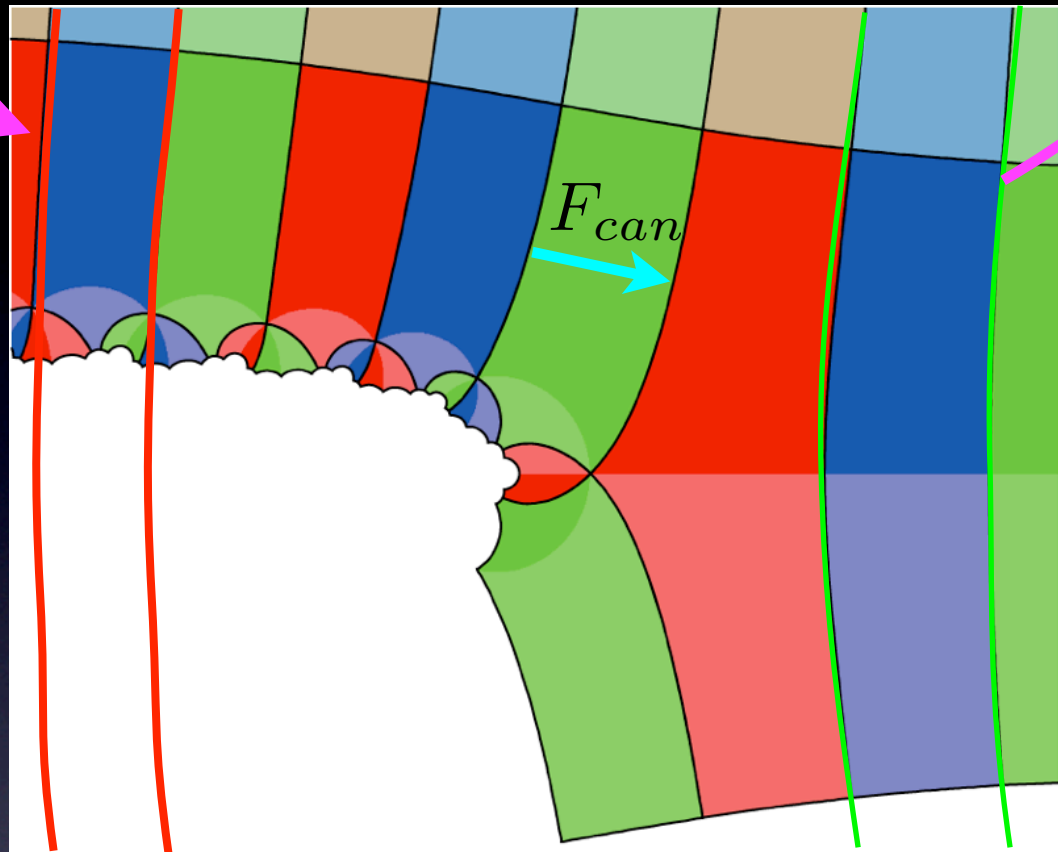
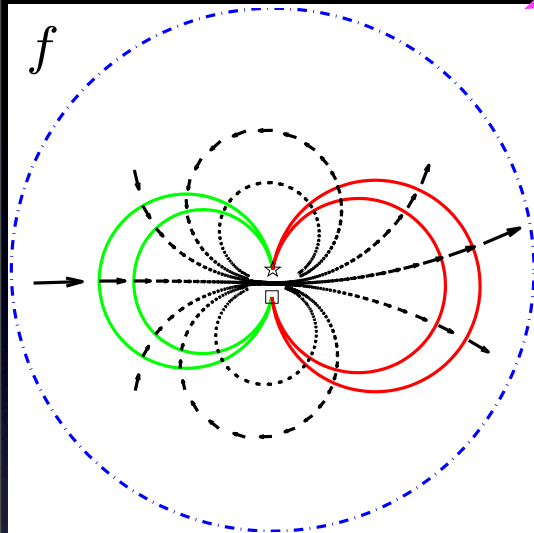
With many estimates, one can show that this much of the pattern is preserved.

For small $\alpha \neq 0$, not every detail of the pattern is preserved.

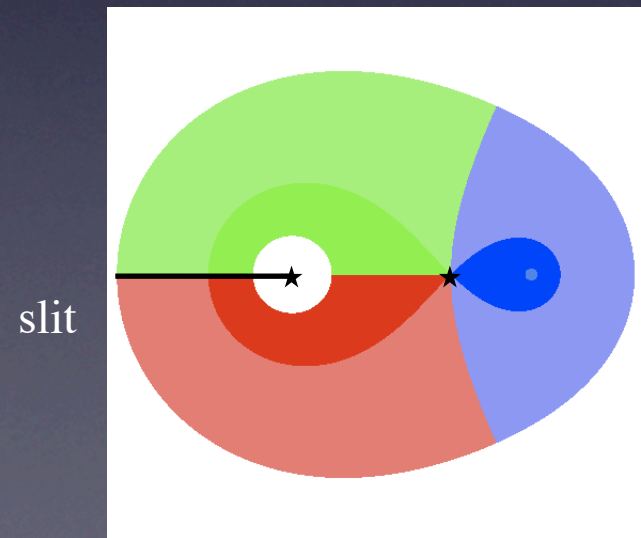


With many estimates, one can show that this much of the pattern is preserved.

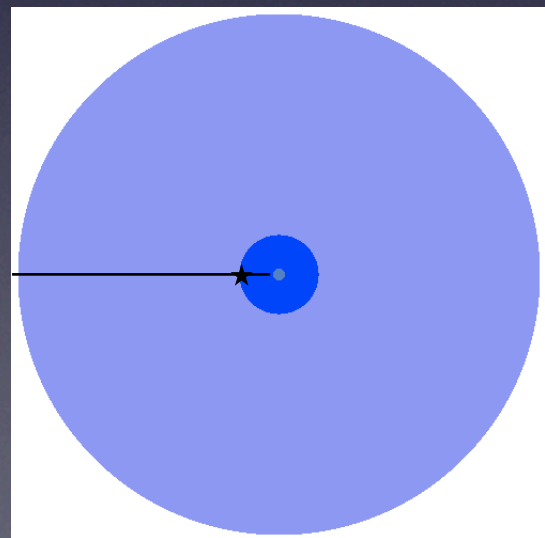
For small $\alpha \neq 0$, not every detail of the pattern is preserved.



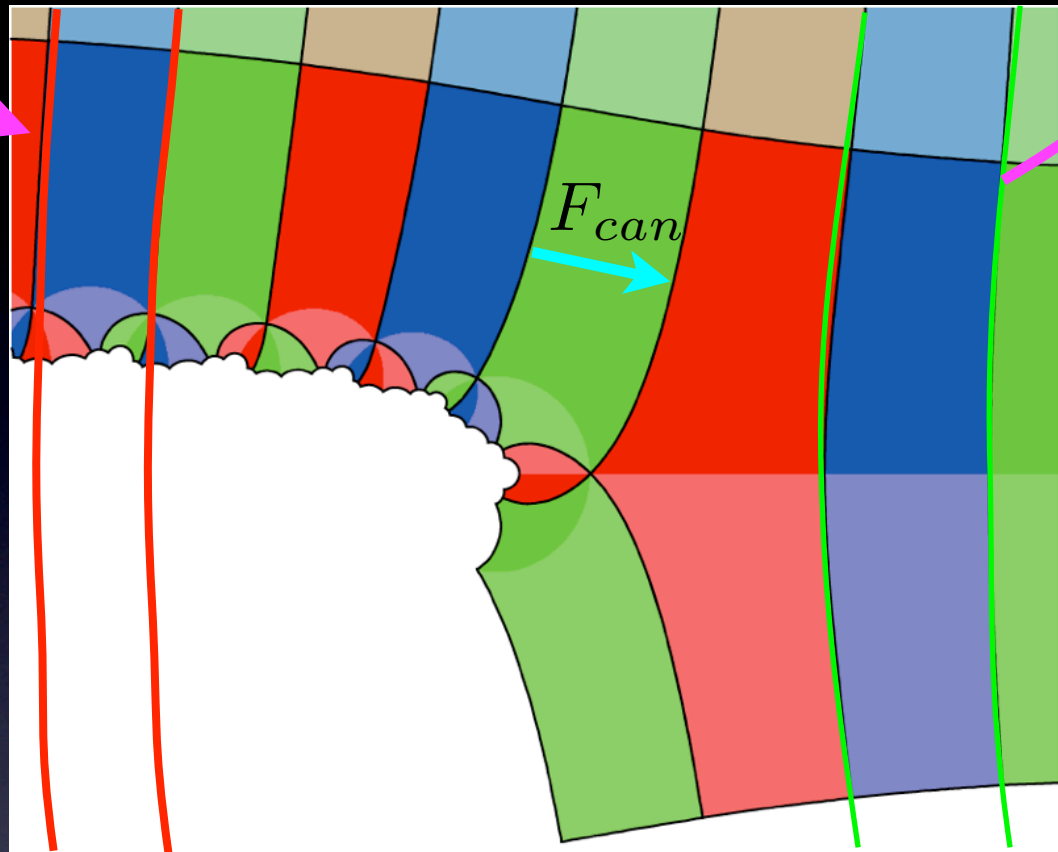
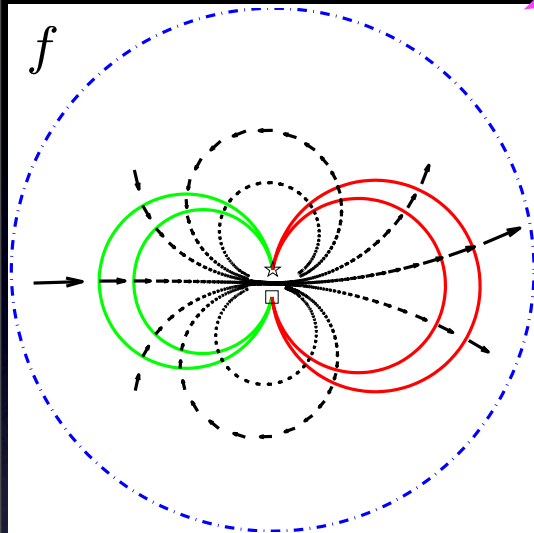
With many estimates, one can show that this much of the pattern is preserved.



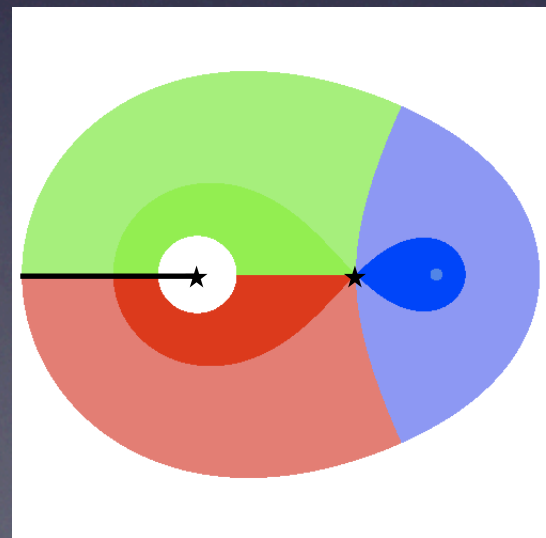
$\rightarrow P$



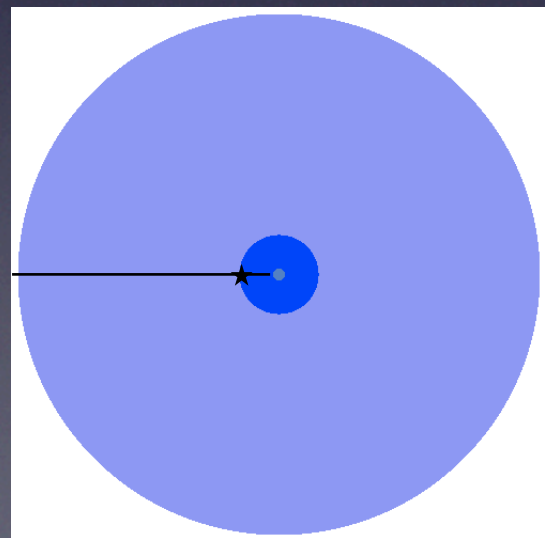
For small $\alpha \neq 0$, not every detail of the pattern is preserved.



With many estimates, one can show that this much of the pattern is preserved.



$\rightarrow P$



The renormalization has the same covering type

Applications of Inou-S. Theorem

Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

For simplicity, let us consider the sector renormalization:

Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

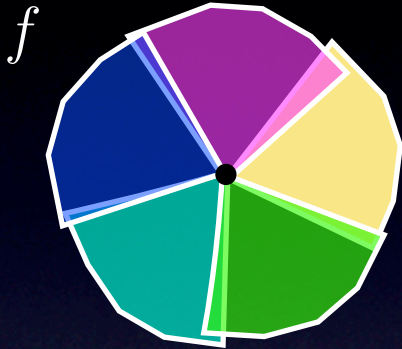
For simplicity, let us consider the sector renormalization:

f

Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

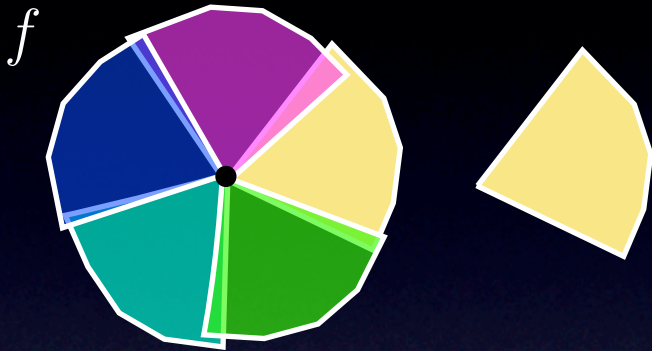
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

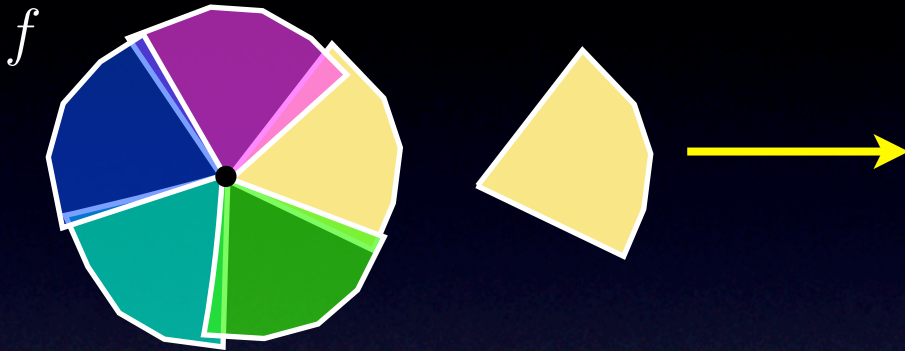
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

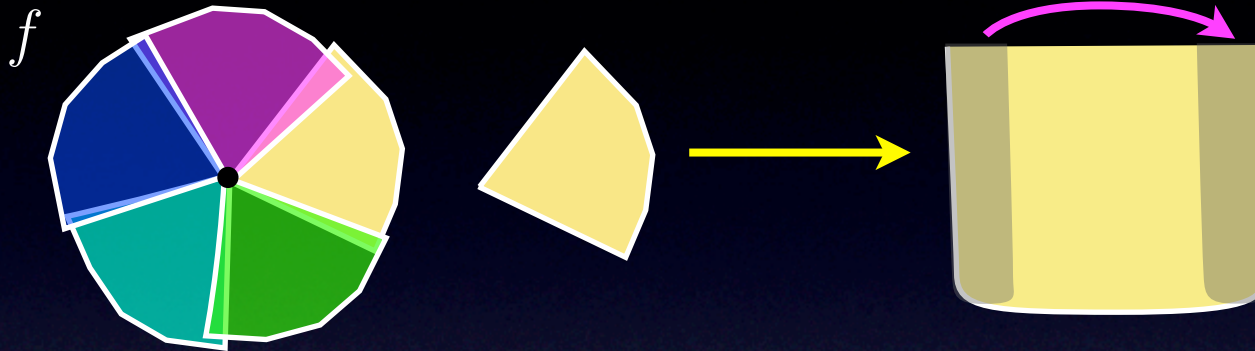
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

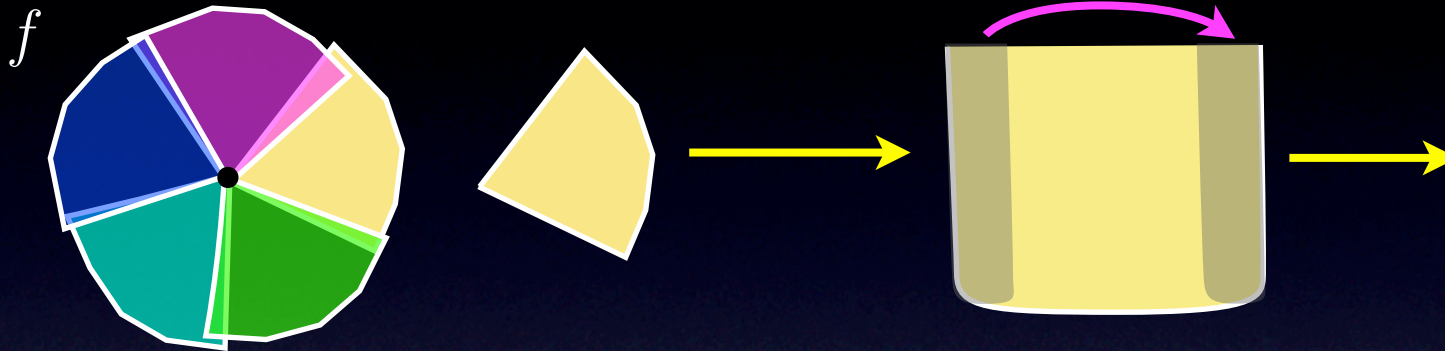
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

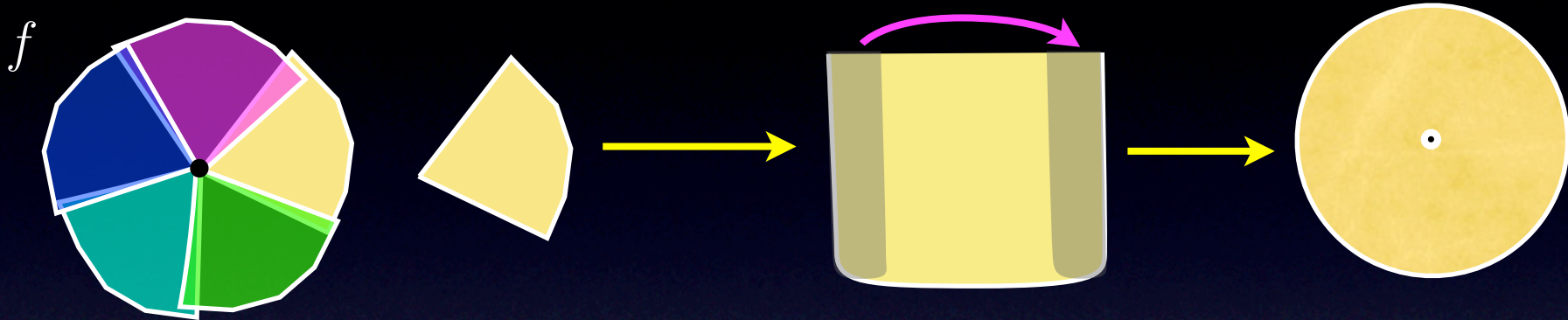
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

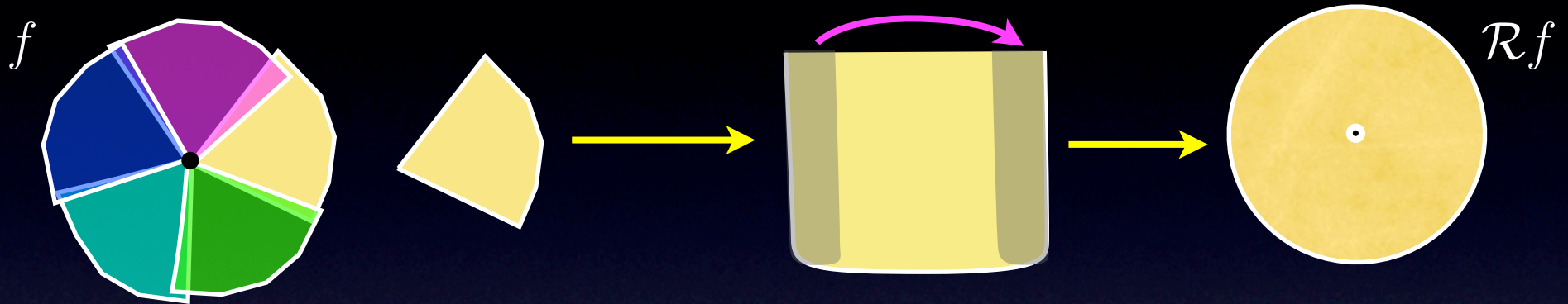
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

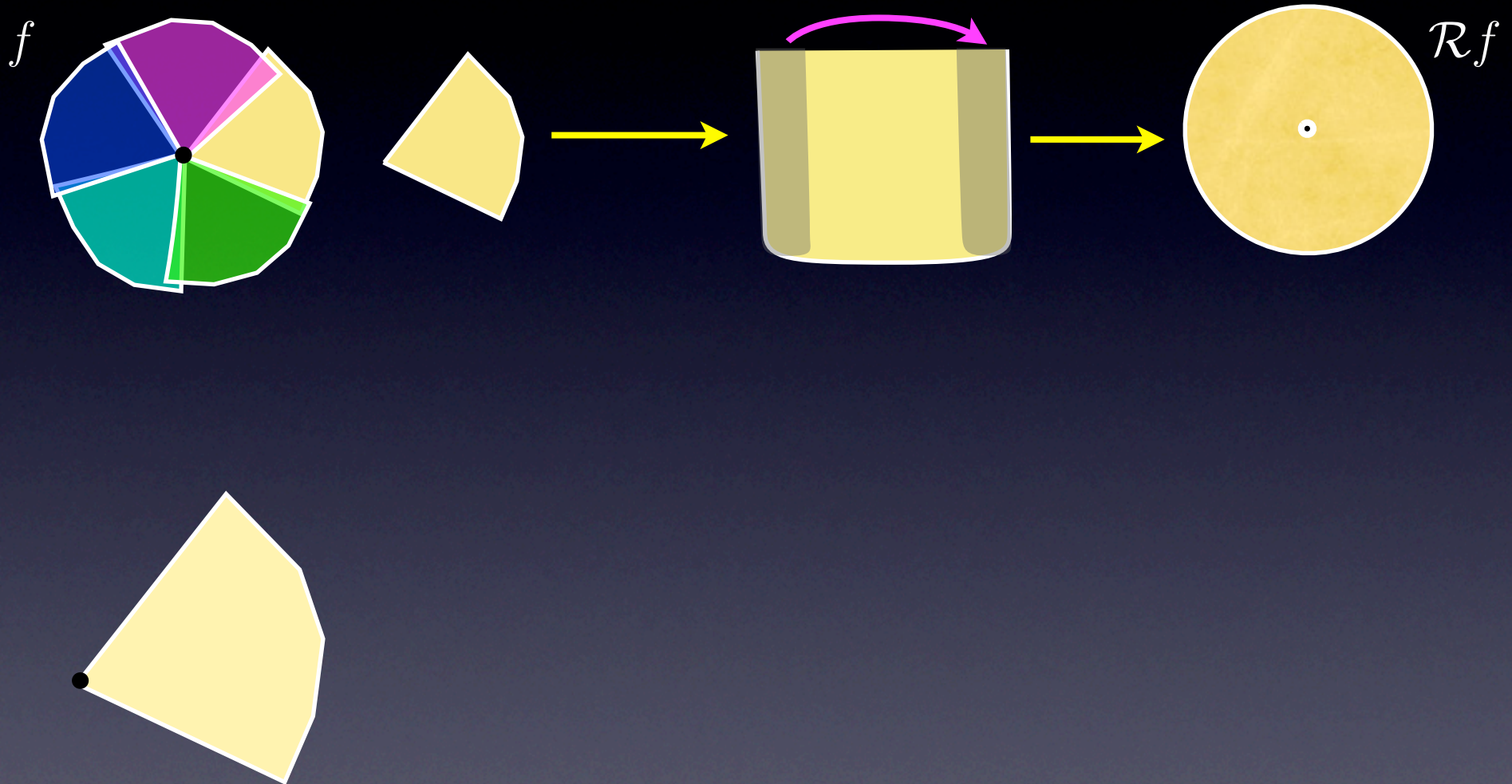
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

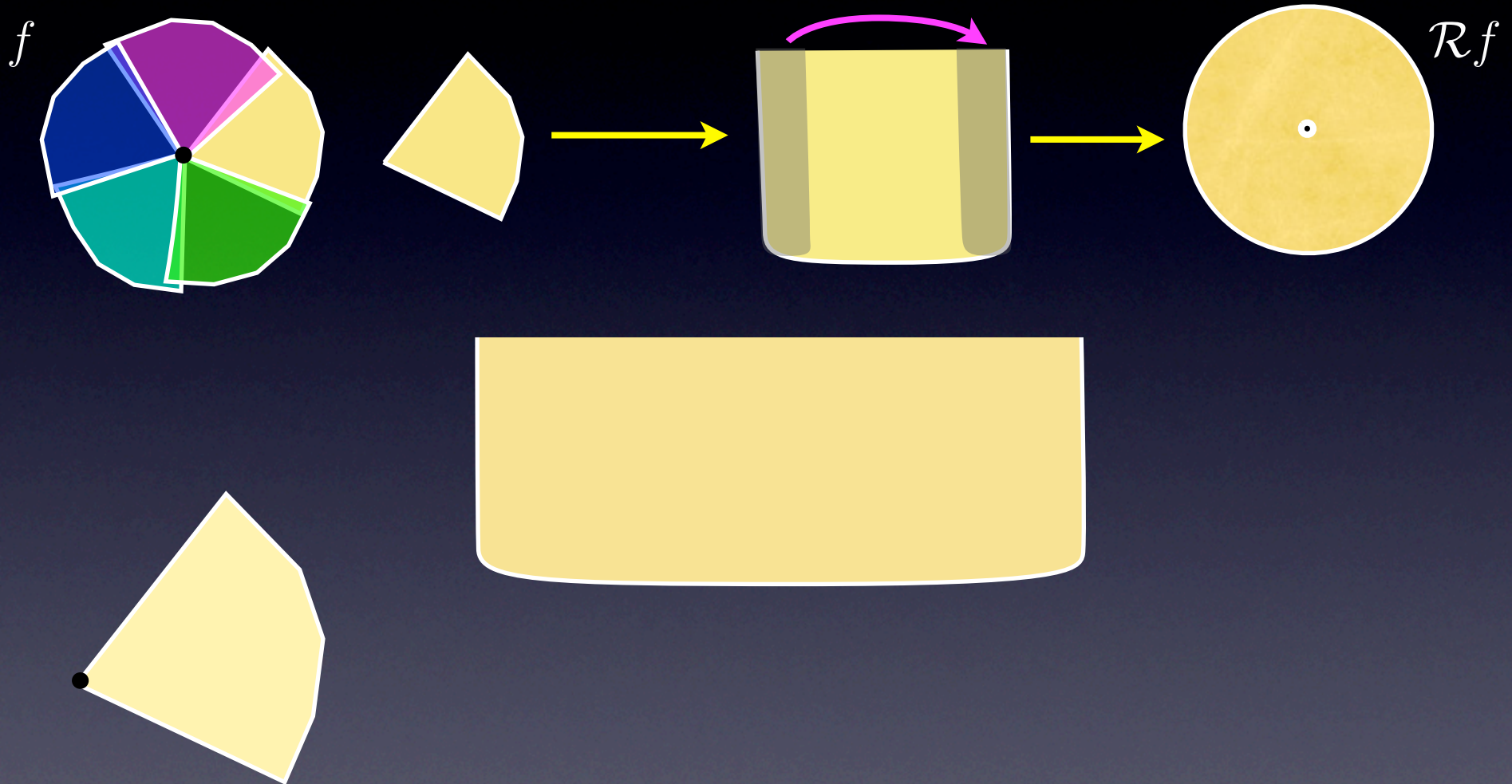
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

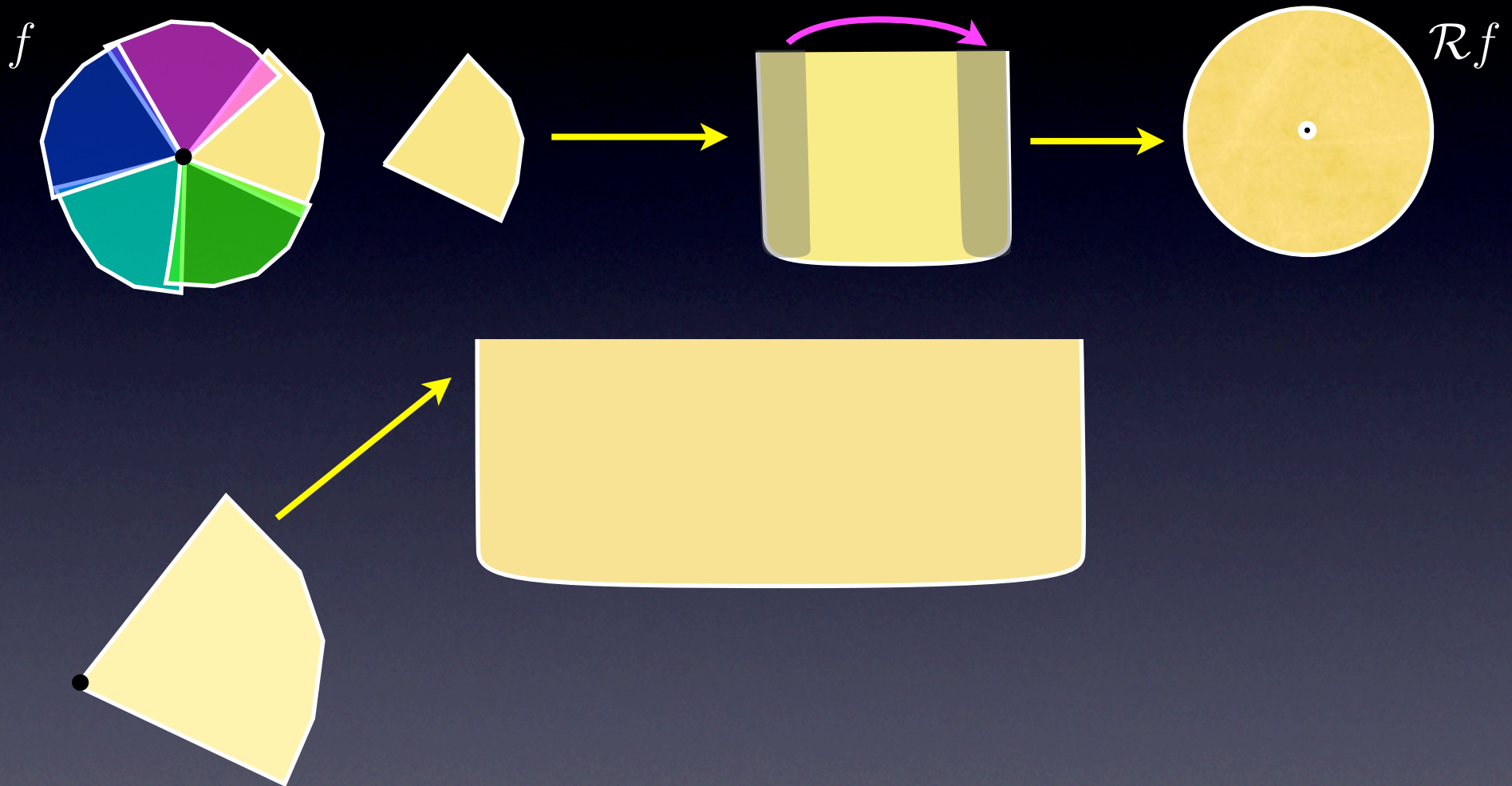
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

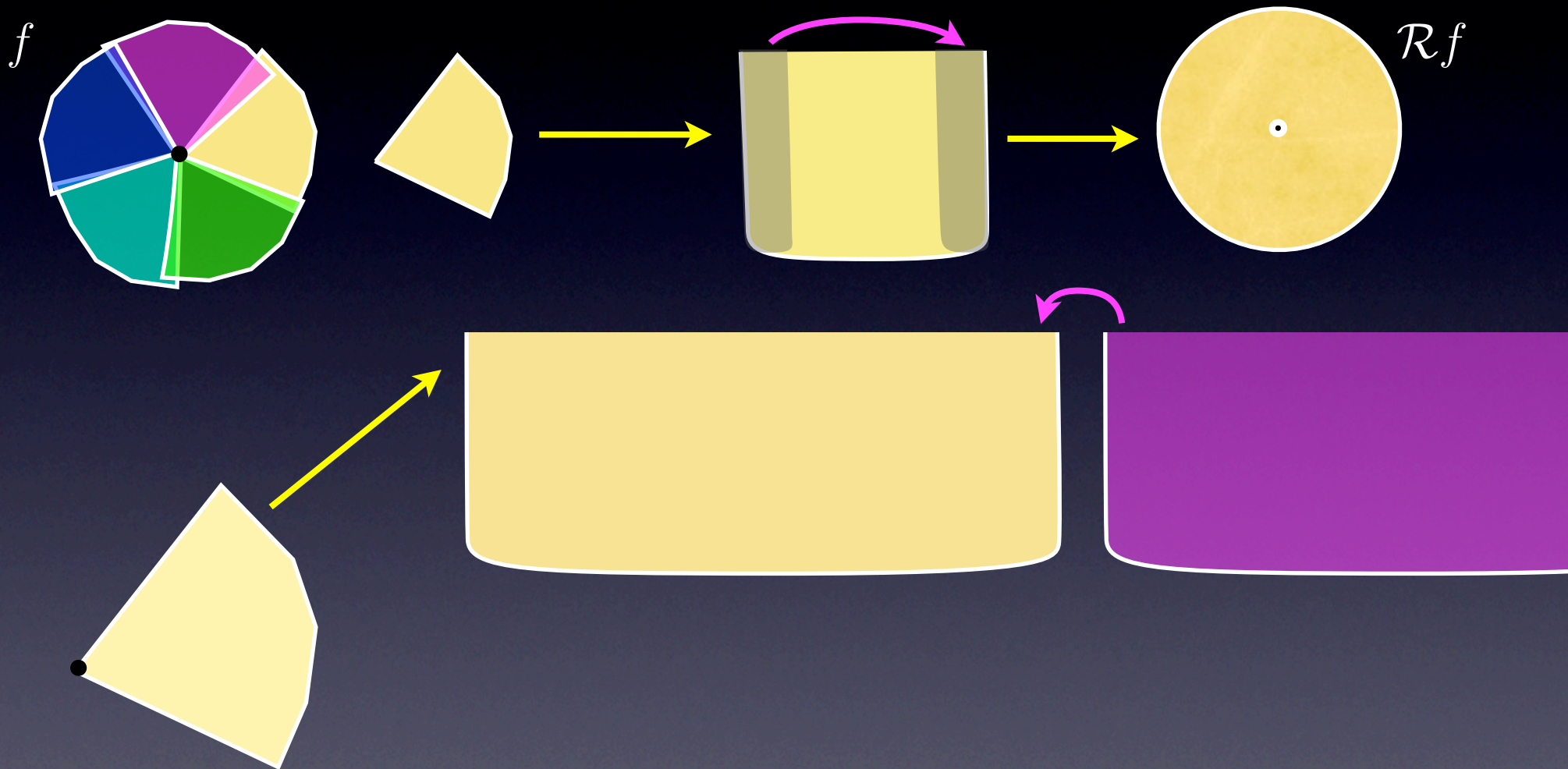
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

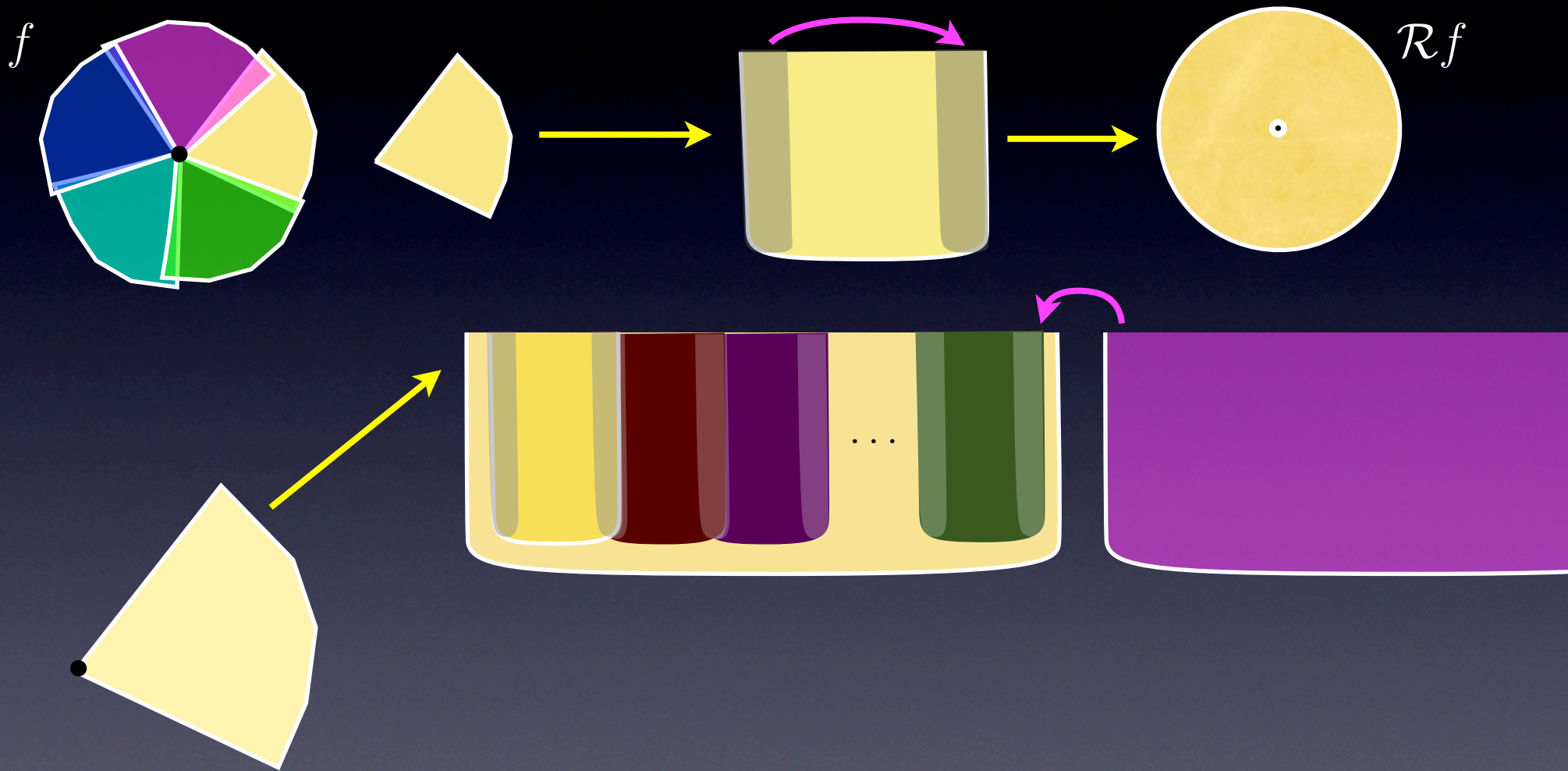
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

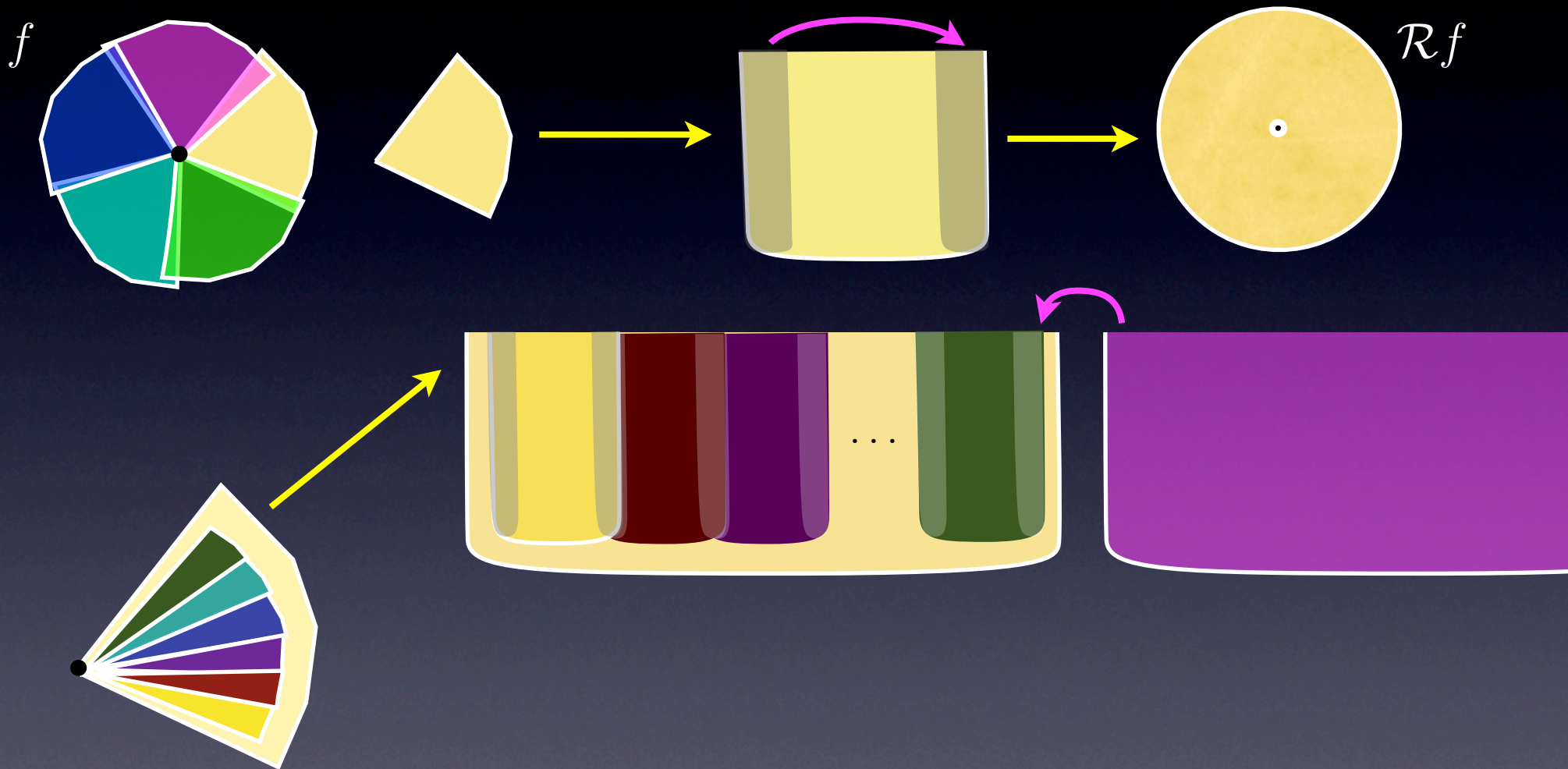
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

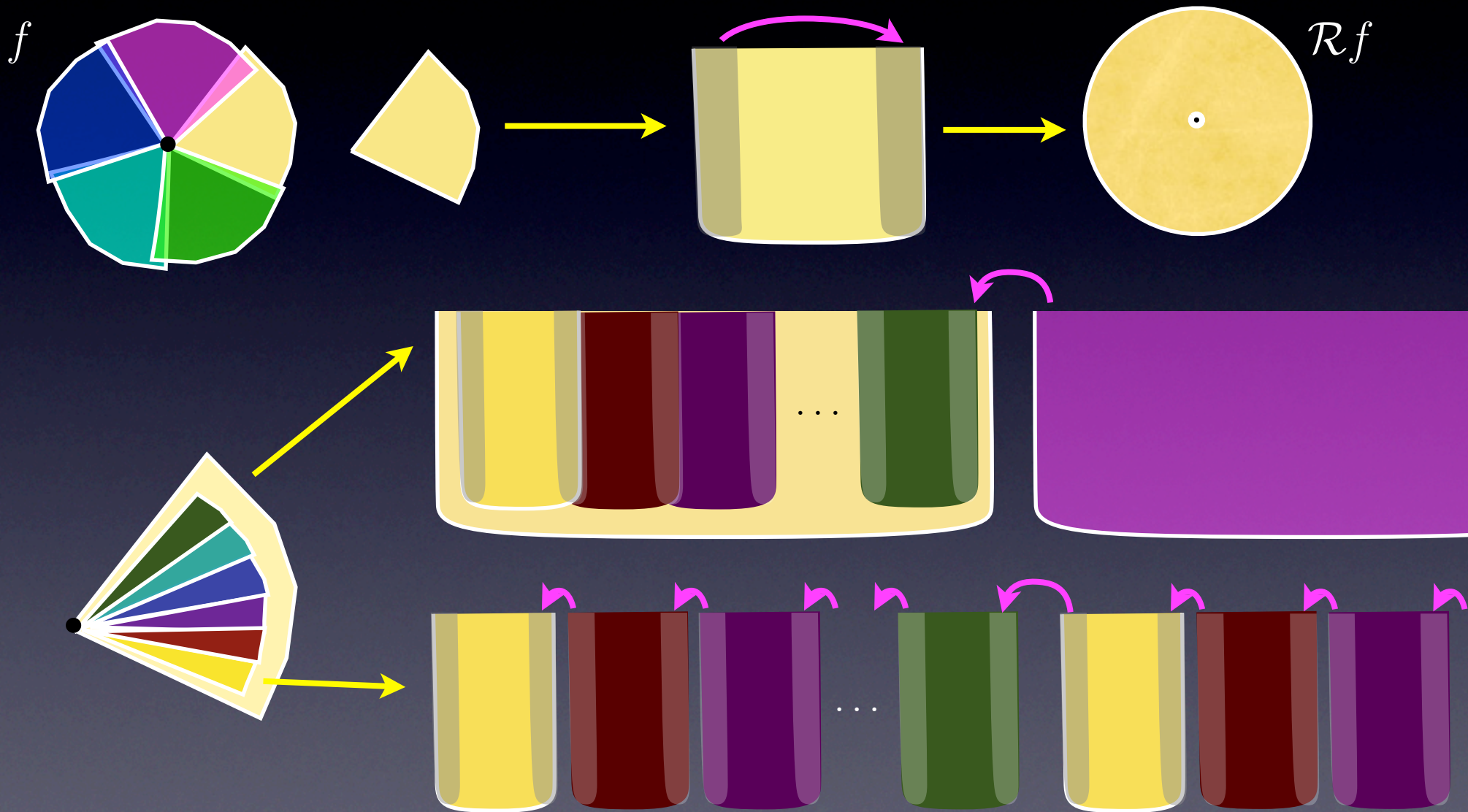
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

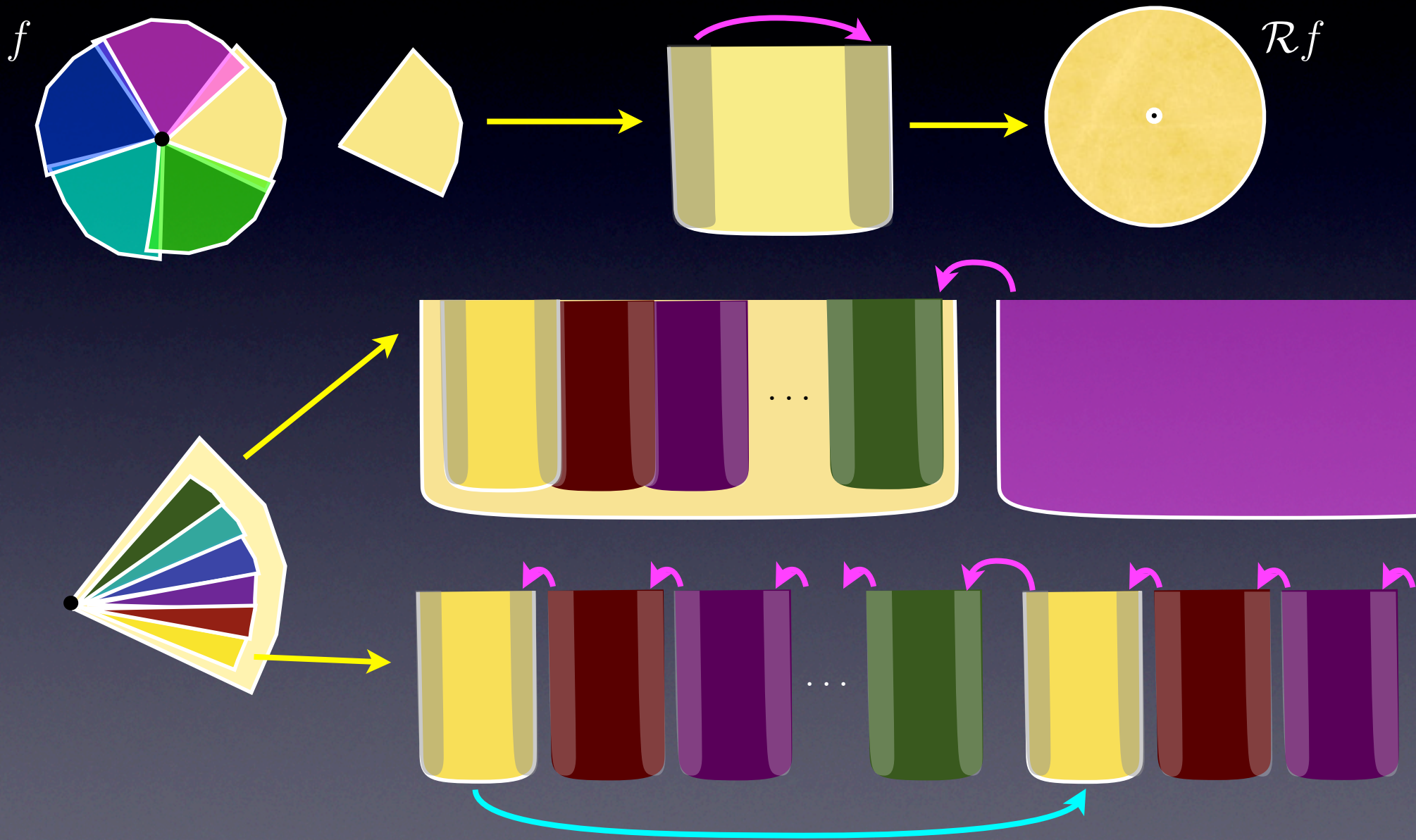
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

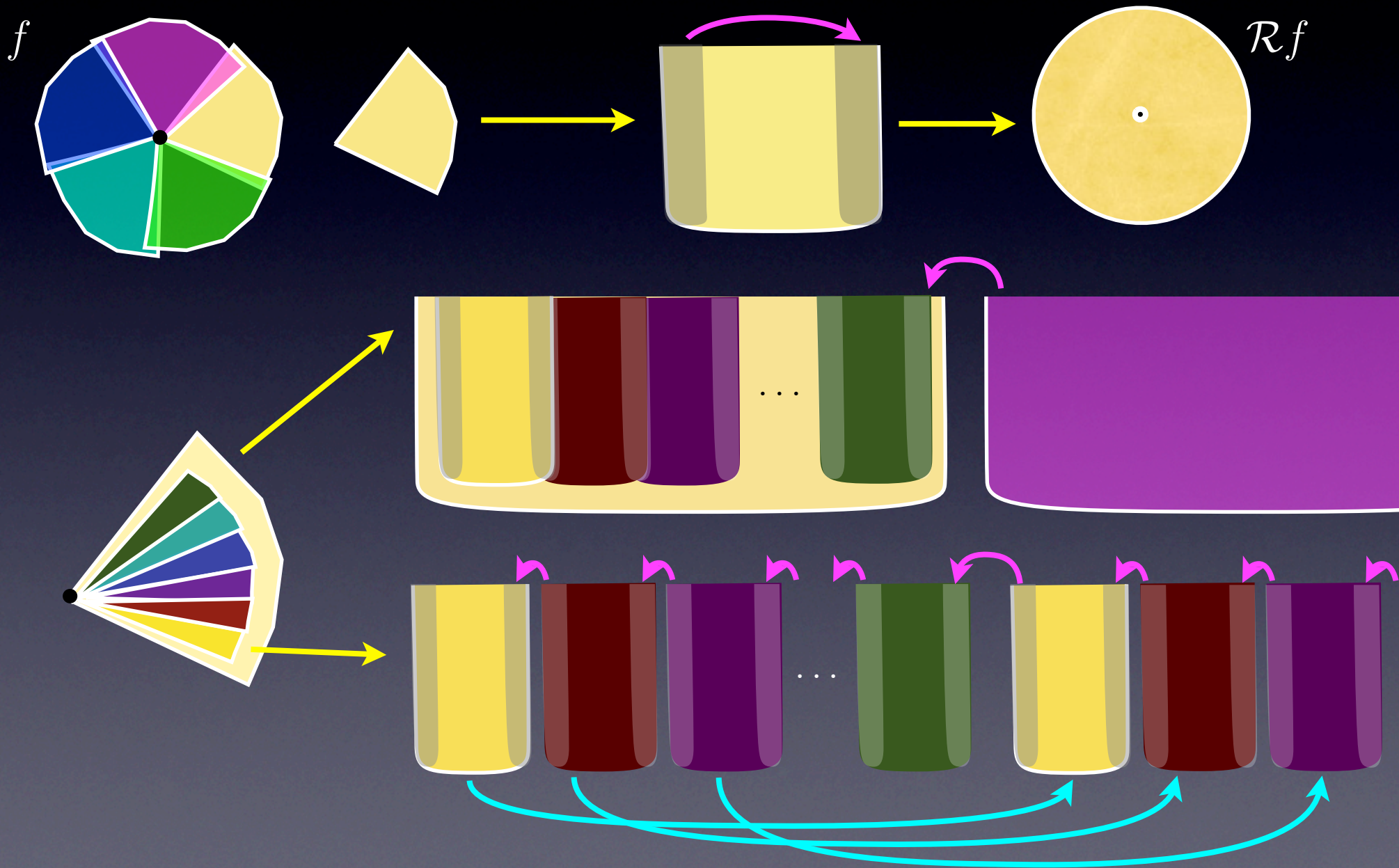
For simplicity, let us consider the sector renormalization:



Applications of Inou-S. Theorem

For applications, we need to see how to reconstruct f from its renormalizations.

For simplicity, let us consider the sector renormalization:



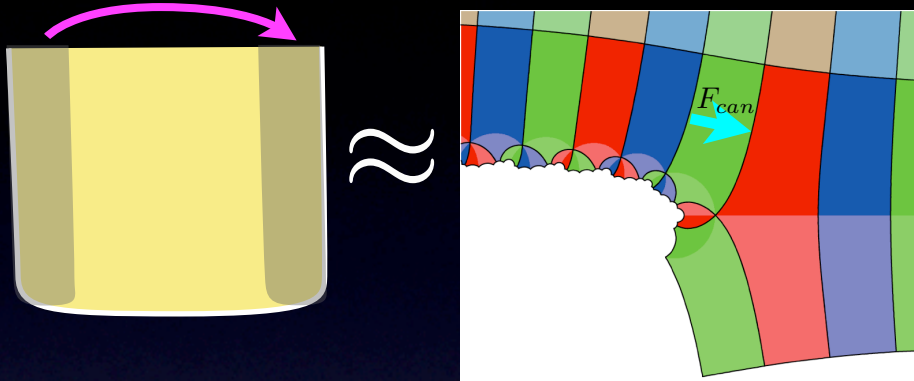
A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as

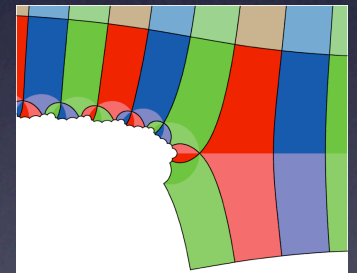
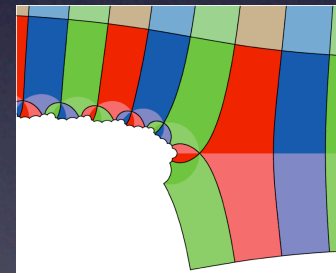
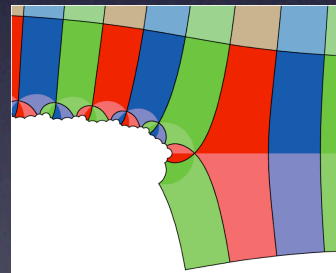
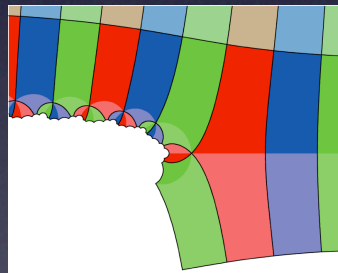
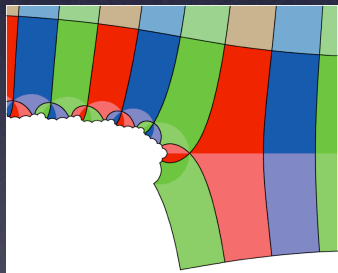
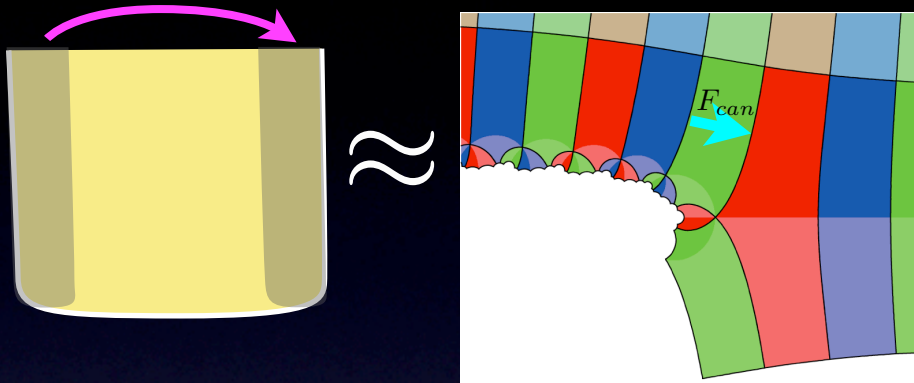
A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as



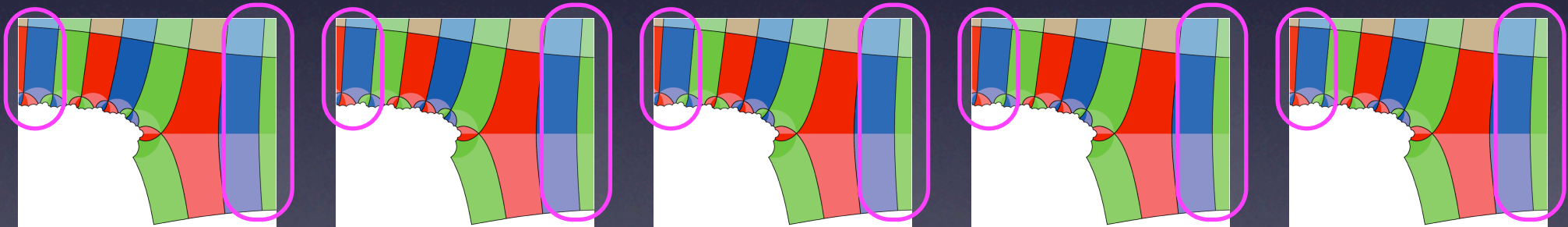
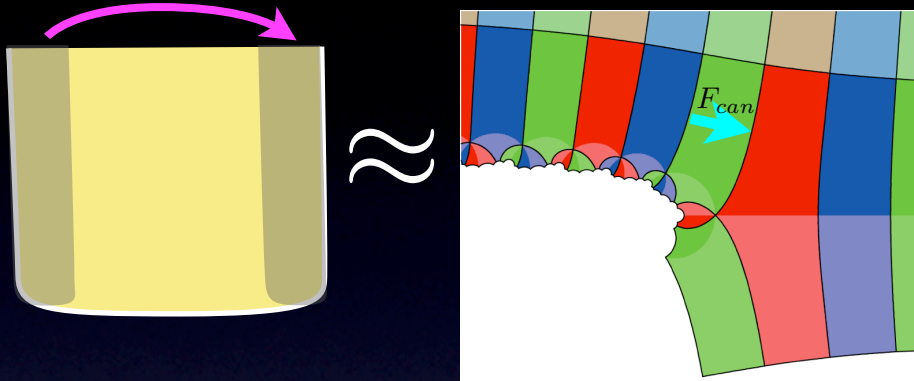
A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as



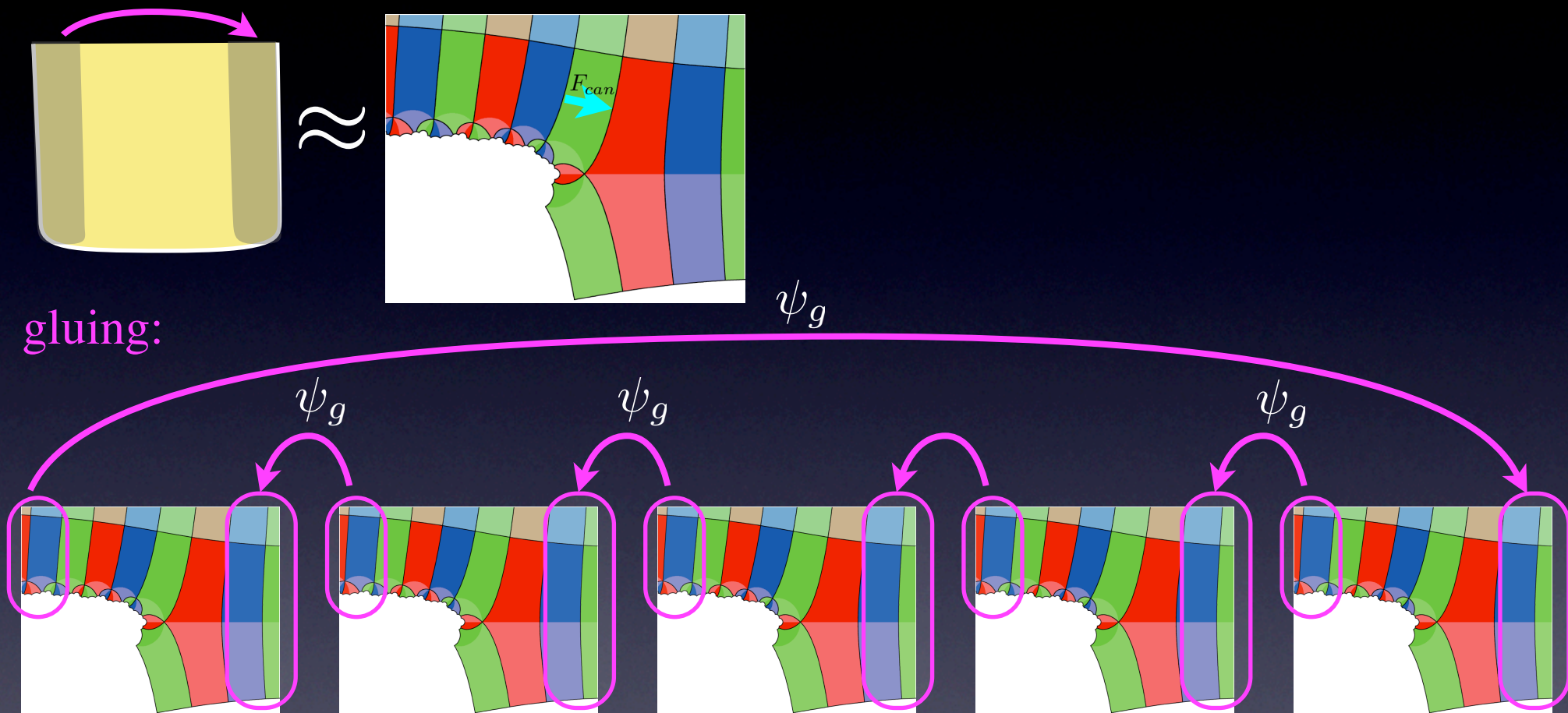
A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as



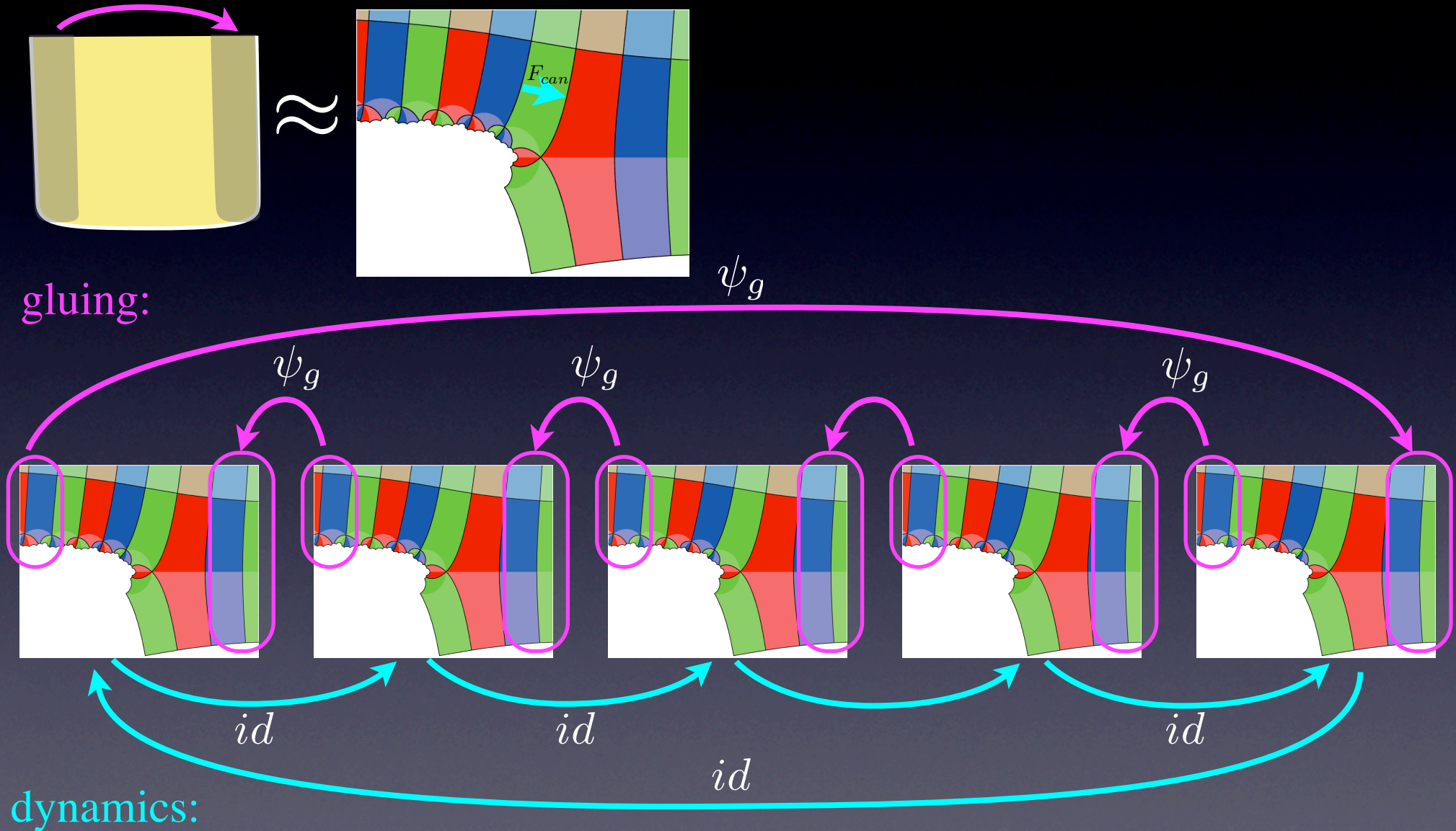
A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as



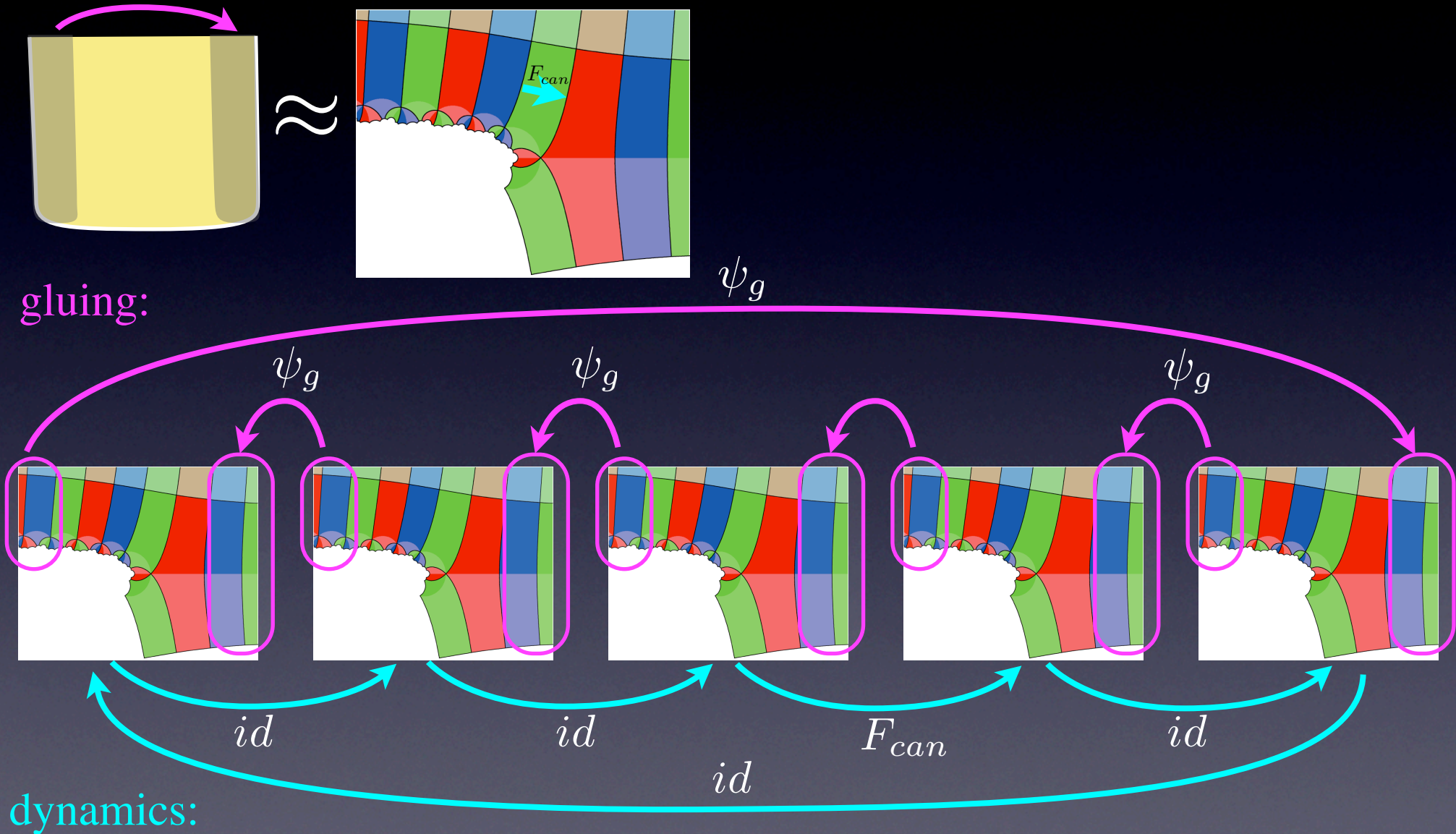
A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as



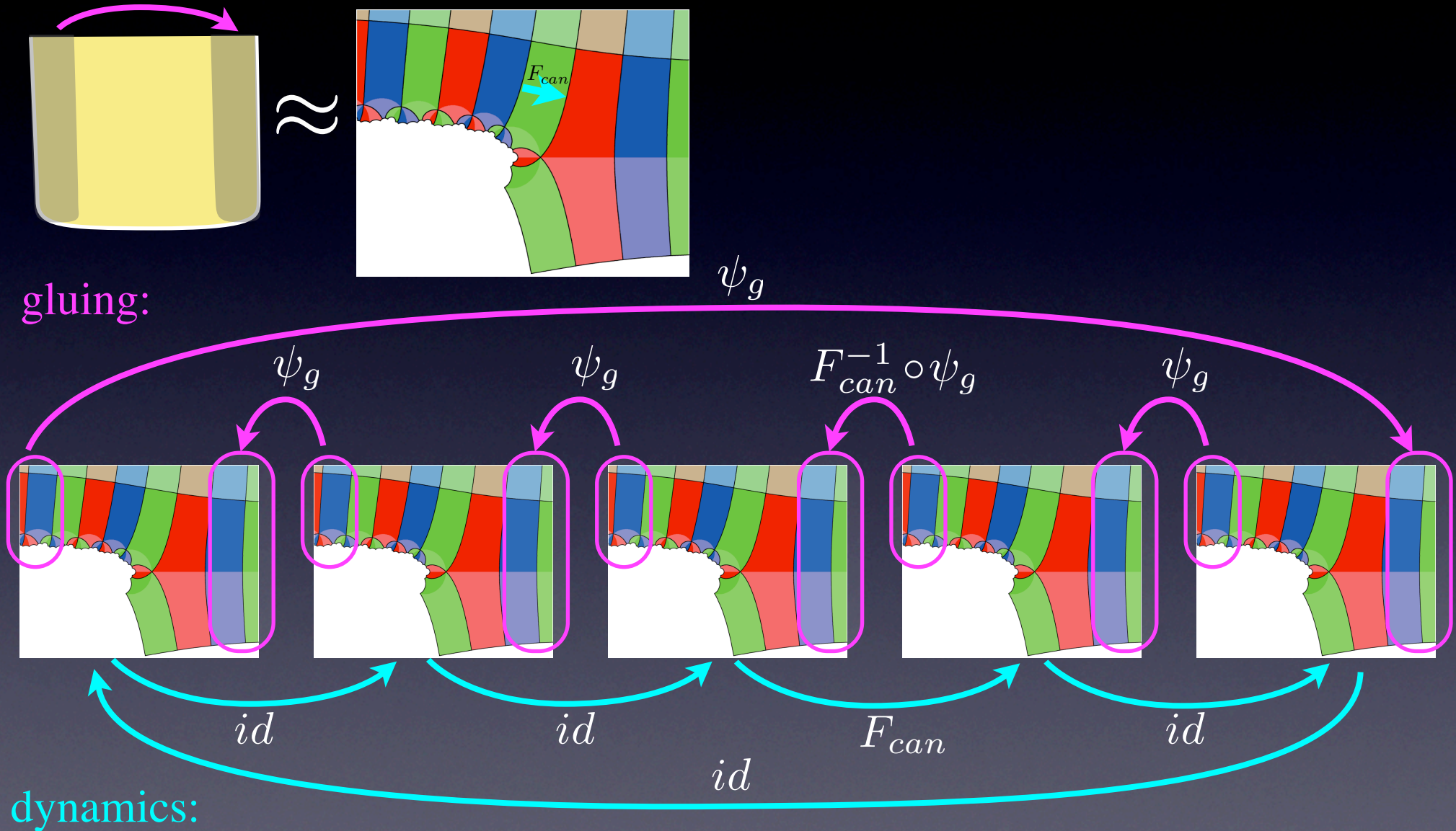
A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as



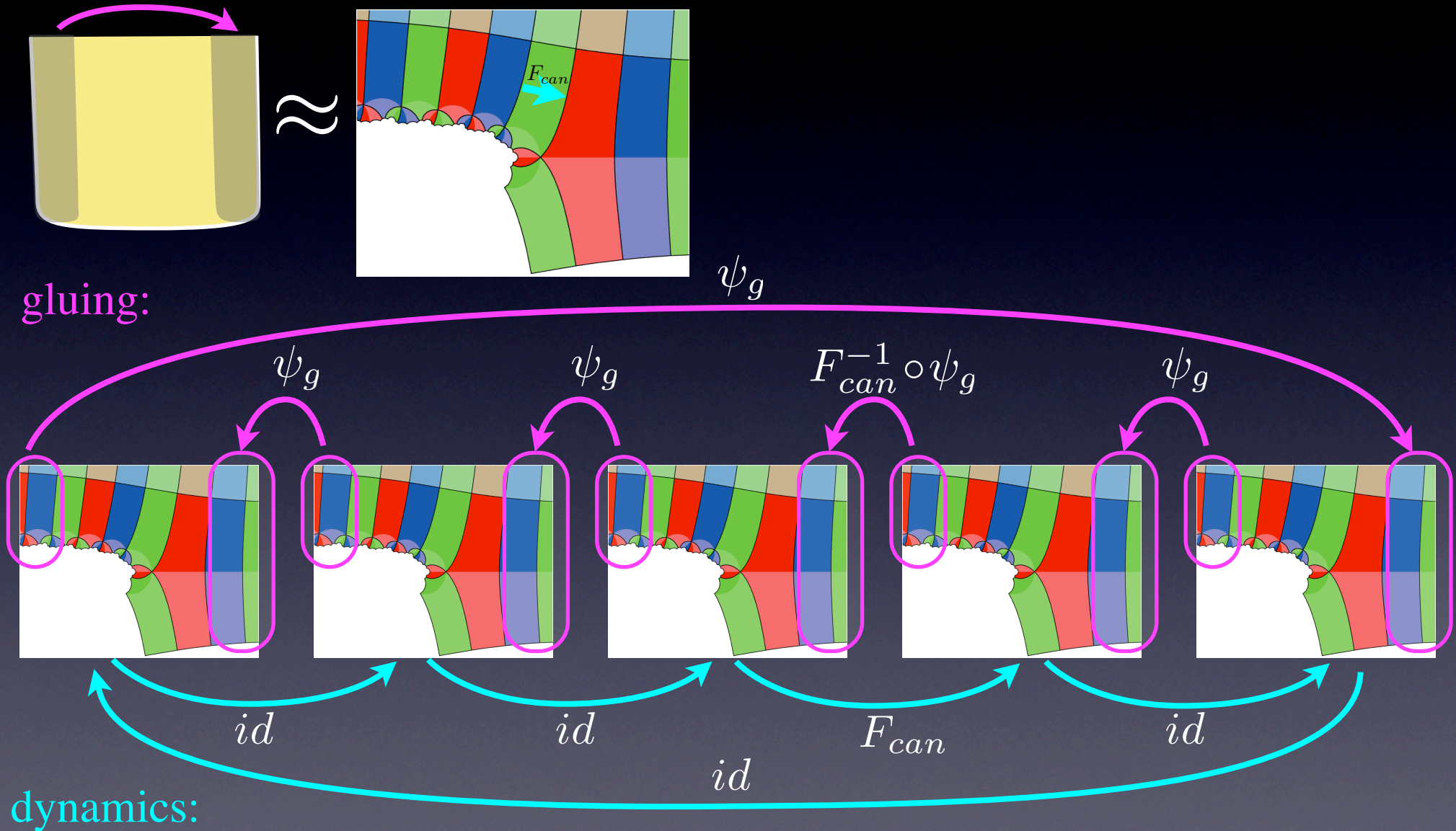
A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as



A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

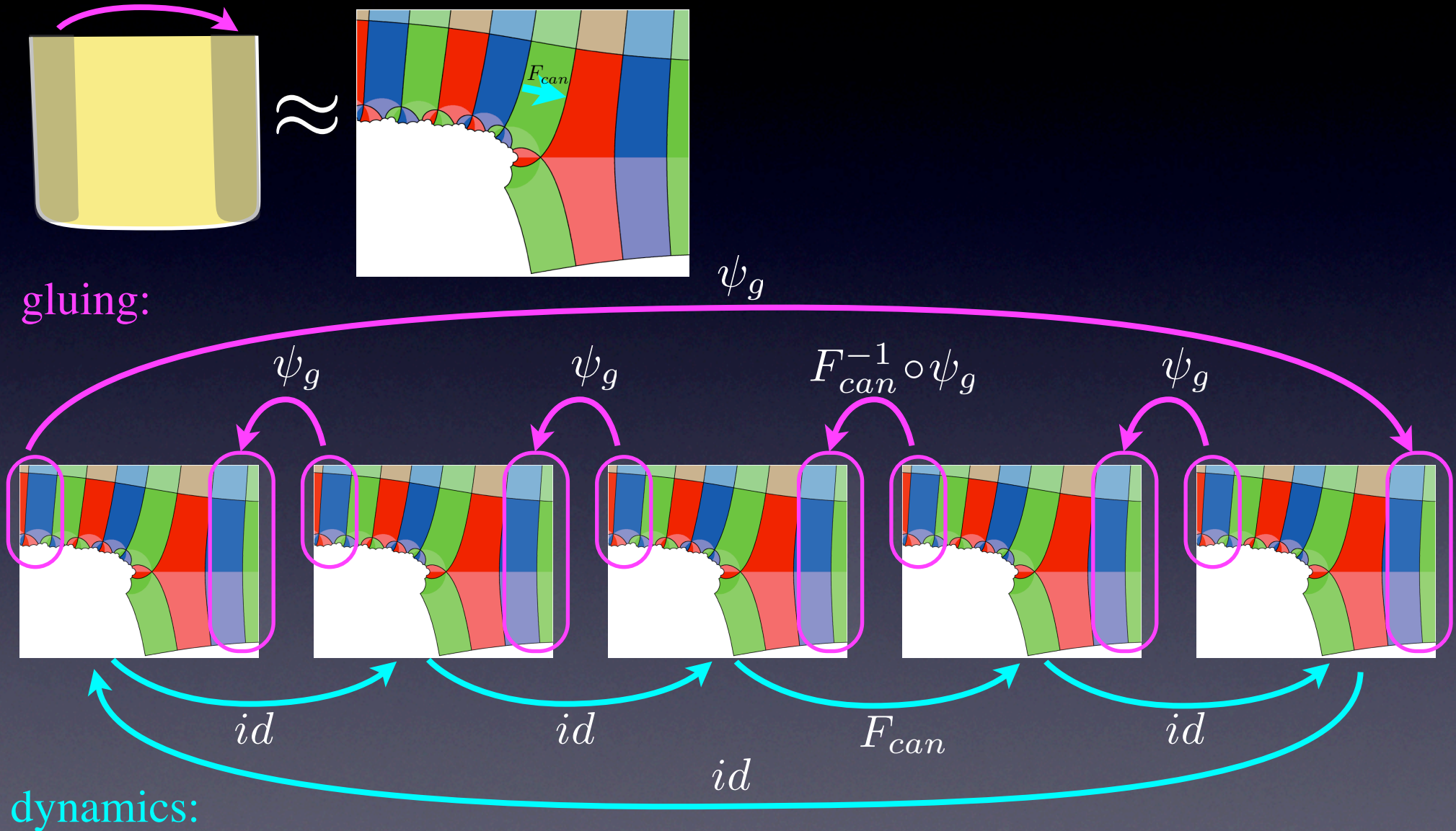
In the case of near-parabolic renormalization, chart can be taken as



1. well-defined after gluing

A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

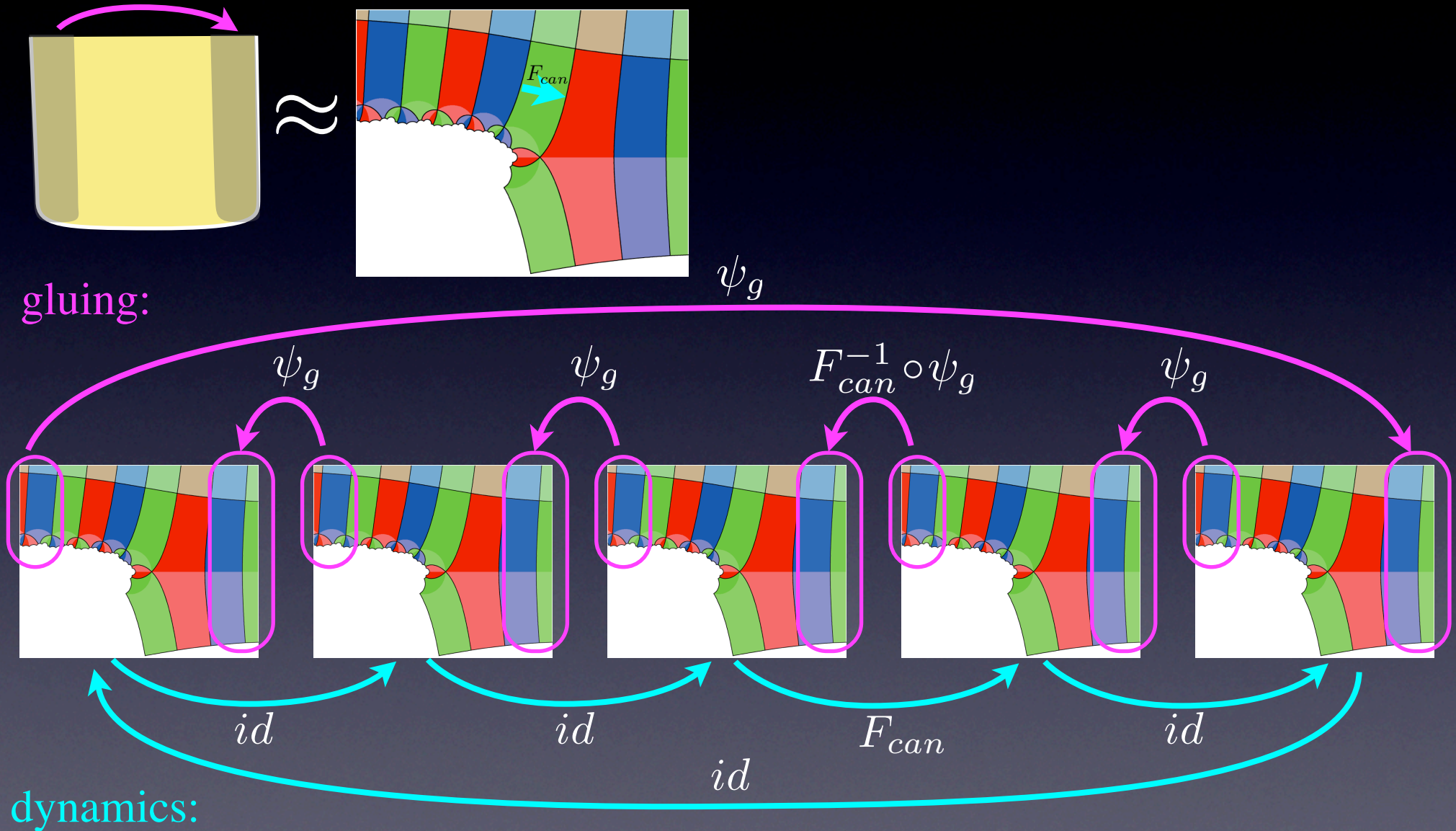
In the case of near-parabolic renormalization, chart can be taken as



1. well-defined after gluing
2. return map is F_{can} modulo ψ_g

A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

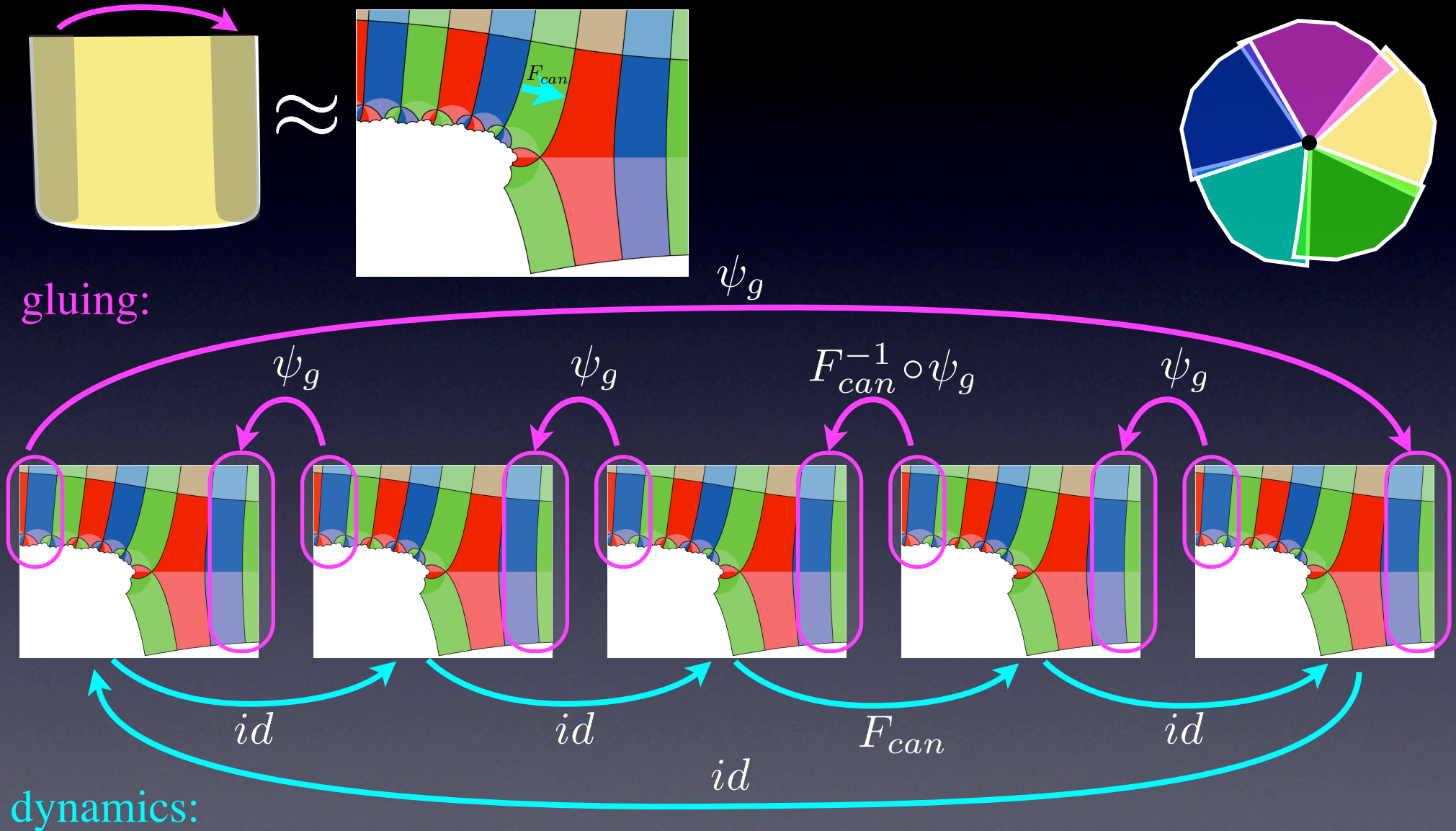
In the case of near-parabolic renormalization, chart can be taken as



1. well-defined after gluing
2. return map is F_{can} modulo ψ_g
3. this picture embeds into f

A punctured neighborhood of the fixed point is covered by *dynamical charts*. model maps on the charts and consistent gluings.

In the case of near-parabolic renormalization, chart can be taken as



1. well-defined after gluing
2. return map is F_{can} modulo ψ_g
3. this picture embeds into f

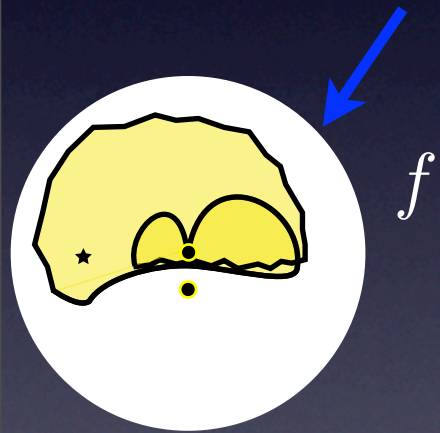
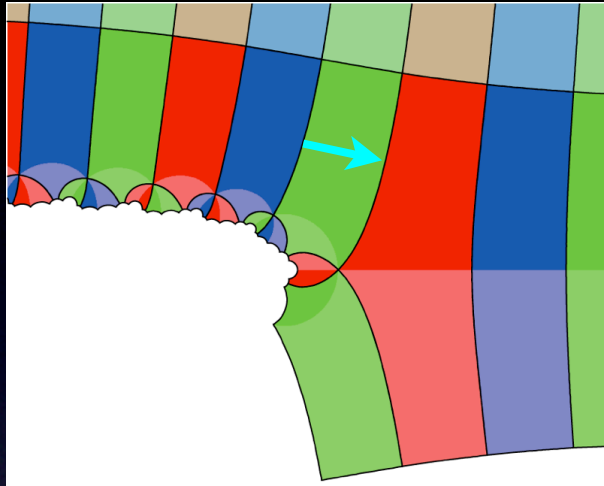
Construction of successive charts

Construction of successive charts



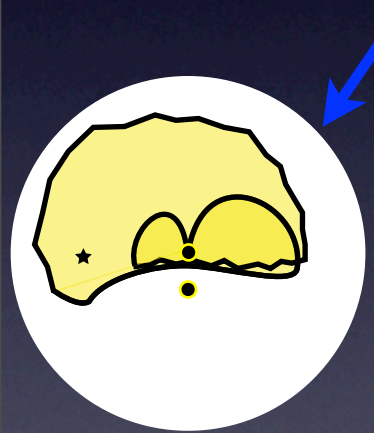
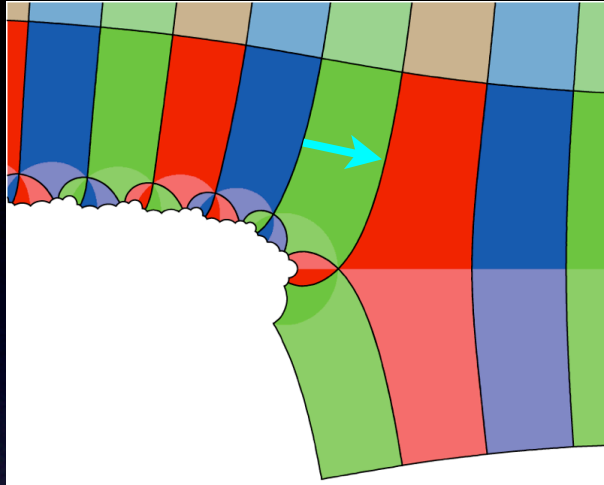
Construction of successive charts

Ω_f



Construction of successive charts

Ω_f



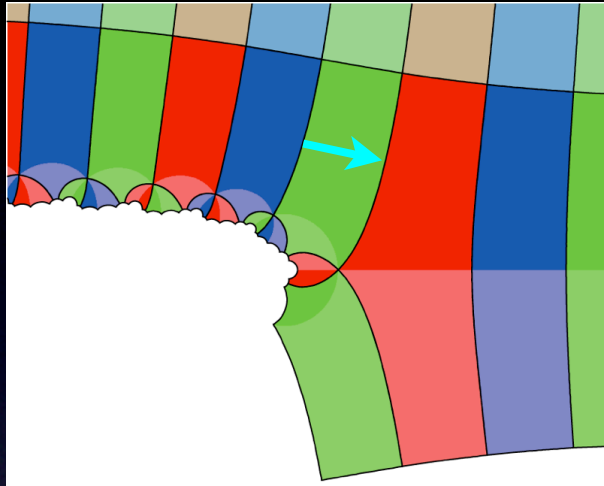
f



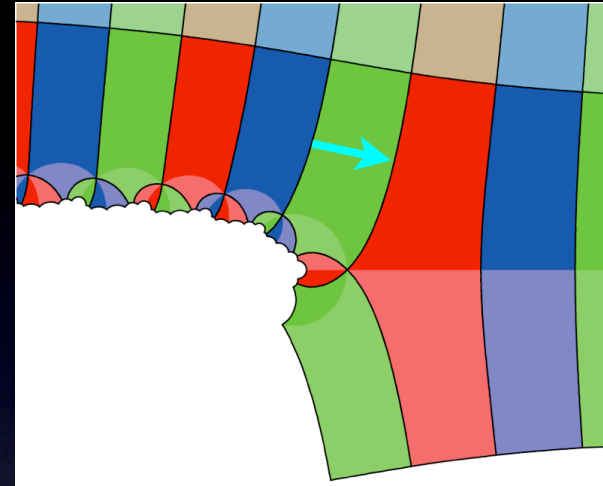
$\mathcal{R}f$

Construction of successive charts

Ω_f



$\Omega_{\mathcal{R}f}$



f

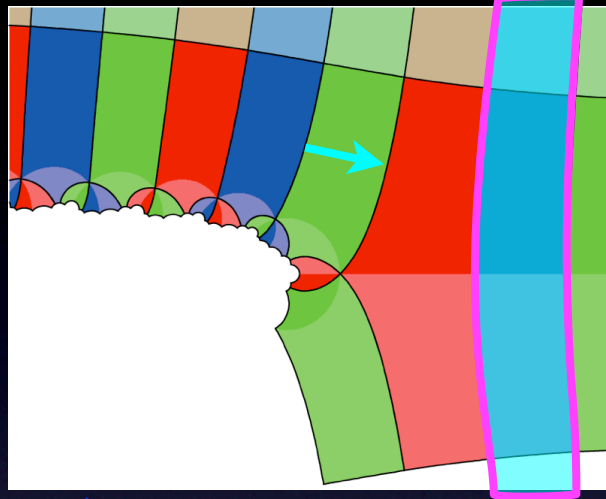


$\mathcal{R}f$

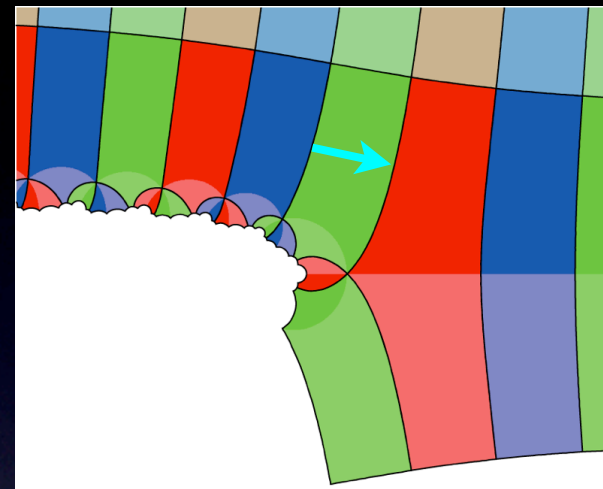


Construction of successive charts

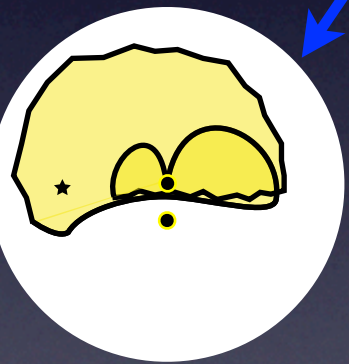
Ω_f



$\Omega_{\mathcal{R}f}$



f



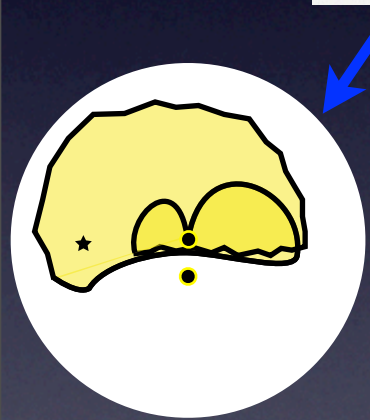
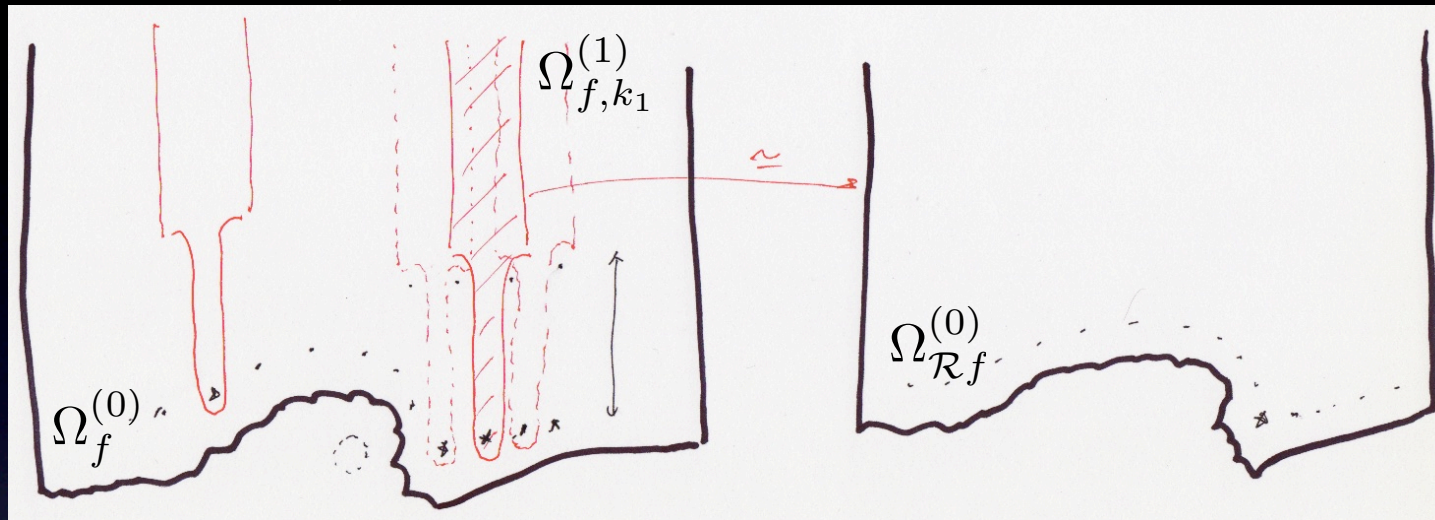
$\mathcal{R}f$



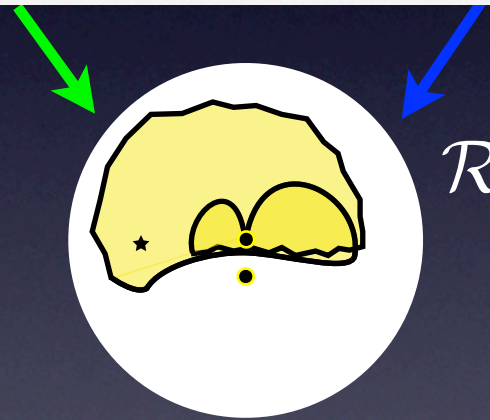
Construction of successive charts

Ω_f

$\Omega_{\mathcal{R}f}$



f

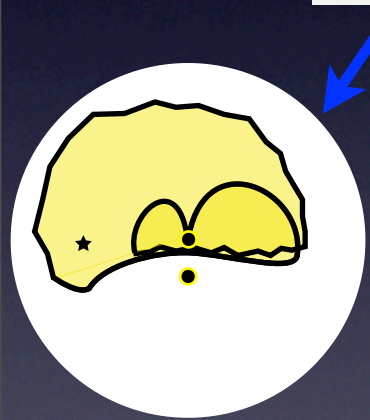
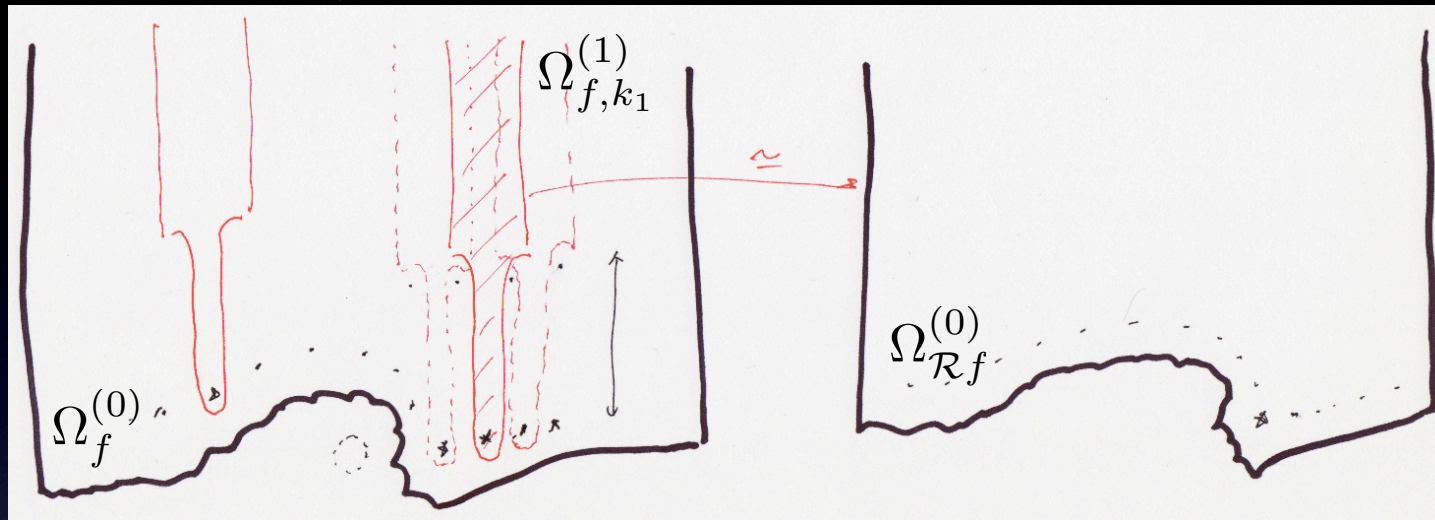


$\mathcal{R}f$

Construction of successive charts

Ω_f

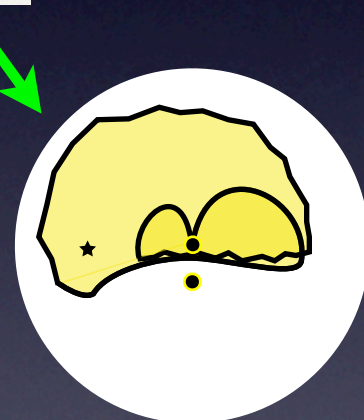
$\Omega_{\mathcal{R}f}$



f



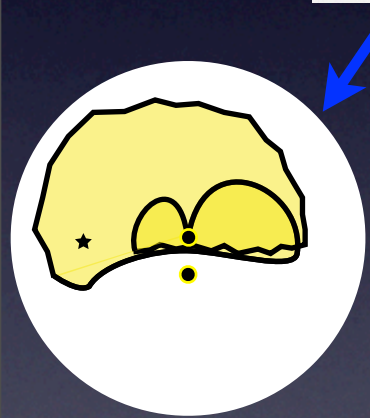
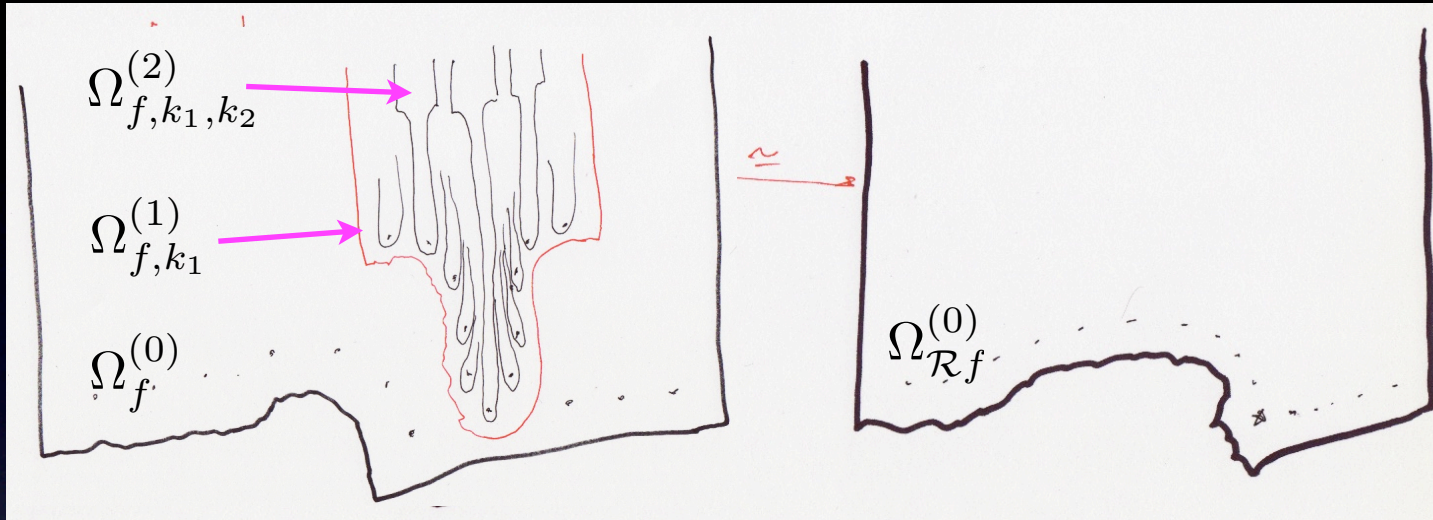
$\mathcal{R}f$



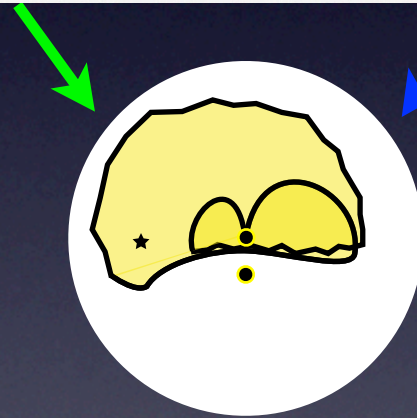
Construction of successive charts

Ω_f

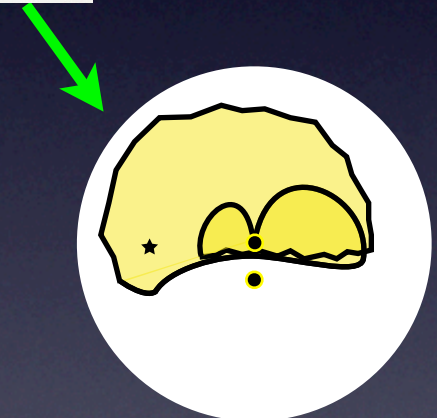
$\Omega_{\mathcal{R}f}$



f



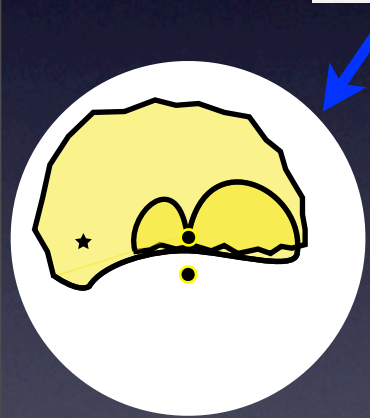
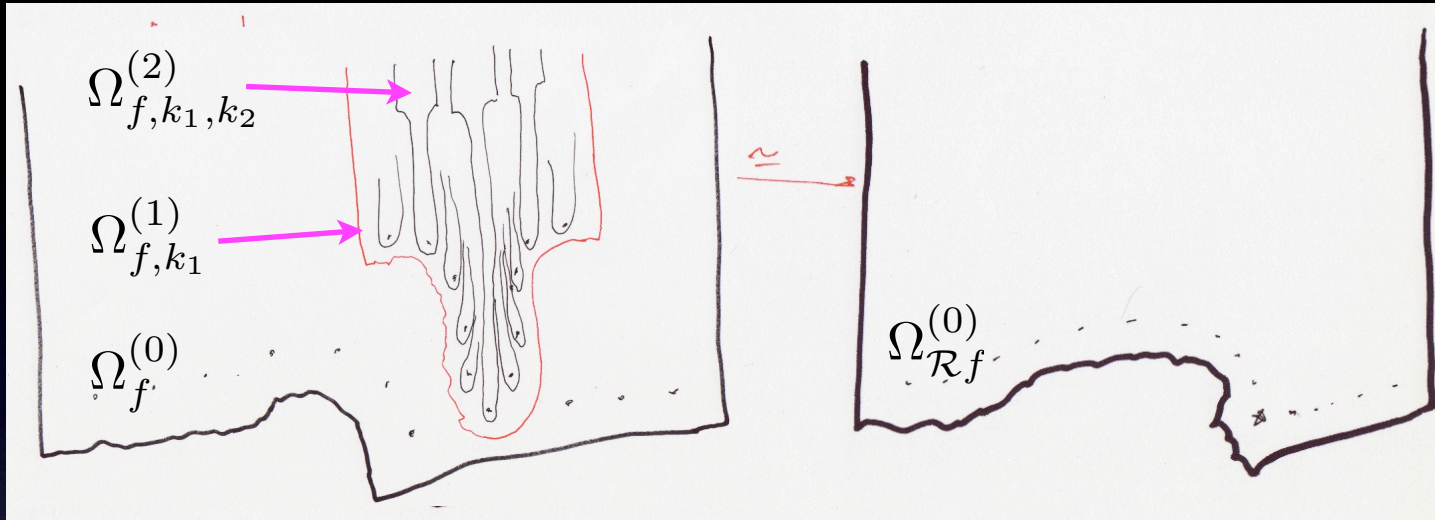
$\mathcal{R}f$



Construction of successive charts

Ω_f

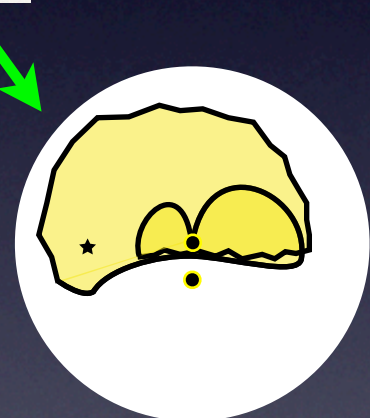
$\Omega_{\mathcal{R}f}$



f



$\mathcal{R}f$

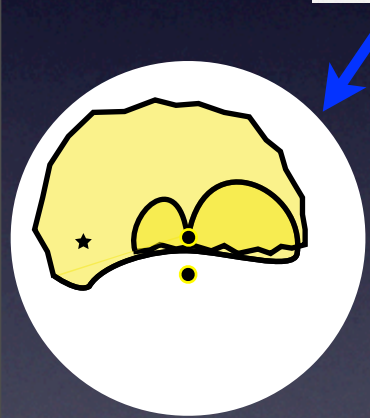
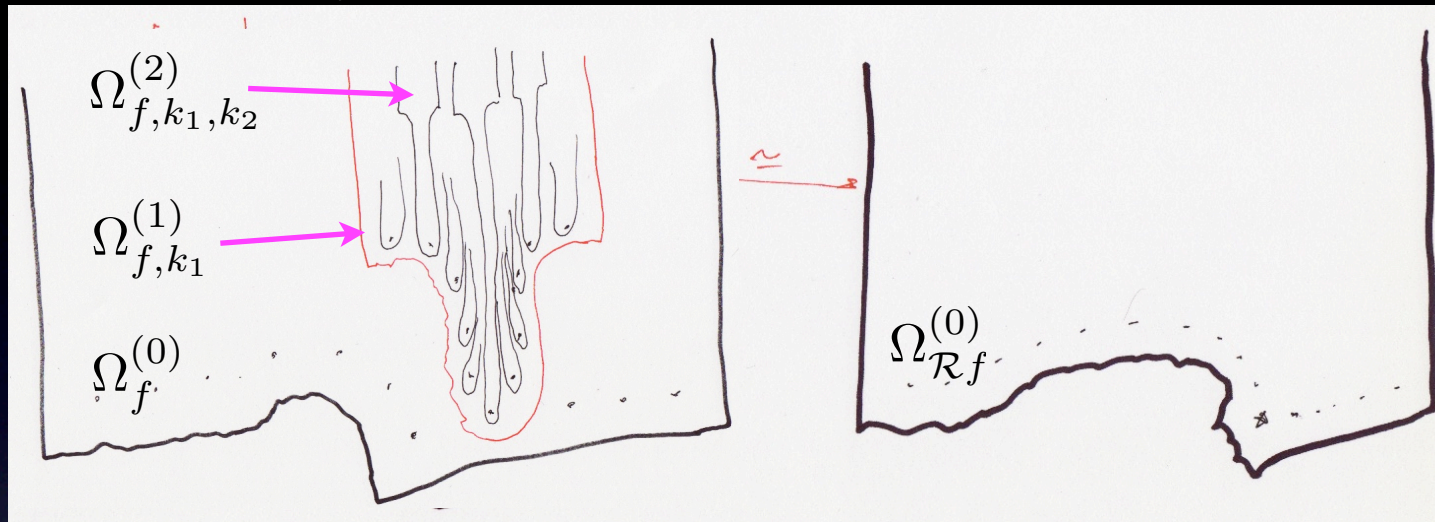


$$\Omega^{(0)} \supset \Omega_{k_1}^{(1)} \supset \dots \supset \Omega_{k_1, k_2, \dots, k_n}^{(n)} \supset \Omega_{k_1, k_2, \dots, k_n, k_{n+1}}^{(n+1)} \supset \dots$$

Construction of successive charts

Ω_f

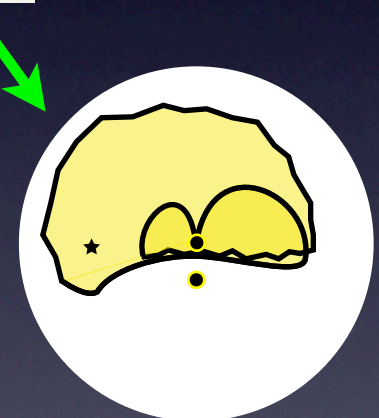
$\Omega_{\mathcal{R}f}$



f



$\mathcal{R}f$



$$\Omega^{(0)} \supset \Omega_{k_1}^{(1)} \supset \dots \supset \Omega_{k_1, k_2, \dots, k_n}^{(n)} \supset \Omega_{k_1, k_2, \dots, k_n, k_{n+1}}^{(n+1)} \supset \dots$$

each $\Omega_{k_1, k_2, \dots, k_n}^{(n)}$ is isomorphic to truncated checkerboard pattern $\Omega_{\mathcal{R}^n f}$

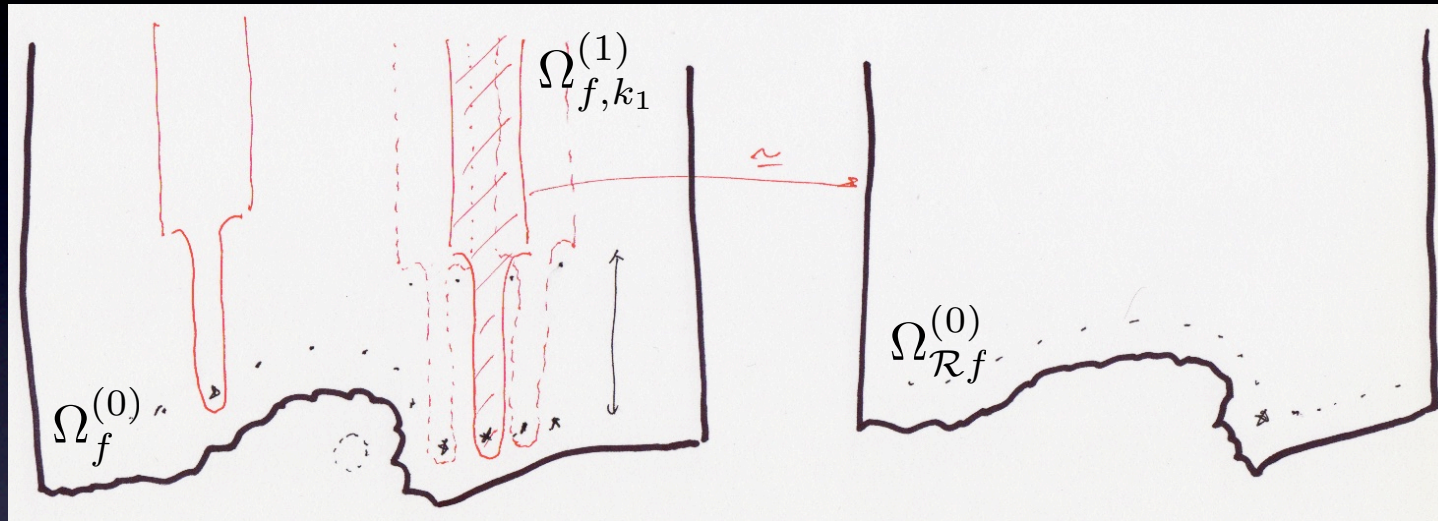
Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.

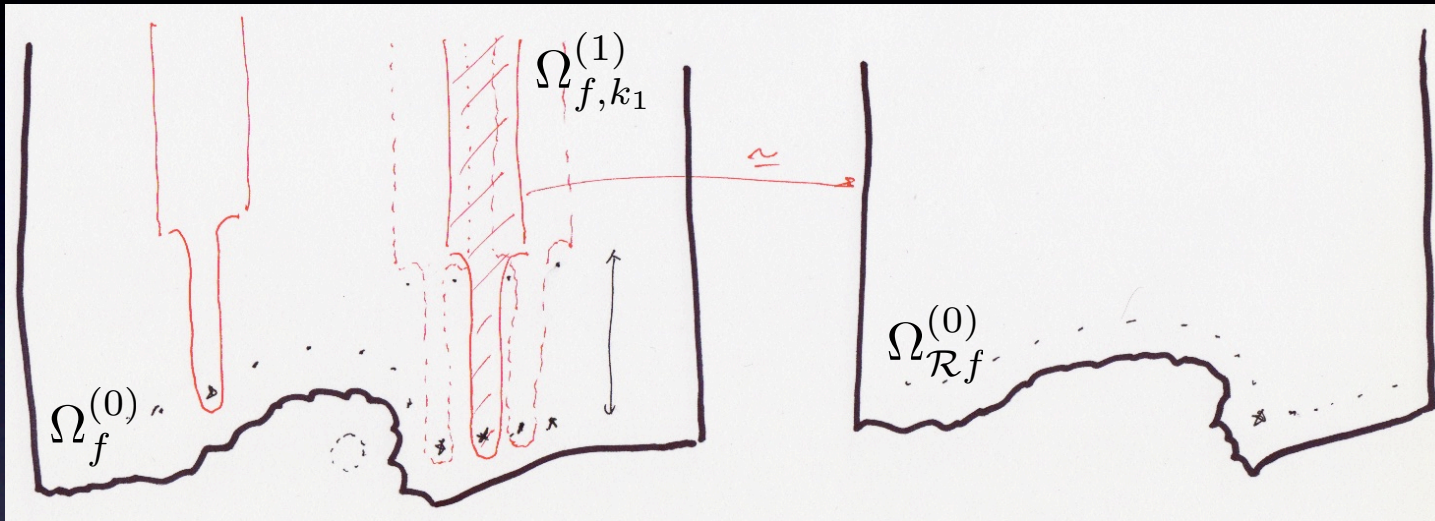
Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.



Theorem: If $f(z) = e^{2\pi i\alpha} z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

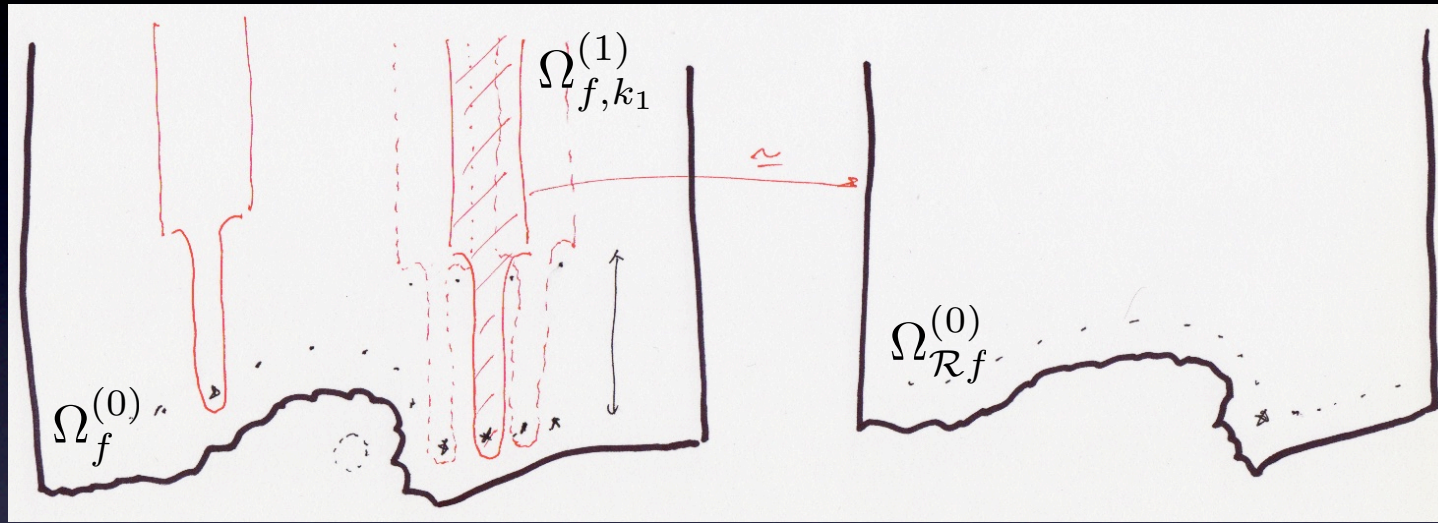
Idea of the proof: Construct approximating curves and show that they converge.



—
approximation
of bdry of
Siegel disk

Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.

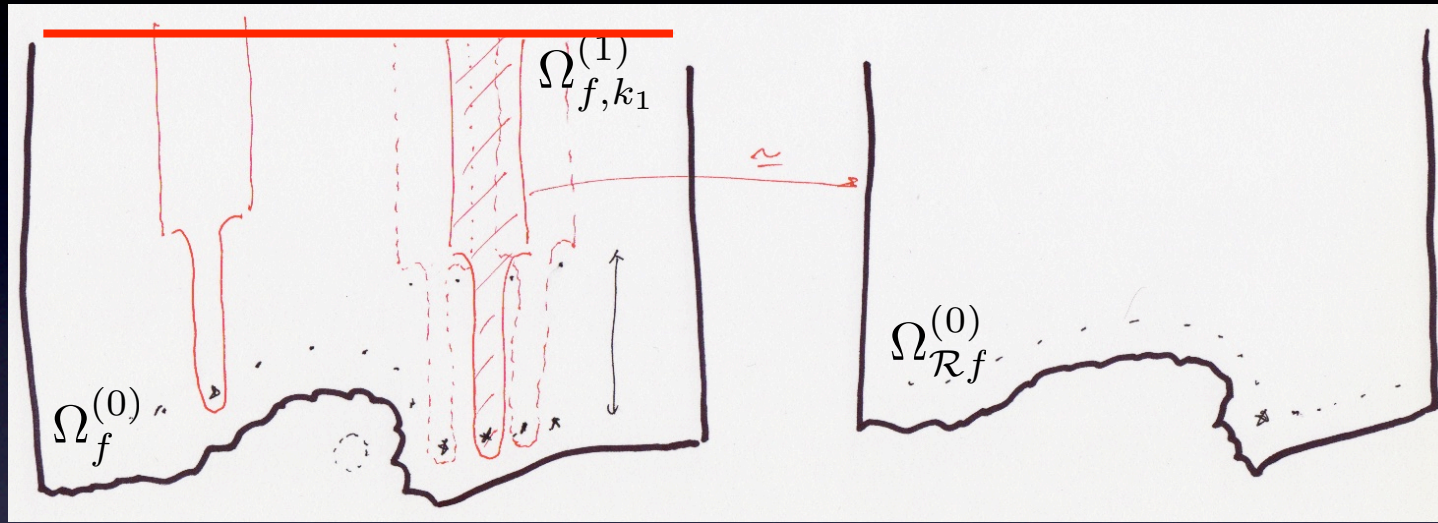


—
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.

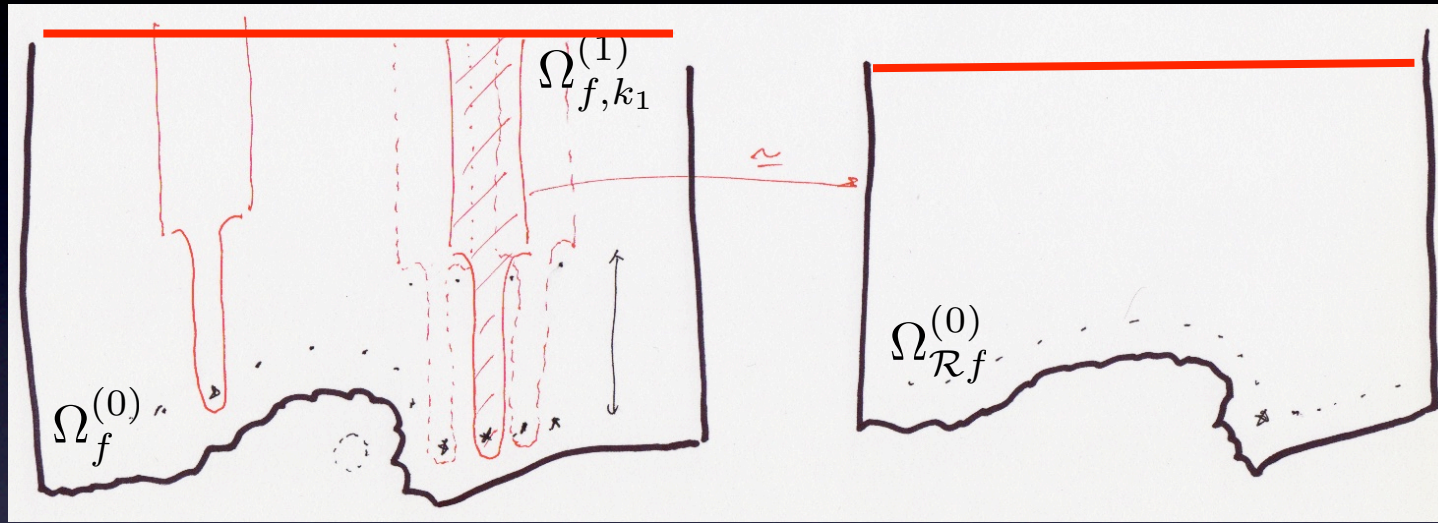


approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.

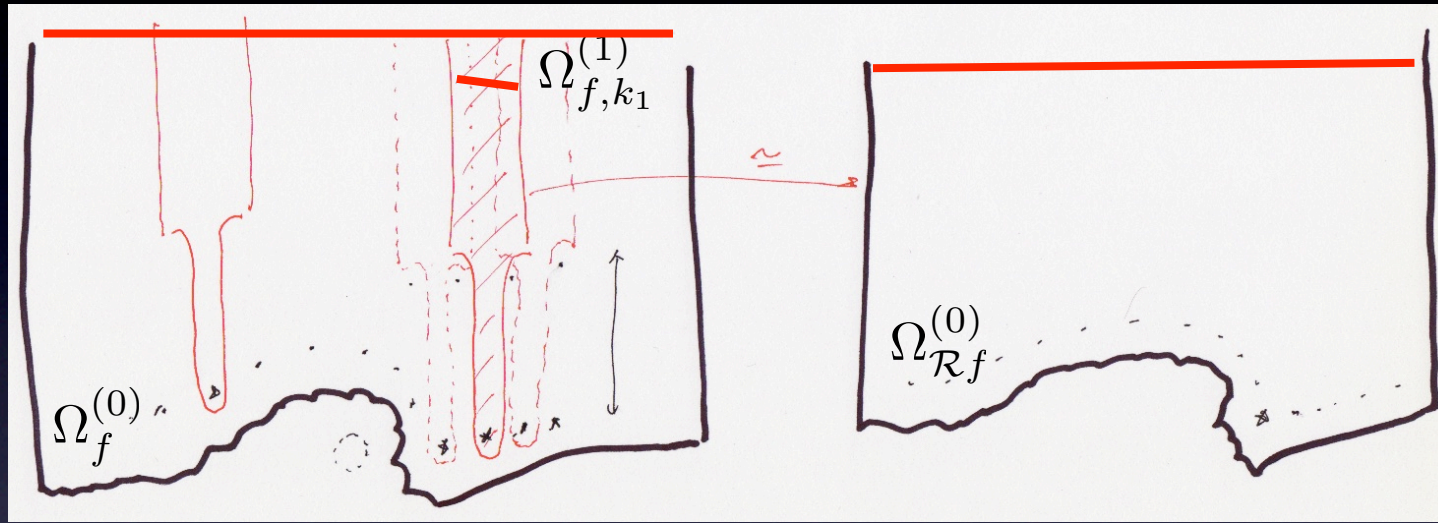


—
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.

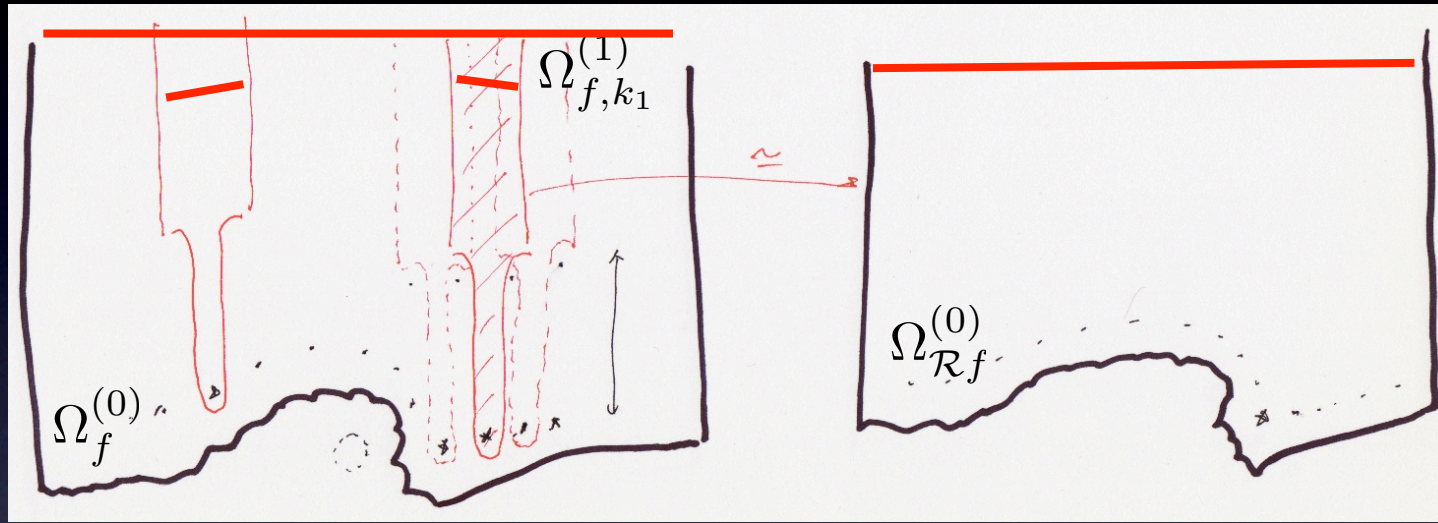


—
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.

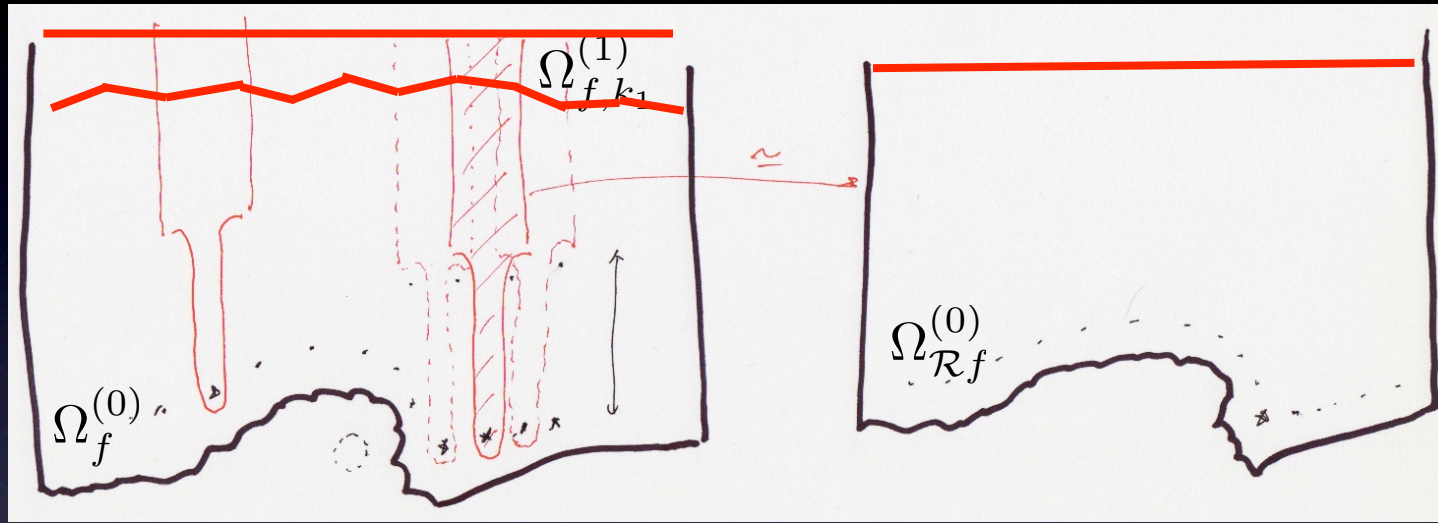


—
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

Theorem: If $f(z) = e^{2\pi i\alpha} z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.

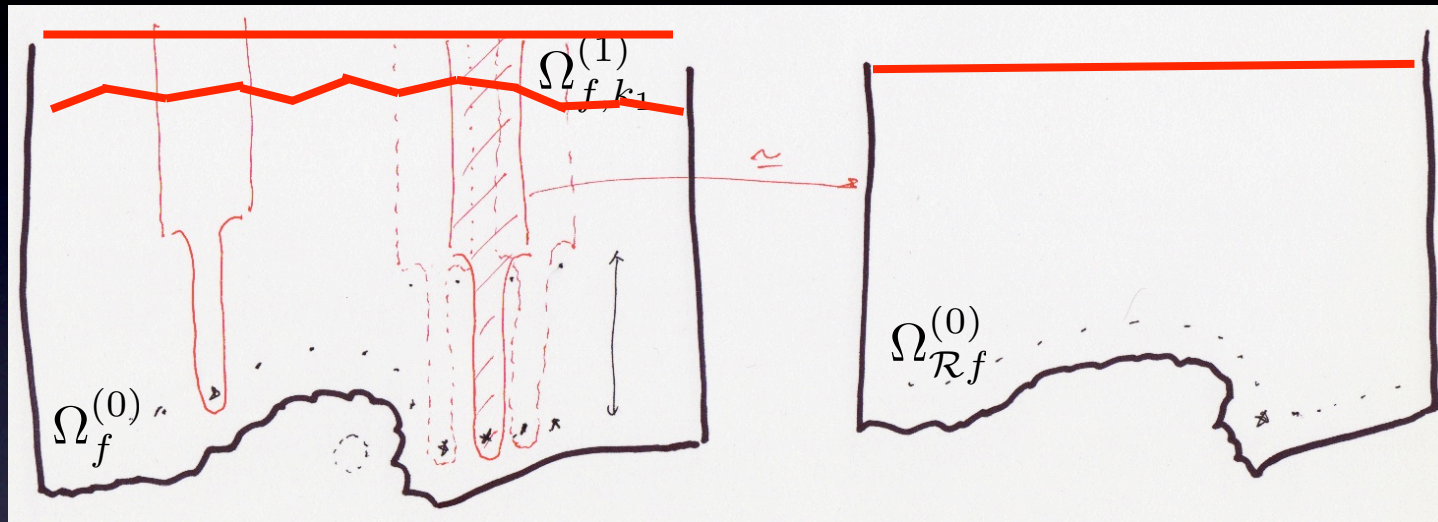


—
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

Theorem: If $f(z) = e^{2\pi i\alpha} z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.



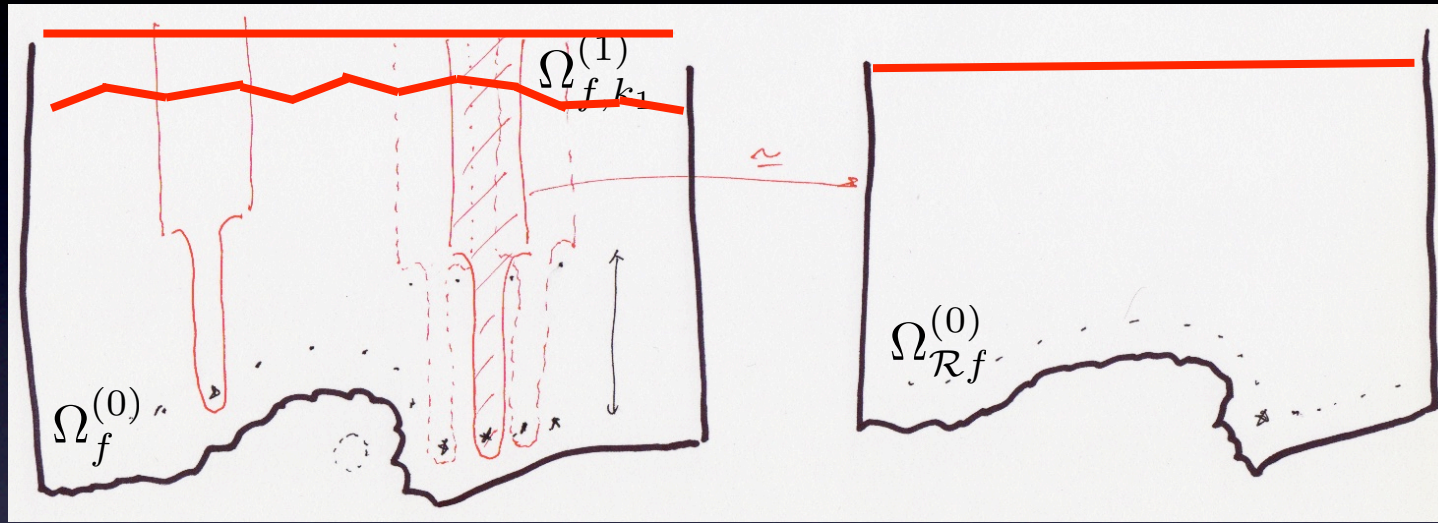
—
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

height of the approx. curve: $h_n = B(\alpha_n)$

Theorem: If $f(z) = e^{2\pi i\alpha} z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.



—
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

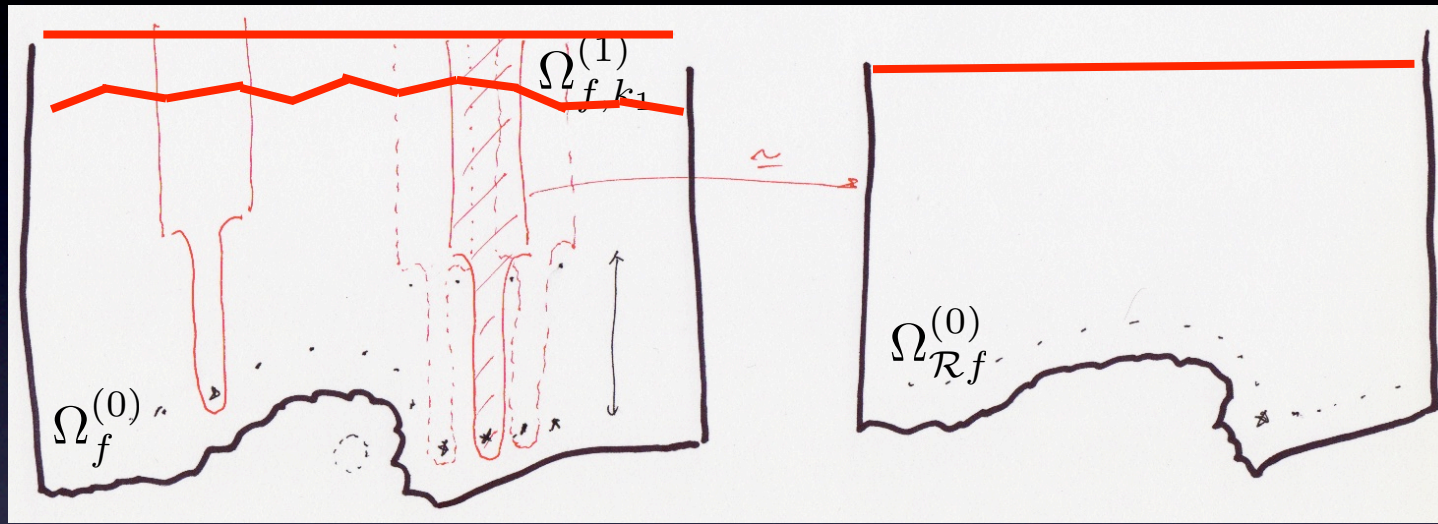
height of the approx. curve: $h_n = B(\alpha_n)$

mapping $\Omega_{f,k_1}^{(1)} \rightarrow \Omega_{Rf}^{(0)}$ is uniformly expanding

(with respect to Poincaré metrics of $\Omega_f^{(0)}$ and $\Omega_{Rf}^{(0)}$)

Theorem: If $f(z) = e^{2\pi i\alpha} z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.



—
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

height of the approx. curve: $h_n = B(\alpha_n)$

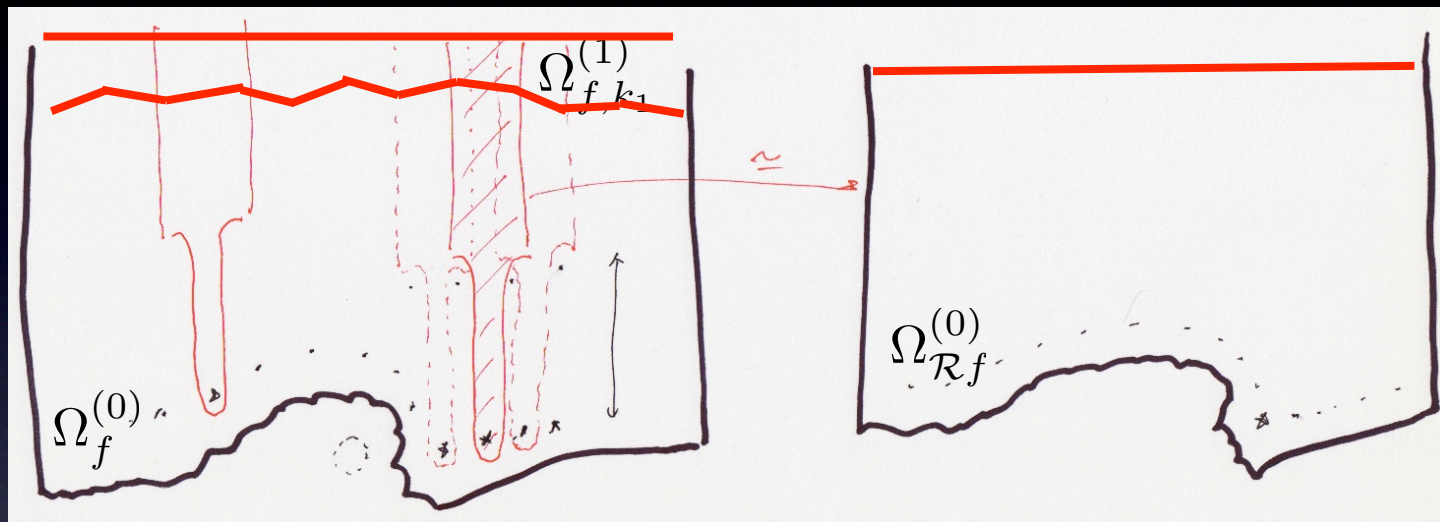
mapping $\Omega_{f,k_1}^{(1)} \rightarrow \Omega_{Rf}^{(0)}$ is uniformly expanding

(with respect to Poincaré metrics of $\Omega_f^{(0)}$ and $\Omega_{Rf}^{(0)}$)

One only needs to work in model space (i.e. truncated checkerboard pattern)

Theorem: If $f(z) = e^{2\pi i\alpha}z + z^2$ with α of sufficiently high type and Brjuno, then the boundary of Siegel disk is a Jordan curve.

Idea of the proof: Construct approximating curves and show that they converge.



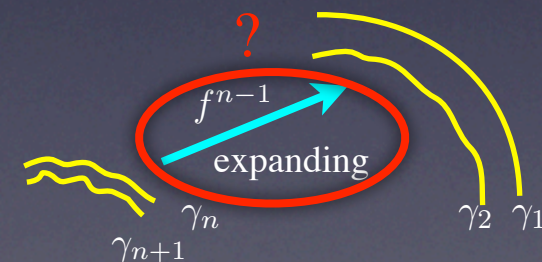
approximation
of bdry of
Siegel disk

Construct approximate boundary curves by joining segments in $\Omega_{f,k_1,\dots,k_n}^{(n)}$ in its canonical coordinate. They converge exponentially.

height of the approx. curve: $h_n = B(\alpha_n)$

mapping $\Omega_{f,k_1}^{(1)} \rightarrow \Omega_{Rf}^{(0)}$ is uniformly expanding

(with respect to Poincaré metrics of $\Omega_f^{(0)}$ and $\Omega_{Rf}^{(0)}$)



One only needs to work in model space (i.e. truncated checkerboard pattern)

Happy Birthday, Jack!